

# Non-Conservative $H^{\frac{1}{2}-}$ Weak Solutions of the Incompressible 3D Euler Equations

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# Main Result

We consider the incompressible 3D Euler equations for an unknown velocity  $u : \mathbb{R} \times \mathbb{T}^3 \rightarrow \mathbb{R}^3$  and pressure  $p : \mathbb{R} \times \mathbb{T}^3 \rightarrow \mathbb{R}$

$$\begin{aligned}\partial_t u + \operatorname{div}(u \otimes u) + \nabla p &= 0 \\ \operatorname{div} u &= 0\end{aligned}$$

## Theorem

Fix  $\beta \in (0, 1/2)$ . For any divergence-free, mean-zero vector fields  $v_{\text{start}}, v_{\text{end}} \in L^2(\mathbb{T}^3)$ , any  $T > 0$ , and any  $\varepsilon > 0$ , there exists a weak solution  $v \in C([0, T]; H^\beta(\mathbb{T}^3))$  to the 3D Euler equations such that  $\|v(0) - v_{\text{start}}\|_{L^2}, \|v(T) - v_{\text{end}}\|_{L^2} < \varepsilon$ .

**Motivation:** K41, intermittency (see Vlad's talk!)

**Goal of the talk:** (1) - Setup and Main Difficulties, (2) - Key heuristic estimates, (3) - Putting everything together

# Setup and Main Difficulties: Inductive Assumptions

- Induction on  $q$ :  $(u_q, \mathring{R}_q)$  solve the Euler-Reynolds system

$$\begin{aligned}\partial_t u_q + \operatorname{div}(u_q \otimes u_q) + \nabla p_q &= \operatorname{div} \mathring{R}_q \\ \operatorname{div} u_q &= 0\end{aligned}$$

- Spatial frequency parameter  $\lambda_q = a^{(b^q)}$
- Amplitude parameter  $\delta_q = \lambda_q^{-2\beta}$
- Very *heuristic* inductive assumptions:

$$\left\| \mathring{R}_q \right\|_{C_t^0 L_x^1} \leq \delta_{q+1}, \quad \left\| \nabla u_q \right\|_{C_t^0 L_x^2} \leq \delta_q^{1/2} \lambda_q, \quad \left\| w_{q+1} \right\|_{C_t^0 L_x^2} \leq \delta_{q+1}^{1/2}$$

# Set-Up and Main Difficulties: The Increment

- o Addition of the increment  $w_{q+1}$  shows that

$$\begin{aligned}\operatorname{div} \dot{R}_{q+1} - \nabla(\rho_{q+1} - \rho_q) &= \operatorname{div}(w_{q+1} \otimes w_{q+1} + \dot{R}_q) \\ &\quad + (\partial_t + u_q \cdot \nabla)w_{q+1} \\ &\quad + w_{q+1} \cdot \nabla u_q\end{aligned}$$

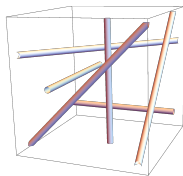
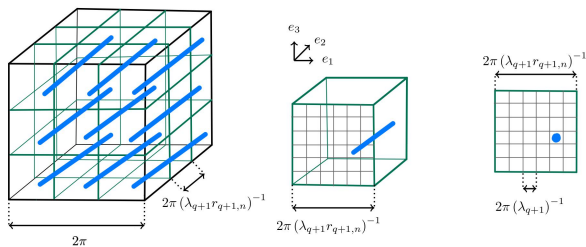
- o Ansatz for  $w_{q+1}$  is

$$w_{q+1} \approx \sum_{\xi} a_{\xi} \left( \dot{R}_q \right) \mathbb{W}_{\xi, r_{q+1}}$$

- o  $a_{\xi}$  is such that

$$\sum_{\xi} a_{\xi}^2 \left( \dot{R}_q \right) \xi \otimes \xi = \rho(x, t) \operatorname{Id} - \dot{R}_q$$

