

# Lecture 4: Laminations, general

Note Title

4/25/2011

coefficient systems, and Hyperbolic geometry

① Today, we discuss how to construct cluster algebras from surfaces with general coefficients.

i.e., we want  $(m+n) \times n$  extended matrix  $\tilde{B} = \begin{matrix} n \\ m \\ n \end{matrix} \begin{bmatrix} B \\ C \end{bmatrix}$  where  $C$  is general.

References: Fomin-Thurston, Cluster Algebras from Surfaces II: Lambda Lengths

For this, we use multi-laminations.

Def: We define an integral unbounded measured lamination  $L$  on a marked surface  $(S, M)$  as a finite collection of pairwise non-intersecting curves, each of which has no self-intersections, modulo isotopy (relative to  $M$ ) such that:

Each curve must be either:

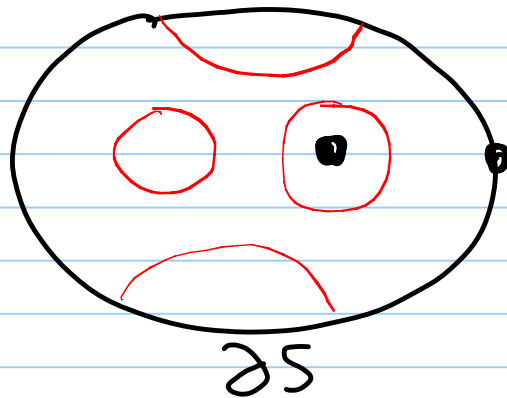
- i) a closed curve
- ii) a curve connecting two unmarked points on  $\partial S - M$
- iii) a curve with one endpoint on  $\partial S - M$  and one endpoint spiraling into a puncture (clockwise OR counter-clockwise), or
- iv) a curve with both endpoints spiraling into a puncture.

② The following curves are disallowed:

- i) a closed curve bounding an unpunctured or once punctured disc,
- ii) a curve with two endpoints on  $\partial S$  which is isotopic to a boundary arc containing 0 or 1 marked point



Disallowed



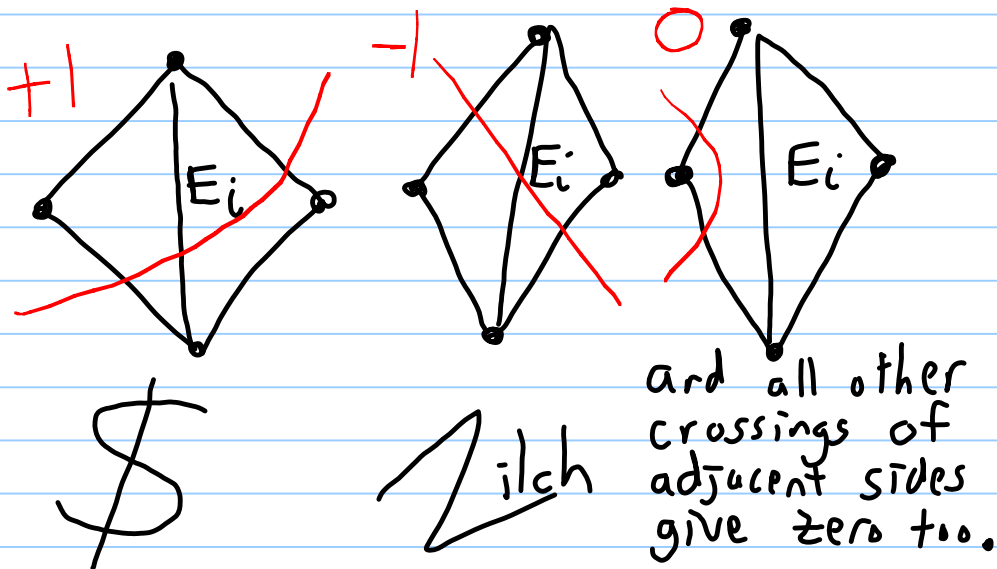
③ Def:

(Tropical)

We assign a Shear Coordinate  $\in \mathbb{Z}$  to each arc  $E_i \in T$  of a triangulation w.r.t. a choice of lamination:

$$b_{E_i}(T, L) \text{ for each } E_i \in T.$$

We look at a quadrilateral inscribing  $E_i$  (in ideal triangulation  $T$ ) for each curve of lamination  $L$  cutting through the quadrilateral, we calculate a contribution to the shear coordinate. Adding them all up gives the appropriate contribution.

