

L^p estimates for $\bar{\Phi}_2^4$.

Ref: Shen-Zhu-Zhu paper on perturbation theory. Lemma 2.1.

Consider $\bar{\Phi}_2^4$ equation

$$\bar{\Phi} = z + \Upsilon,$$

$$(\partial_t - \Delta + m)z = 3$$

$$\begin{aligned} (\partial_t - \Delta + m)\Upsilon = & -\frac{1}{2}(\Upsilon^3 + 3z\Upsilon^2 \\ & + 3z^2\Upsilon \\ & + z^3) \end{aligned}$$

Here and throughout, omit ε parameter. (If you want, can imagine this all takes place in the continuum,

although we may not have covered the stochastic estimates for $z, :z^2:, :z^3:$ in the continuum.)

Simplification: unit volume, so that we can avoid dealing with weights.

Goal: get estimates on Y in terms of

$z, :z^2:, :z^3:,$

which we control using probability theory.

Lemma. Let $p \geq 1$. We have that

$$\frac{1}{2p} \partial_t (\|Y\|_{L^p}^{2p}) = -(2p-1) \|Y^{p-1} \nabla Y\|_{L^2}^2 - \|Y^{p-2} \nabla(|Y|^2)\|_{L^2}^2$$

$$-m \|Y\|_{L^{2p}}^{2p} - \frac{\lambda}{2} \|Y\|_{L^{2(p+1)}}^{2(p+1)}$$

$$- \frac{\lambda}{2} \int_{\mathbb{T}^2} G |Y|^{2p-1}$$

where $G := ZY^2 + :Z^2:Y + :Z^3:$.

Remark. To get estimates on Y , we will need to bound the bad terms by the good terms. The good terms are good b/c they come with a sign.

Note if $G \equiv 0$, then we would obtain that $\|Y\|_{L^{2p}}$ is decreasing in time.

Pf. We write

$$(\partial_t - \Delta + m)\psi = -\frac{1}{2}\psi^3 - \frac{1}{2}G.$$

Thus, we compute

$$\begin{aligned} \frac{1}{2p} \partial_t (\psi^{2p}) &= \psi^{2p-1} \partial_t \psi \\ &= \psi^{2p-1} \left(\Delta \psi - m\psi - \frac{1}{2}\psi^3 - \frac{1}{2}G \right) \\ &= \psi^{2p-1} \Delta \psi - m\psi^{2p} - \frac{1}{2}\psi^{2(p+1)} \\ &\quad - \frac{1}{2}G\psi^{2p-1}. \end{aligned}$$

We have that (implicit sum over j)

$$\begin{aligned} \frac{1}{2p} \Delta (\psi^{2p}) &= \frac{1}{2p} \partial_j^i \partial_j (\psi^{2p}) \\ &= \partial_j^i (\psi^{2p-1} \partial_j \psi) \\ &= (2p-1) \psi^{2p-2} \partial_j^i \psi \partial_j \psi \end{aligned}$$

$$\begin{aligned}
 & + \psi^{2p-1} \partial^j \partial_j \psi \\
 & = (2p-1) \psi^{2(p-1)} |\nabla \psi|^2 \\
 & \quad + \psi^{2p-1} \Delta \psi.
 \end{aligned}$$

Inserting this identity for $\psi^{2p-1} \Delta \psi$,
we obtain

$$\begin{aligned}
 \frac{1}{2p} (\partial_t - \Delta)(\psi^{2p}) & = -(2p-1) \psi^{2(p-1)} |\nabla \psi|^2 \\
 & - m \psi^{2p} - \frac{\lambda}{2} \psi^{2(p+1)} - \frac{1}{2} G \psi^{2p-1}.
 \end{aligned}$$

