

# Lecture 1 : motivation

Euclidean QFT: ↗ goal

probability measures  $\mu$  on  $S'(\mathbb{R}^d)$

Stochastic quantization: ↗ tool

Gaussian  $W \xrightarrow{F_\mu} \phi$

law  $\phi = \mu$

(e.g.  $\phi$  solves some stoch diff equ with  $W$ )

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- physical system: {states}      {observables}
- Classical mech:  $\underset{\downarrow}{\text{Eg: } (q, p)}$       functions like  $E = q^2 + p^2$ .
- Quantum mechanics :

Hilbert space  $H := \{ \text{states } A \}$

Operators = {observables  $w$ }

"measurement":  $w(A) = \langle w, Aw \rangle$

- Dynamics :

Time translation <sup>unitary</sup> group  $U(t)$

$U(t) = e^{-itE}$  where  $E$ : hamiltonian/energy

$w(t) = U(t)w$  "Schrödinger picture"

$A(t) = U(t)^{-1}A U(t)$  "Heisenberg picture"

relation:  $w(t)(A) = w(A(t))$  Exercise

$$\begin{aligned} & \langle w(t), Aw(t) \rangle \\ &= \langle U(t)w, A U(t)w \rangle \\ &= \langle w, A(t)w \rangle \end{aligned}$$

key quantities:  $w(A_1(t_1) \dots A_n(t_n))$

Assume:  $\exists \varphi \in \mathcal{H}$  (vacuum, ground state)

$$U(t)\varphi = \varphi$$

$w(A_1(t_1) \dots A_n(t_n))$

$$= \langle \varphi, U(t_1)^{-1} A_1 U(t_1) \dots U(t_n)^{-1} A_n U(t_n) \varphi \rangle$$

$$= \langle \varphi, A_1 U(t_1 - t_2) A_2 \dots U(t_{n-1} - t_n) A_n \varphi \rangle$$

Example:  $\mathcal{H} = L^2(\mathbb{R}, \rho)$   $\rho \sim N(0, \frac{1}{2a})$

operators: e.g. multiplication by  $C(\mathbb{R})$

To describe a dynamic, recall:

Hermite polynomials  $H_n(x)$  is orthogonal basis of  $\mathcal{H}$ .

$$U(t) H_n = e^{-i\alpha n t} H_n \quad \text{check: normalized?}$$

$$\text{Ground state } \varphi = H_0 = 1$$

one can also describe this dynamic as follows:

$$U(t) = e^{-it\bar{E}} \quad E \text{ is energy operator}$$

where  $\bar{E}$  is an operator with is diagonalized by  $H_n$

$$E H_n = \alpha_n H_n$$

TA session:  $E$  is ODE operator  $-\partial_x^2 + x\partial_x$  (up to constants)

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"Euclidean" = "going to Imaginary time"

replace  $t$  by  $-it$ .  $e^{-it\bar{E}} \rightarrow e^{-tE}$

"Wick rotation"  
"analytic continuation"

( $t \geq 0$ )

Continuous Semigroup of Contractions

$$\langle \varphi, A_1 K(t_1, t_2) A_2 \dots, K(t_{n-1}, t_n) A_n \varphi \rangle$$

"Schwinger functions"

- ① Can they recover quantum system?
- ② They are correlation functions of a prob measure.  
(of a random function)

$$\text{Ex: } dX_t = -\alpha X_t dt + dB_t$$

$$X_0 \sim N(0, \frac{1}{2\alpha}) \quad t \in \mathbb{R}$$

Prob measure  $\mu$  on  $C(\mathbb{R})$

$$X_t = \int_{-\infty}^t e^{-\alpha(t-s)} dB_s \quad \checkmark \text{ random field on } \mathbb{R}$$

$$\begin{aligned} E(X_t X_{t'}) &= E\left[\int_0^t e^{-\alpha(t-s)} dB_s \int_0^{t'} e^{-\alpha(t'-s)} dB_s\right] \\ t > t' &= e^{-\alpha(t-t')} \int_{-\infty}^{t'} e^{-2\alpha(t'-s)} ds = \frac{e^{-\alpha(t-t')}}{2\alpha} \end{aligned}$$

$$\begin{aligned} E[f(X_t) | X_0 = x] &= E\left[f\left(e^{-\alpha t} x + N(0, \frac{1-e^{-2\alpha t}}{2\alpha})\right)\right] \\ &=: (K(t)f)(x) \quad \text{Semigroup} \end{aligned}$$