# Set-Up and Main Difficulties: Intermittent Mikado/Pipe Flows



## Set-Up and Main Difficulties: Type 1 Oscillation Error

$$\begin{aligned}\operatorname{div} \left( w_{q+1} \otimes w_{q+1} + \dot{R}_q \right) &= \operatorname{div} \left( \sum_{\xi} a_{\xi}^2(\dot{R}_q) \mathbb{W}_{\xi, r_{q+1}} \otimes \mathbb{W}_{\xi, r_{q+1}} + \dot{R}_q \right) \\ &= \operatorname{div} \left( \sum_{\xi} a_{\xi}^2(\dot{R}_q) \mathbb{P}_{=0} \left( \mathbb{W}_{\xi, r_{q+1}} \otimes \mathbb{W}_{\xi, r_{q+1}} \right) + \dot{R}_q \right) \\ &\quad + \sum_{\xi} \nabla a_{\xi}^2(\dot{R}_q) \mathbb{P}_{\neq 0} \left( \mathbb{W}_{\xi, r_{q+1}} \otimes \mathbb{W}_{\xi, r_{q+1}} \right)\end{aligned}$$

## Set-Up and Main Difficulties: The Transport Error

$$\begin{aligned}(\partial_t + u_q \cdot \nabla) \left( a_\xi \left( \mathring{R}_q \right) \mathbb{W}_{\xi, r_{q+1}} \right) &= (\partial_t + u_q \cdot \nabla) \left( a_\xi \left( \mathring{R}_q \right) \right) \mathbb{W}_{\xi, r_{q+1}} \\ &\quad + a_\xi \left( \mathring{R}_q \right) (\partial_t + u_q \cdot \nabla) \left( \mathbb{W}_{\xi, r_{q+1}} \right)\end{aligned}$$

## Set-Up and Main Difficulties: Type 2 Oscillation Error

$$\operatorname{div} (a_{\xi} a_{\xi'} \mathbb{W}_{\xi, r_{q+1}}(\Phi_{\xi}) \otimes \mathbb{W}_{\xi', r_{q+1}}(\Phi_{\xi'}))$$



# Heuristic Estimates: Velocity and Stress Cutoff Functions

- **Cutoffs:** We need knowledge of  $\nabla u_q$  to determine a timescale for Lagrangian coordinate systems, and we need knowledge of the local size of  $\dot{R}_q$  to enact cancellation
- **Heuristic Definition:** Level set decomposition
  
- **Issues:** How to estimate material derivatives?

## Heuristic Estimates: Three General Principles

- $\mathring{R}_q$  lives at frequency  $\lambda_q = \lambda_{q+1} \cdot \frac{\lambda_q}{\lambda_{q+1}}$  and has amplitude  $\delta_{q+1} = \frac{\delta_{q+1} \lambda_q}{\lambda_{q+1} \cdot \frac{\lambda_q}{\lambda_{q+1}}}$
- $\mathring{R}_{q+1}$  lives at frequency  $\lambda_{q+1} = \lambda_{q+1} \cdot 1$  and has amplitude  $\delta_{q+2} = \frac{\delta_{q+1} \lambda_q}{\lambda_{q+1} \cdot 1}$
- “Intermediate” errors satisfy “intermediate” estimate

**Goal:** Build a *sub*-perturbation  $w_{q+1,n}$  which cancels  $\mathring{R}_{q,n}$  and produces errors which either are small enough to be absorbed into  $\mathring{R}_{q+1}$  or satisfy estimates as before but with  $r_{q+1,n'} > r_{q+1,n}$ .

## Heuristic Estimates: 3<sup>rd</sup> General Principle (Type 2 Osc.)

$$\left\| \nabla^M \mathring{R}_{q,n} \right\|_{L^1} \leq \frac{\delta_{q+1} \lambda_q}{\lambda_{q+1} r_{q+1,n}} (\lambda_{q+1} r_{q+1,n})^M$$

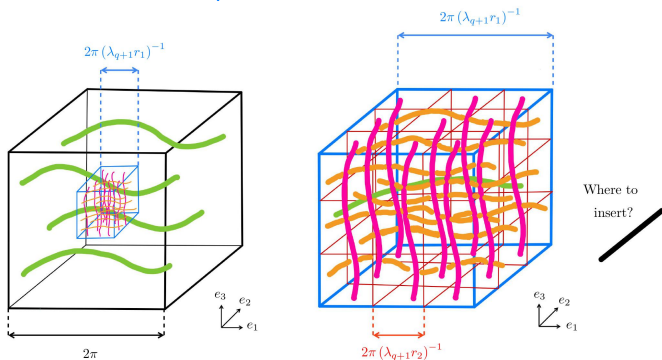
**Claim:** Correcting this error by pipe flows with intermittency  $r_{q+1,n'}$  for  $r_{q+1,n'} = r_{q+1,n}^{3/4}$  produces *new* Type 2 oscillation errors which vanish

