

Lecture 5: Converge to equilibrium

Last time we showed
existence of infinite volume ϕ^f measure
by tightness limit of μ_n as $n \rightarrow \infty$

To show uniqueness of infinite volume ϕ^f measure.
One could argue by the following steps

Step 1: On the whole \mathbb{Z}^d
$$dX = (AX - X^3) dt + dB \quad (A = \Delta - m^2)$$
has unique (strong) solution.

(Write its semigroup by P_t)

Step 2: Invariant measure of P_t is unique.
(call it μ)

Step 3: Every tight limit of μ_n must be μ .

Step 1 is "Standard":

(e.g. tight limit in $C(\mathbb{R}_+, \mathbb{R}^{\mathbb{Z}^d})$
identify it as solution to the right martingale problem
Pathwise uniqueness by Gronwall type argument)

Step 3: just argue that every tight limit of μ_n is an invariant meas of P_t

Step 2 is the "crucial step" and relies on some additional structure of P_t
We will show "convexity" is sufficient

To better understand the role of convexity

consider general $dX = (AX - \frac{i}{\hbar} U'(X))dt + dB$

corresponding to $e^{-\frac{i}{\hbar} \int (\nabla\phi)^2 + m^2\phi^2 + U(\phi)} \prod_x d\phi_x$

Imagine two solutions Z_1, Z_2 .

with initial conditions $Z_1(0), Z_2(0)$

$$\begin{aligned} \partial_t Z_1 &= AZ_1 - \frac{1}{2}U'(X_1) & X_j &= Y_j + Z_j \\ \partial_t Z_2 &= AZ_2 - \frac{1}{2}U'(X_2) & j &= 1, 2. \end{aligned}$$

$$H = Z_2 - Z_1$$

$$\begin{aligned} \partial_t H - AH &= \frac{1}{2}(U'(X_1) - U'(X_2)) \\ &= \frac{1}{2}(U'(X_1) - U'(X_1 + H + K)) \stackrel{\text{def}}{=} Q \end{aligned}$$

where $K = Y_2 - Y_1$

$$Q = -\frac{1}{2} \int_0^1 d\tau U''(X_1 + \tau(H+K))(H+K)$$

Assume: $U''(\varphi) \geq -2\chi$ for some $\chi \in \mathbb{R}_+$

e.g. $U(\varphi) = \beta\varphi^2 + \lambda\varphi^4 \quad \lambda > 0$

This implies $\varphi \mapsto U(\varphi) + 2\chi \sum \varphi^2$ is convex.

$$\text{Let } G \stackrel{\text{def}}{=} \frac{1}{2} \int_0^1 d\tau U''(X_1 + \tau(H+K)) \geq -\chi$$

Then $Q = -G(H+K)$

Test eqn with $p^2 H$ for some p

RHS:

$$\begin{aligned} \sum p^2 H \cdot Q &= -\sum p^2 H \cdot G H - \sum p^2 H \cdot G K \\ &\leq C_\delta \|P G K\|_{L^2}^2 + \delta \|P H\|_{L^2}^2 + \sum p^2 H^2 \cdot \chi \\ &= C_\delta \|P G K\|_{L^2}^2 + (\delta + \chi) \|P H\|_{L^2}^2 \end{aligned}$$

Take $\delta \leq \chi$

$$\frac{1}{2} \partial_t \|P H\|_{L^2}^2 + \sum p^2 H (m^2 - \Delta) H \leq C \|P G K\|_{L^2}^2 + 2\chi \|P H\|_{L^2}^2$$

Consider $F(t) = e^{ct} \|P H\|_{L^2}^2$.

Then,

$$\begin{aligned} \frac{1}{2} \partial_t (e^{ct} \|P H\|_{L^2}^2) &= \frac{c}{2} e^{ct} \|P H\|_{L^2}^2 + \frac{1}{2} e^{ct} \partial_t \|P H\|_{L^2}^2 \\ (1) \leq -e^{ct} \sum p^2 H (m^2 - 2\chi - \frac{c}{2} - \Delta) H &+ C e^{ct} \|P G K\|_{L^2}^2 \end{aligned}$$

Let $\tilde{Q}(t) = C \|P G(t) K(t)\|_{L^2}^2$

choose $p(x) = e^{-\theta|x|}$

For $|\varepsilon \theta \leq 1$

$$|\nabla(p^2)| \leq 2\theta p^2$$

↑
check!

$$\underbrace{\sum_{\wedge} \rho^2 H(-\Delta)H} = \sum_{\wedge} \nabla H \nabla(\rho^2 H) = \sum_{\wedge} \nabla H (\nabla \rho^2 \cdot H + \rho^2 \nabla H)$$

$$\Rightarrow \sum_{\wedge} (\nabla H)^2 \cdot \rho^2 - \underbrace{40 \sum_{\wedge} |H| \cdot |H| \cdot \rho^2}_{\leq c\theta^2 \sum_{\wedge} \rho^2 H^2 + \frac{1}{2} \sum_{\wedge} \rho^2 (\nabla H)^2} \quad (2)$$

$$\Rightarrow \frac{1}{2} \sum_{\wedge} \rho^2 (\nabla H)^2 - c\theta^2 \sum_{\wedge} \rho^2 H^2$$