Compute:

$$\begin{aligned} K(t) e^{i\lambda x} &= E\left[e^{i\lambda(e^{-\alpha t} x + N(0, \frac{1-e^{-2\alpha t}}{2\alpha}))}\right] \\ &= e^{i\lambda e^{-\alpha t} - \frac{\lambda^2(1-e^{-2\alpha t})}{2\alpha}} \end{aligned}$$

$$\begin{aligned} \text{Thus } K(t) \underbrace{e^{i\lambda x + \frac{\lambda^2}{4\alpha}}}_{\sum_{n \geq 0} \frac{(i\lambda)^n}{n!} H_n(x)} &= e^{i\tilde{\lambda} x + \frac{\tilde{\lambda}^2}{4\alpha}} \\ (\tilde{\lambda} = \lambda e^{-\alpha t}) \quad \underbrace{\quad}_{\sum_{n \geq 0} \frac{(i\lambda)^n}{n!} e^{-\alpha nt} H_n} \end{aligned}$$

$$\Rightarrow K(t) H_n = e^{-\alpha n t} H_n$$

This is exactly  $U(t)$  in imaginary time!

More interestingly, correlations are just Schwinger functions!

For  $A_1, A_2 \dots A_n \in C(\mathbb{R})$  (observables)

$$\begin{aligned}
 & E[A_1(X_{t_1}) \dots A_n(X_{t_n})] \\
 &= E[A_1(X_{t_1}) \dots E[A_n(X_{t_n}) | X_{t_{n-1}}]] \\
 &= E[A_1(X_{t_1}) \dots \underbrace{A_{n-1}(X_{t_{n-1}})}_{\text{---}} \underbrace{K(t_n - t_{n-1})}_{\text{---}} \underbrace{A_n(X_{t_{n-1}})}_{\text{---}}] \\
 &= \dots \\
 &= E[(A_1 K(t_2 - t_1) A_2 \dots K(t_n - t_{n-1}) A_n)(X_{t_1})] \\
 &= \langle \varphi, A_1 K(t_2 - t_1) \dots K(t_n - t_{n-1}) A_n \varphi \rangle_{L^2(\mathbb{R}, f)} \quad \begin{matrix} \text{Stationary} \\ f = N(0, \frac{1}{2\alpha}) \end{matrix}
 \end{aligned}$$

In general, Schwinger functions must satisfy some properties to recover quantum system.

e.g.:  $\langle \varphi, A_2 K(t) A_1 A_1 K(t) A_2 \varphi \rangle \geq 0$

"reflection positivity" (RP)

Osterwalder-Schrader theorem (baby version)

RP + certain regularity property  $\xrightarrow{\text{recover}} \text{quantum data}$   
 $(H, U(t) \text{ etc})$

$$RP: \mathbb{E} [\overline{\Theta F(X)} F(X)] \geq 0 \quad \forall \text{ functionals } F$$

$\uparrow f\text{lip } X: X(t) \rightarrow X(t-t)$        $F(X) = \sum_k c_k e^{i\lambda_k X(t_k)}$

Remark: above OU field  $X$  is RP.

Beyond Gaussian:

$$\mathbb{Q}_T(d\omega) = \frac{1}{Z_T} e^{-\int_{-T}^T V(X_s(\omega)) ds} \underbrace{P(d\omega)}_{\text{OU}}$$

reflection positive:

$$G = e^{-\int_0^T V(X_s) ds}$$

$$\Theta G = e^{-\int_0^T V(X_{-s}) ds} \quad \boxed{\mathbb{E}_{\mathbb{Q}_T} (\overline{\Theta F} F) \stackrel{?}{\geq} 0}$$

$$e^{-\int_{-T}^T V(X_s) ds} = \Theta G \cdot G = \overline{\Theta G} \cdot G$$

$$\mathbb{E}_{\mathbb{Q}_T} [\overline{\Theta F} F] = \frac{1}{Z_T} \mathbb{E}_P [\overline{\Theta F} F \underbrace{e^{-\int_{-T}^T V(X_s(\omega)) ds}}_{\overline{\Theta G} \cdot G}]$$

$$= \frac{1}{Z_T} \mathbb{E}_P [\overline{\Theta(FG)} FG] \stackrel{\text{RP of } P}{\geq} 0$$