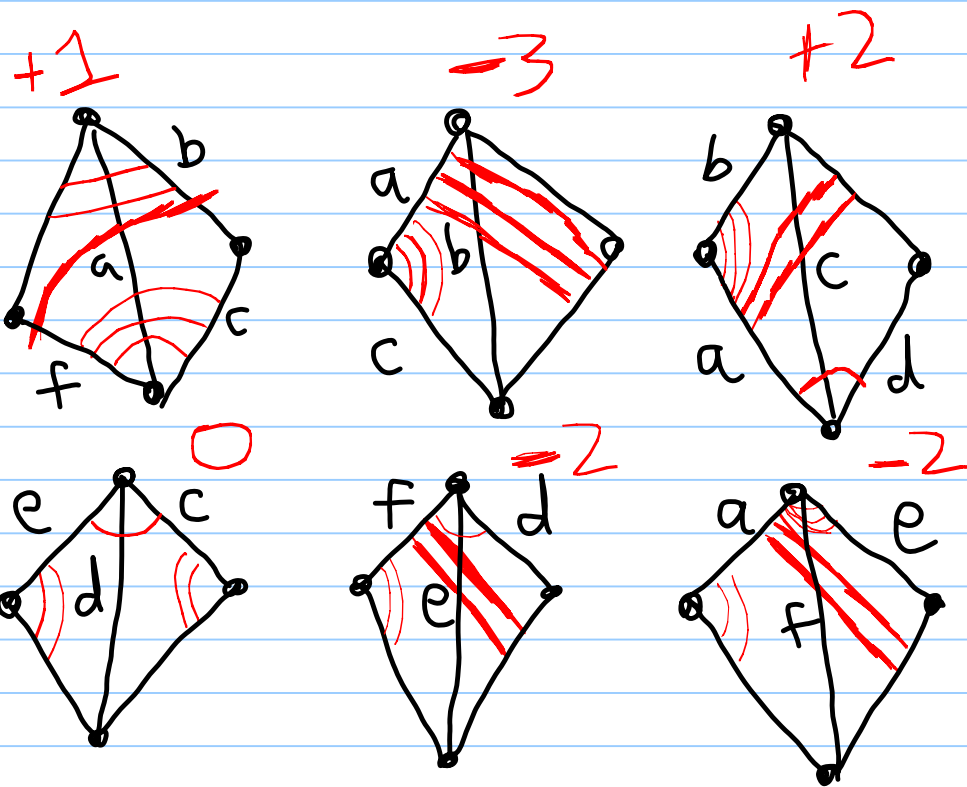
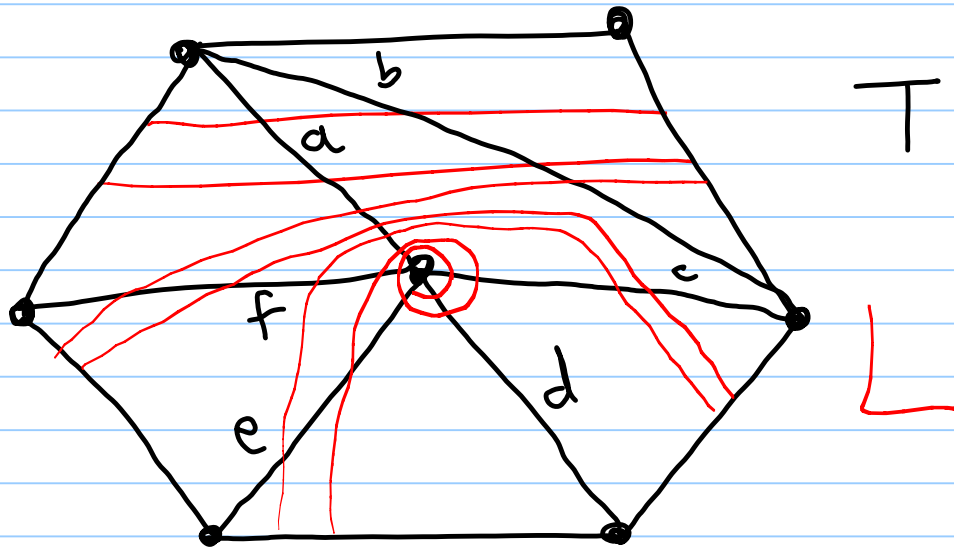