Then

$$\frac{1}{2} \partial_t (e^{ct} \|PH\|_{\ell^2}^2) + e^{ct} (m^2 - 2\chi - \frac{c}{2} - c\theta^2) \sum_{\wedge} \rho^2 H^2$$

$$+ \frac{1}{2} e^{ct} \sum_{\wedge} \rho^2 (\nabla H)^2$$

$$\stackrel{(2)}{\leq} \frac{1}{2} \partial_t (e^{ct} \|PH\|_{\ell^2}^2) + e^{ct} \sum_{\wedge} \rho^2 H (m^2 - \frac{c}{2} - 2\chi) H$$

$$+ e^{ct} \sum_{\wedge} \rho^2 H (-\Delta) H$$

$$\stackrel{(1)}{\leq} e^{ct} \tilde{Q}(t)$$

Assume $\boxed{m^2 - 2\chi - \frac{c}{2} - c\theta^2 \geq 0}$?

$$\frac{1}{2} \partial_t (e^{ct} \|PH\|_{\ell^2}^2) \leq e^{ct} \tilde{Q}(t)$$

$$\Rightarrow \frac{1}{2} e^{ct} \|PH\|_{\ell^2}^2 \leq \frac{1}{2} \|PH(0)\|_{\ell^2}^2 + \int_0^t e^{cs} \tilde{Q}(s) ds$$

$$\Rightarrow \|PH(t)\|_{\ell^2}^2 \leq e^{-ct} \|PH(0)\|_{\ell^2}^2 + 2 \int_0^t e^{-c(t-s)} \tilde{Q}(s) ds$$

(★)

Where

$$\begin{aligned}\tilde{Q}(s) &= C \|P G(s) K(s)\|_{\ell^2}^2 \\ &= C \sum_{\wedge} P^2 \left(\int_0^1 d\tau U''(X_{\wedge}^{(s)} + \tau(H^{(s)} + K^{(s)})) \right)^2 K(s)^2 \\ &\leq C \int_0^1 \sum_{\wedge} P^2 \cdot U''(X_{\wedge}^{(s)} + \tau(H^{(s)} + K^{(s)}))^2 \cdot K(s)^2 d\tau\end{aligned}$$

need to explain how two processes are coupled.

$$\mathbb{E} \tilde{Q}(s) \leq \sum_{\wedge} \int_0^1 P^2 \cdot \mathbb{E} \left[U''(X_{\wedge}^{(s)} + \tau(H^{(s)} + K^{(s)}))^4 \right]^{\frac{1}{2}} \mathbb{E} \left[K(s)^4 \right]^{\frac{1}{2}} d\tau$$

$\mathbb{E} [\text{poly}(X_{\wedge} + \tau(H + K))^4] \leq C$ uniform in $M \rightarrow \infty$

$$\mathbb{E} \tilde{Q}(s) \leq \sum_{\wedge} P^2 \cdot \mathbb{E} [K(\omega)^4]^{\frac{1}{2}} \quad (K \text{ in stationary})$$

By (*) we have:

Lem: Under above assumptions.

$$\mathbb{E} \|P H_t\|_{\ell^2}^2 \leq e^{-ct} \mathbb{E} \|P H(\omega)\|_{\ell^2}^2 + C \sum_{\wedge} P^2 \mathbb{E} [K(\omega)^4]^{\frac{1}{2}}$$

Remark: As $t \rightarrow \infty$

$$\text{If } \sum_{\wedge} P^2 \mathbb{E} [K(\omega)^4]^{\frac{1}{2}} \rightarrow 0.$$

$$\text{We have } \mathbb{E} \|P H_t\|_{\ell^2}^2 \rightarrow 0$$

Write P_t : semigroup for dynamic X on \mathbb{Z}^d .

Suppose, μ, ν are invariant measures of P_t .

One can find $\mu_m \rightarrow \mu$

$\nu_m \rightarrow \nu$

We have solution X_m to SDE starting from μ_m
Semigroup P_t^M

Similar as step 1. $(X_m)_m$ is tight. Limit satisfy SDE on \mathbb{Z}^d
with i.c. μ

By Step 1. $X_m \rightarrow X$

So, $\int P_t^M F d\mu_m \longrightarrow \int P_t F d\mu \quad (M \rightarrow \infty)$

\parallel
 $E F(X_m^{(t)}) \longrightarrow E F(X(t)) \quad (\star)$

Similarly: $\int P_t^M F d\nu_m \longrightarrow \int P_t F d\nu \quad (M \rightarrow \infty)$

Therefore

$$\left| \int F d\mu - \int F d\nu \right|$$

$$= \left| \int P_t F d\mu - \int P_t F d\nu \right| \quad (\text{invariance of } \mu, \nu)$$

$$= \lim_{M \rightarrow \infty} \left| \int P_t^M F d\mu_m - \int P_t^M F d\nu_m \right| \quad (\text{by } (\star))$$

$$= \lim_{M \rightarrow \infty} \inf_{\pi \in C(\mu_M P_t^M, \nu_M P_t^M)} \left| \int (\tilde{F}(x) - \tilde{F}(y)) d\pi(x, y) \right|$$

$$\leq C \lim_{M \rightarrow \infty} W_2(\mu_M P_t^M, \nu_M P_t^M)$$

↓ Wasserstein distance

from the $t \rightarrow \infty$
decay result
shown above

$$W_2(\mu, \nu) \stackrel{\text{def}}{=} \inf_{\pi \in C(\mu, \nu)} E_{\pi}$$

$$\leq C e^{-ct} \lim_{M \rightarrow \infty} W_2(\mu_M, \nu_M) \rightarrow 0$$