The desired result now follows by
integrating in space. \square

Now, suppose z, z^2, z^3 were in L^∞ . Then we could easily bound

$$\left| \int z Y^{2p+1} \right| \leq \|z\|_\infty \|Y\|_{L^{2p+1}}^{2p+1}$$

$$\leq \gamma \|Y\|_{L^{2p+2}}^{2p+2} + C_\gamma \|z\|_\infty^{2p+2}$$

(apply Young, i.e. $xy \leq \frac{x^q}{q} + \frac{y^{q'}}{q'}$ where

$$1 = \frac{1}{q} + \frac{1}{q'},$$
$$q = \frac{2p+2}{2p+1}, \quad q' = 2p+2).$$

Can handle other terms similarly,
and obtain

$$\partial_t \|Y\|_{L^{2p}}^{2p} + m \|Y\|_{L^{2p}}^{4p} + p \|Y^{p-1} \nabla Y\|_{L^{2p+2}}^2 + \lambda \|Y\|_{L^{2p+2}}^{2p+2} \leq C \lambda \max(\|z^1\|_{L^\infty}, \|z^2\|_{L^\infty}, \|z^3\|_{L^\infty}).$$

We would then have bounded the size of Y by the sizes of

$$z, z^2, z^3,$$

which was our goal.

However, z, z^2, z^3 are not in L^∞ . OTOH, they are "almost" in L^∞ , since they lie in

$$B_{\infty, \infty}^{-k} \text{ for any } k > 0.$$

Thus, we should be able to perturb our previous argument, by making use of the good term $\|Y^{p-1} \nabla Y\|_{L^2}^2$, which controls derivatives of Y . We turn

to this next. We will need the following

Analytic lemmas.

Lemma. Let $k > 0$. For $f \in \mathcal{B}_{\infty, \infty}^{-k}$, $g \in \mathcal{B}_{\infty, \infty}^k$, we have that

$$|\int fg| \lesssim \|f\|_{\mathcal{B}_{\infty, \infty}^{-k}} \|g\|_{\mathcal{B}_{1, 1}^k}.$$

Pf. Have that

$$\int fg = \sum_{|j-j'| \lesssim 1} \int \Delta_j f \Delta_{j'} g,$$

thus

$$|\int fg| \leq \sum_{|j-j'| \lesssim 1} \|\Delta_j f\|_{L^\infty} \|\Delta_{j'} g\|_{L^1}$$

$$\lesssim \|f\|_{B_{\infty, \infty}^{-k}} \sum_j 2^{jk} \|\Delta_j g\|_{L^1}$$

$$\lesssim \|f\|_{B_{\infty, \infty}^{-k}} \|g\|_{B_{1,1}^k}. \quad \square$$

Lemma. For $\alpha > 0$, we have that

$$\|f\|_{B_{1,1}^{-\alpha}} \lesssim_{\alpha} \|f\|_{L^1}.$$

Pf. For any j , we have that

$$\|\Delta_j f\|_{L^1} \lesssim \|f\|_{L^1}$$

(by, say, Young's convolution inequality, since

$$\Delta_j f = \check{\psi}_j * f, \text{ and}$$

\check{e}_j is L^1 -normalized. \square

Lemma. We have that

$$\|g\|_{B_{1,1}^k} \leq \|g\|_{B_{1,1}^{\frac{1}{2}-k}}^{\frac{1}{2}} \|g\|_{B_{1,1}^{3k}}^{\frac{1}{2}}$$

Pf. Cauchy-Schwarz. \square

Lemma. We have that

$$\|f\|_{B_{1,1}^{1-\gamma}} \leq \|\nabla f\|_{B_{1,1}^{-\gamma}} + \|f\|_{L^1}.$$

Pf. By Bernstein's inequality, for $j \geq 0$, we have that

$$\|\Delta_j f\|_{L^1} \leq 2^j \|\Delta_j \nabla f\|_{L^1}.$$

Thus

$$\|f\|_{\mathcal{B}_{1,1}^{-\gamma}} \lesssim \|f\|_{L^1}$$

$$+ \sum_{j \geq 0} 2^{j(1-\gamma)} 2^{-j} \|\Delta_j \nabla f\|_{L^1}$$

$$\lesssim \|f\|_{L^1} + \|\nabla f\|_{\mathcal{B}_{1,1}^{-\gamma}}. \quad \square.$$