Def: A multi-lamination  $(L_1, \dots, L_m)$  is a collection of  $m$  laminations.

We make  $\tilde{B}$  by  $\begin{bmatrix} B \\ C \end{bmatrix}$ , each row of

$$C = [b_{E_1}(T, L_j), b_{E_2}(T, L_j), \dots, b_{E_n}(T, L_j)]$$

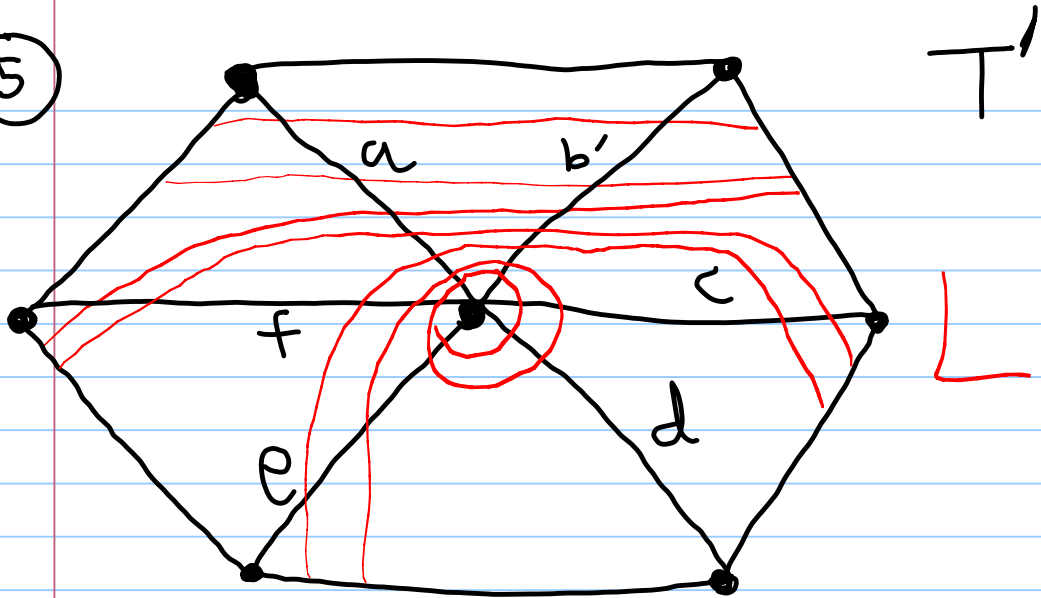
④ Example (Fig 3) of Fomin-Thurston)