**Proof:** Strategy - Intermittency parameter  $r_{q+1,n'}$  allows for  $(r_{q+1,n'})^{-2}$  disjoint placements

Step 1 - Checkerboard Cutoff Functions and “Pipe Wavelets”

# Heuristic Estimates: 3<sup>rd</sup> General Principle (Type 2 Osc. cont'd)

## Diagram of Stress Decomposition and Pipe Wavelets



## Heuristic Estimates: 3<sup>rd</sup> General Principle (Type 2 Osc. cont'd)

$$\left\| \nabla^M \hat{R}_{q,n} \right\|_{L^1} \leq \frac{\delta_{q+1} \lambda_q}{\lambda_{q+1} r_{q+1,n}} (\lambda_{q+1} r_{q+1,n})^M$$

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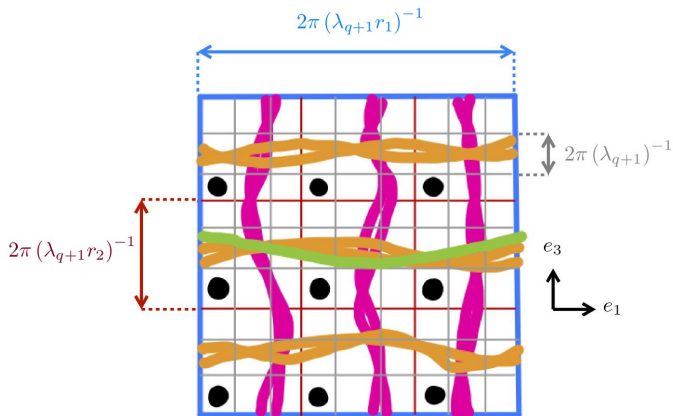
**Proof: (cont'd)**

Step 2 - Counting existing pipes

How many pipes which were periodized to scale  $(\lambda_{q+1} r_{q+1,n'})^{-1}$  and then deformed by a Lipschitz diffeomorphism inhabit a set of diameter  $(\lambda_{q+1} r_{q+1,n})^{-1}$ ?

# Heuristic Estimates: 3<sup>rd</sup> General Principle (Type 2 Osc. cont'd)

## Projection and Counting of Existing Pipes



## Heuristic Estimates: 3<sup>rd</sup> General Principle (Type 2 Osc. cont'd)

$$\left\| \nabla^M \hat{R}_{q,n} \right\|_{L^1} \leq \frac{\delta_{q+1} \lambda_q}{\lambda_{q+1} r_{q+1,n}} (\lambda_{q+1} r_{q+1,n})^M$$

**Claim:** Correcting this error by pipe flows with intermittency  $r_{q+1,n'}$  for  $r_{q+1,n'} = r_{q+1,n}^{3/4}$  produces new Type 2 oscillation errors which vanish

**Proof: (cont'd)**

Step 3 - Covering a single existing pipe

How many grid squares of size  $\lambda_{q+1}^{-1}$  does it take to cover a deformed pipe of thickness  $\lambda_{q+1}^{-1}$  which lives in a set of diameter  $(\lambda_{q+1} r_{q+1,n})^{-1}$ ?

## Heuristic Estimates: 3<sup>rd</sup> General Principle (Type 2 Osc. cont'd)

$$\left\| \nabla^M \hat{R}_{q,n} \right\|_{L^1} \leq \frac{\delta_{q+1} \lambda_q}{\lambda_{q+1} r_{q+1,n}} (\lambda_{q+1} r_{q+1,n})^M$$

**Claim:** Correcting this error by pipe flows with intermittency  $r_{q+1,n'}$  for  $r_{q+1,n'} = r_{q+1,n}^{3/4}$  produces *new* Type 2 oscillation errors which vanish