With these lemmas in hand, we return to estimating Υ . First,

$$\lambda \left| \int z \Upsilon^{2p+1} \right|$$

$$\lesssim \lambda \|z\|_{\mathcal{B}_{0,\infty}^{-k}} \|\Upsilon^{2p+1}\|_{\mathcal{B}_{1,1}^k}$$

$$\lesssim \lambda \|z\|_{B_{\infty, \rho}^{-k}} \|Y^{2p+1}\|_{B_{1,1}^{-k}}^{\frac{1}{2}} \|Y^{2p+1}\|_{B_{1,1}^{3k}}^{\frac{1}{2}}$$

$$\lesssim \lambda \|z\|_{B_{\infty, \rho}^{-k}} \|Y^{2p+1}\|_{L^2}^{\frac{1}{2}} \|Y^{2p+1}\|_{B_{1,1}^{3k}}^{\frac{1}{2}}$$

$$\leq \lambda \|z\|_{B_{\infty, \rho}^{-k}} \|Y\|_{L^{2p+2}}^{\frac{2p+1}{2}} \|Y^{2p+1}\|_{B_{1,1}^{3k}}^{\frac{1}{2}}$$

We estimate

$$\|Y^{2p+1}\|_{B_{1,1}^{3k}} \leq \|\nabla(Y^{2p+1})\|_{L^1} + \|Y^{2p+1}\|_{L^1}$$

$$\lesssim p \|Y^{2p} \nabla Y\|_{L^1} + \|Y^{2p+1}\|_{L^1}$$

$$\leq P \| \psi^{p+1} \|_{L^2} \| \psi^{p-1} \nabla \psi \|_{L^2} \\ + \| \psi \|_{L^{2p+2}}^{2p+1} \cdot$$

From this, obtain

$$\lambda \left| \int z \psi^{2p+1} \right|$$

$$\leq P \lambda \| z \|_{B_{\infty,0}^{-k}} \| \psi \|_{L^{2p+2}}^{\frac{2p+1}{2}} \| \psi \|_{L^{2p+2}}^{\frac{p+1}{2}} \\ \| \psi^{p-1} \nabla \psi \|_{L^2}^{\frac{1}{2}}$$

$$+ \lambda \| z \|_{B_{\infty,0}^{-k}} \| \psi \|_{L^{2p+2}}^{2p+1} \cdot$$

Estimate by Young

$$\text{[Yellow highlight]} \leq \gamma \lambda \|Y\|_{L^{2p+2}}^{2p+2} + C \gamma \lambda \|Z\|_{B_{\infty,1}^{-\kappa}}^{2p+2}.$$

Estimate

$$\text{[Purple highlight]} = p \lambda \|Z\|_{B_{\infty,1}^{-\kappa}} \|Y\|_{L^{2p+2}}^{\frac{3p+2}{2}} \|Y^{p-1} \nabla Y\|_{L^2}^{\frac{1}{2}}$$

$$\leq \gamma p \|Y^{p-1} \nabla Y\|_{L^2}^2 \left(\begin{array}{l} g = \frac{4}{3} \\ g' = 4 \end{array} \right) + C \lambda^{\frac{4}{3}} p \|Z\|_{B_{\infty,1}^{-\kappa}}^{\frac{4}{3}} \|Y\|_{L^{2p+2}}^{2p+2 - \frac{2}{3}}.$$

$$\leq \gamma p \|\Upsilon^{p-1} \nabla \Upsilon\|_{L^2}^2$$

$$\left(\begin{array}{l} g = \frac{2p+2}{2p+2-\frac{2}{3}} \\ g' = \frac{2p+2}{\frac{2}{3}} \end{array} \right)$$

$$+ \gamma \lambda \|\Upsilon\|_{L^{2p+2}}^{2p+2}$$

$$+ C p^{3(p+1)} \lambda^{\left(\frac{4}{3} - \frac{1}{g'}\right)g'} \|Z\|_{L^{\frac{4}{3}}\left(\frac{2p+2}{\frac{2}{3}}\right)}$$