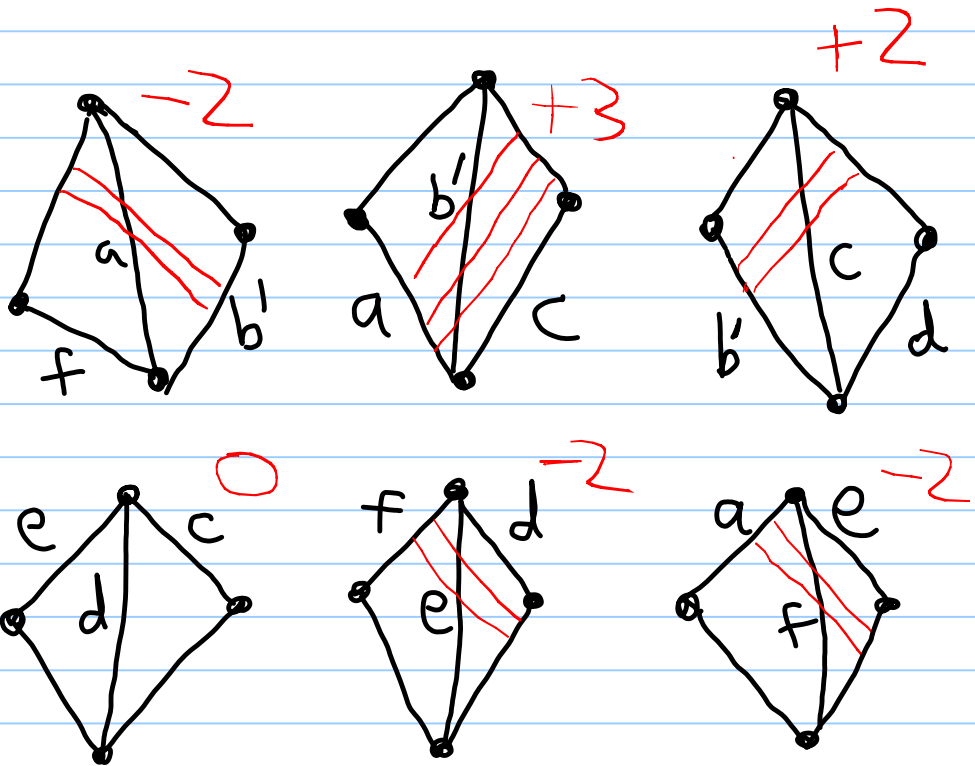
$$b_a(T, L) = 1, b_b(T, L) = -3, \dots$$

Let us now flip  $b \mapsto b'$   
to get triangulation  $T'$  with  
 $L$  fixed:

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$T'$   
 $L$



Notice:  $b_{b'}(T', L) = -b_b(T, L)$

$$b_a(T', L) = b_a(T, L) - \max(-b_b(T, L), 0)$$

$$b_c(T', L) = b_c(T, L) + \max(b_b(T, L), 0)$$

$$b_{\gamma}(T', L) = b_{\gamma}(T, L) \text{ o.w.}$$

$$\textcircled{6} \tilde{B}(T) = \begin{array}{c} a \\ b \\ c \\ d \\ e \\ f \\ L \end{array} \begin{array}{c} a \quad b \quad c \quad d \quad e \quad f \\ \begin{bmatrix} 0 & 1 & -1 & 0 & 0 & 1 \\ -1 & 0 & 1 & 0 & 0 & 0 \\ 1 & -1 & 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 0 & -1 & 0 \\ 0 & 0 & 0 & 1 & 0 & -1 \\ -1 & 0 & 0 & 0 & 1 & 0 \end{bmatrix} \end{array}$$

$\mu_b$

$$\tilde{B}(T') =$$

observe:

Agrees with  
how tropical shear  
coordinates  
change under  
mutating  $T$  to  $T'$ .

$$\begin{array}{c} a \\ b' \\ c \\ d \\ e \\ f \\ L \end{array} \begin{array}{c} a \quad b \quad c \quad d \quad e \quad f \\ \begin{bmatrix} 0 & -1 & 0 & 0 & 0 & 1 \\ 1 & 0 & -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 0 & -1 & 0 \\ 0 & 0 & 0 & 1 & 0 & -1 \\ -1 & 0 & 0 & 0 & 1 & 0 \end{bmatrix} \end{array}$$

Thm (Fomin-Thurston)