**Proof: (cont'd)**

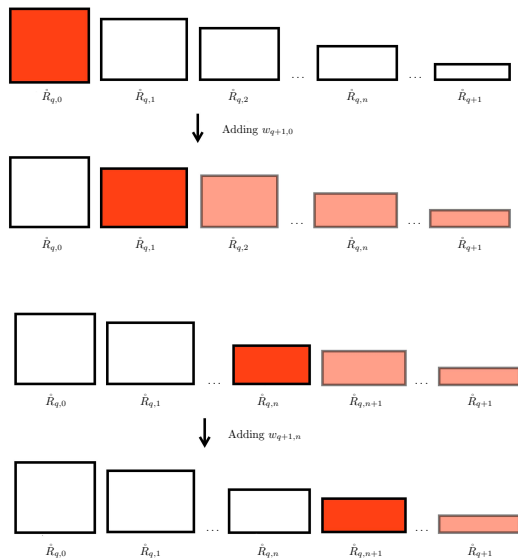
Step 4 - Pigeonhole principle and relative intermittency inequality

How many grid squares did I use to cover the existing pipes, and is this less than the number of options I have for placing the new set?



# Putting Everything Together - Adding Sub-Perturbations

$w_{q+1,n}$



Thank you for your attention!

# Difficulties and Solutions

1. Low-frequency, high-amplitude oscillation errors due to intermittency
1. Higher order stresses and correctors - “two (or more) bites at the apple”

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  2. Sharp higher order material and spatial derivatives, and sharp local timescale information
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# Difficulties and Solutions

1. Low-frequency, high-amplitude oscillation errors due to intermittency
  2. Sharp higher order material and spatial derivatives, and sharp local timescale information
  3. Pipes from overlapping Lagrangian coordinate systems or different “bites at the apple” may intersect
1. Higher order stresses and correctors - “two (or more) bites at the apple”
  2. Cutoff functions which act as a Littlewood-Paley projection in Eulerian and Lagrangian coordinates, with info about position and amplitude
  3. Intermittency allows for shifts, and thus a placement technique

# General Principles

Let  $r_n \in [\lambda_q/\lambda_{q+1}, 1]$ , and let  $\mathring{R}_{q,n}$  be an error living at frequency  $\lambda_{q+1}r_{q+1,n}$  and of size  $\frac{\delta_{q+1}\lambda_q}{\lambda_{q+1}r_{q+1,n}}$  in  $L^1$ .

$$\leftarrow w_{q+1,n} = \sum_{\mathring{R}_{q,n}} a_{\mathring{R}_{q,n}} W_{\mathring{R}_{q,n}, r_{q+1,n}}$$

1. **Effect of intermittency on Type 1:** Correcting this error by a pipe flow with intermittency  $r_{q+1,n'} > r_{q+1,n}$  produces a new Type 1 oscillation error with minimum frequency  $\lambda_{q+1}r_{q+1,n'}$  and size  $\frac{\delta_{q+1}\lambda_q}{\lambda_{q+1}r_{q+1,n'}}$ .

- Want less intermittency!

- Goal:  $r_{q+1,n'} \rightarrow 1$

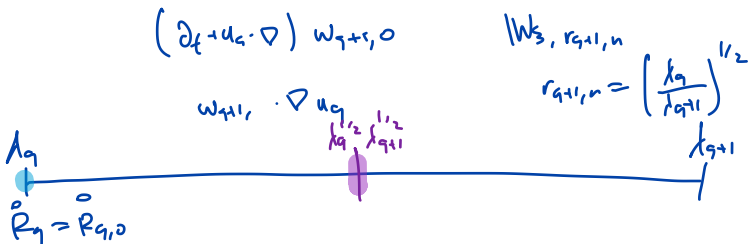
$\Rightarrow$  everything ends up in  $\mathring{R}_{q,n}$ !

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$$W_{q+1,n} = \sum_3 a_3 (P_3^0) |W_3, r_{q+1,n}'$$

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2. **Effect of intermittency on Nash/transport:** Correcting this error by a pipe flow with intermittency  $r_{q+1,n'} = r_{q+1,n}^{1/2}$  produces *new* Nash and transport errors of size  $\delta_{q+2}$  which are commensurate with  $H^{1/2}$  regularity.



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- 3. Effect of intermittency on Type 2:** Correcting this error by pipe flows with intermittency  $r_{q+1,n'} = r_{q+1,n}^{3/4}$  ensures that the *new* Type 2 errors vanish.