Using that

$$\left(\frac{4}{3} - \frac{1}{g'}\right)g' = \frac{4}{3}g' - g'\left(1 - \frac{1}{g'}\right)$$

$$= \frac{4}{3}g' - g' + 1$$

$$= \frac{1}{3}g' + 1$$

$$= \frac{(2p+2)}{\binom{2}{3} 3} + 1$$

$$= p + 2,$$

we finally obtain

$$\leq \gamma p \|Y^{p-1} \nabla Y\|_{L^2}^2$$

$$+ \gamma \lambda \|Y\|_{L^{2p+2}}^{2p+2}$$

$$+ C p^{3(p+1)} \lambda^{p+2} \|z\|_{\text{Besov}^{-k}}^{4(p+1)}.$$

Putting the estimates for

and together,

obtain

$$\frac{\lambda}{2} \left| \int z \psi^{2p+1} \right| \leq \gamma p \|\psi^{p-1} \nabla \psi\|_{L^2}^2$$

$$+ \gamma p \lambda \|\psi\|_{L^{2p+2}}^{2p+2}$$

$$+ C (\lambda + p^{3(p+1)} \lambda^{p+2}) \|z\|_{B_{\infty, \infty}}^{4(p+1)}.$$

Arguing similarly, can obtain

$$\frac{\lambda}{2} \left| \int :z^2: \psi^{2p} \right| \leq \gamma p \|\psi^{p-1} \nabla \psi\|_{L^2}^2$$

$$+ \gamma \lambda \|\psi\|_{L^{2p+2}}^{2p+2}$$

$$+ C \left(\lambda + \rho^{3(p+1)} \lambda^{1 + \frac{p+1}{3}} \right) \Big|_{z^2} \Big|_{B_{\rho p}}^{\frac{4(p+1)}{3} - k}$$

$$\frac{\lambda}{2} \left| \int : z^3 : \Upsilon^{2p-1} \right| \leq \gamma \rho \|\Upsilon^{p-1} \nabla \Upsilon\|_{L^2}^2$$

$$+ \gamma \lambda \|\Upsilon\|_{L^{2p+2}}^{2p+2}$$

$$+ C \left(\lambda + \rho^{3(p+1)} \lambda^{1 + \frac{p+1}{5}} \right) \Big|_{z^3} \Big|_{B_{\rho p}}^{\frac{4(p+1)}{5} - k}$$

In the end, obtain

$$\frac{1}{2p} \partial_t \|\Upsilon\|_{L^{2p}}^{2p} + m \|\Upsilon\|_{L^{2p}}^{2p} + \frac{\rho}{2} \|\Upsilon^{p-1} \nabla \Upsilon\|_{L^2}^2$$

$$+ \frac{\lambda}{4} \|\Upsilon\|_{L^{2(p+1)}}^{2(p+1)}$$

$$\leq C(\lambda + \rho^{\frac{4(p+1)}{p}} \lambda^{p+2})$$

$$\left(\|z\|_{B_{\infty, \rho}^{-k}}^{\frac{4(p+1)}{3}} + \|z^2\|_{B_{\infty, \rho}^{-k}}^{\frac{4(p+1)}{3}} \right.$$

$$\left. + \|z^3\|_{B_{\infty, \rho}^{-k}}^{\frac{4(p+1)}{5}} \right) .$$

From this, can obtain moment bounds for Y . In particular, by integrating over $t \in [0, T]$, obtain

$$\frac{1}{2p} \mathbb{E}[\|Y(T)\|_{L^2}^{2p}]$$

$$+ \int_0^T m \mathbb{E}[\|Y(t)\|_{L^2}^{2p}] dt$$

$$+ \int_0^T \mathbb{E}[\|Y(t)^{p-1} \nabla Y(t)\|_{L^2}^2] dt$$

$$+ \frac{\lambda}{4} \int_0^T \mathbb{E} \left[\|Y(t)\|_{L^{2p+2}}^{2p+2} \right] dt$$

$$\leq \frac{1}{2p} \mathbb{E} \left[\|Y(0)\|_{L^{2p}}^{2p} \right]$$

$$+ C_p T (\lambda + p^{3(p+1)} \lambda^{2+p})$$

If Y is stationary in time, then by dividing by T and sending $T \rightarrow \infty$, we obtain

$$m \mathbb{E} \left[\|Y(0)\|_{L^{2p}}^{2p} \right] + \mathbb{E} \left[\|Y(0)^{p-1} \nabla Y(0)\|_{L^2}^2 \right]$$

$$+ \frac{\lambda}{4} \mathbb{E} \left[\|Y(0)\|_{L^{2p+2}}^{2p+2} \right]$$

$$\leq C_p (\lambda + p^{3(p+1)} \lambda^{2+p}).$$