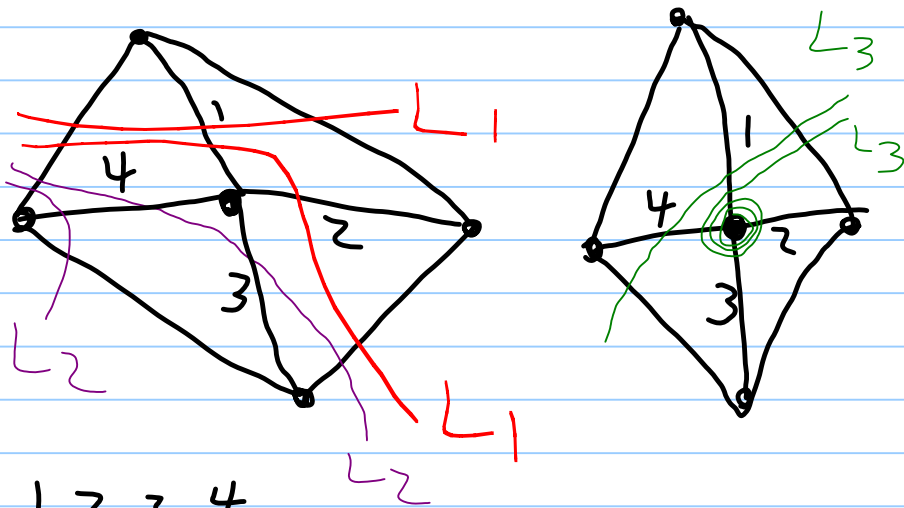
[Based on work by Fock-Goncharov  
and W. Thurston]

If we define  $(m+n) \times n$   $\tilde{B}$ -matrix  
by ideal triangulation  $T = (\tau_1, \dots, \tau_n)$   
and multi-lamination  $(L_1, \dots, L_m)$ ,  
and we mutate  $T \xrightarrow{\mu_k} T'$ ,

then  $\mu_k(\tilde{B}(T, L)) = \tilde{B}(T', L)$ .

Rem: Also an extension to tagged  
triangulations, see Fomin-Thurston.

⑦ Another Example:



	1	2	3	4
1	0	-1	0	1
2	1	0	-1	0
3	0	1	0	-1
4	-1	0	1	0
L1	-1	1	0	0
L2	0	0	-1	1
L3	2	0	0	-1

Recall, the Theorem  
 any row in  $\mathbb{Z}^n$  can  
 be achieved by the  
 choice of some lamination.

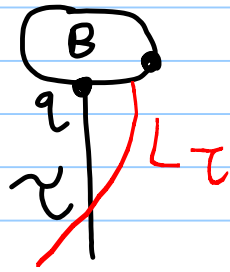
A special case: Principal coeffs

$$\begin{bmatrix} B \\ \hline I_n \end{bmatrix}$$

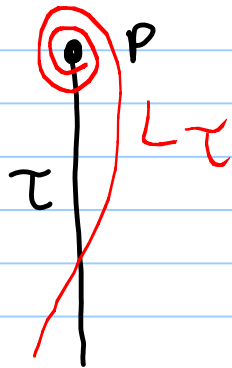
For any tagged arc  $\tau$  in  
 a triangulation, we define  
 the elementary lamination

$L_\tau$  associated to  $\tau$  to be the  
 single curve obtained from  $\tau$   
 by changing  $\tau$ 's endpoints in  
 one of the following 3 ways:

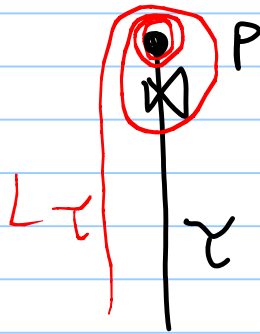
⑧ If  $\gamma$ 's endpoint  $q$  lies on a boundary  $B$ , we move  $q$  slightly "counter-clockwise" around  $B$  so that it misses marked points;



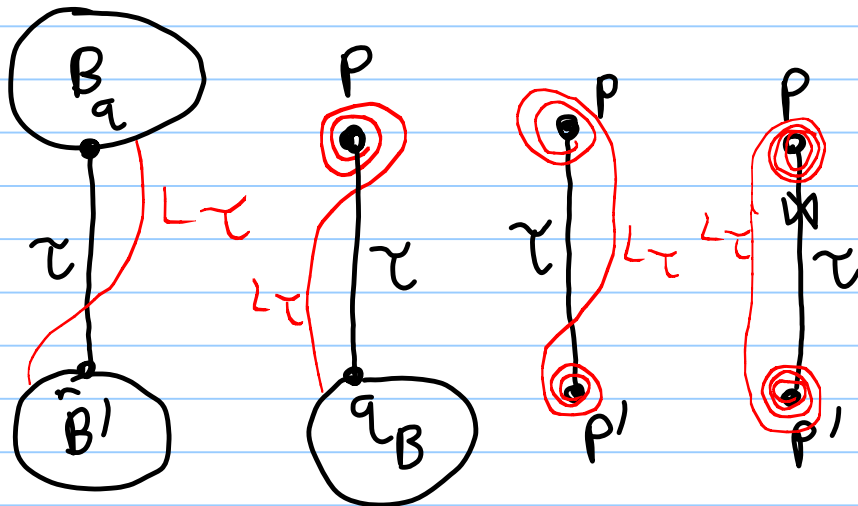
If  $\gamma$ 's endpoint  $p$  lies at a puncture  $\&$   $\gamma$  is tagged plain, we rotate  $L\gamma$  counter-clockwise around  $p$ :



Finally, if  $\gamma$  is notched at  $p_j$ , then we rotate  $L\gamma$  clockwise around  $p$  instead:

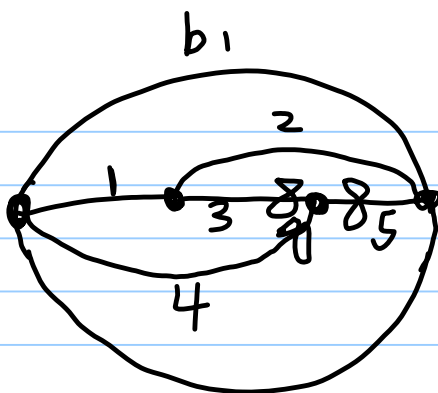


Putting these rules together:

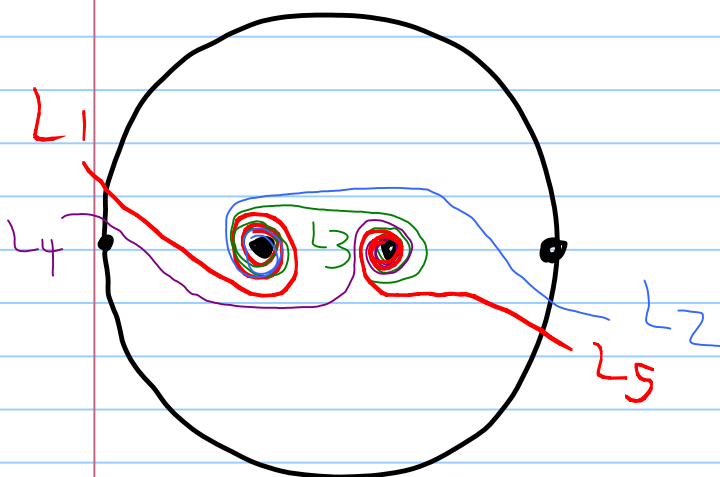




⑨ Example:



gives laminations  $b_2$



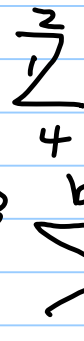
No lamination crosses  $\tau_4$  so crossing quadrilateral in negative formation impossible.

only  $L_1$  crosses  $b_1, \tau_1,$  and  $\tau_3$

so column corresponding to  $\tau_1$  is indeed

$$\begin{matrix} L_1 \\ L_2 \\ L_3 \\ L_4 \\ L_5 \end{matrix} \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$

The other 4 columns give the  $5 \times 5$  identity matrix on the bottom.

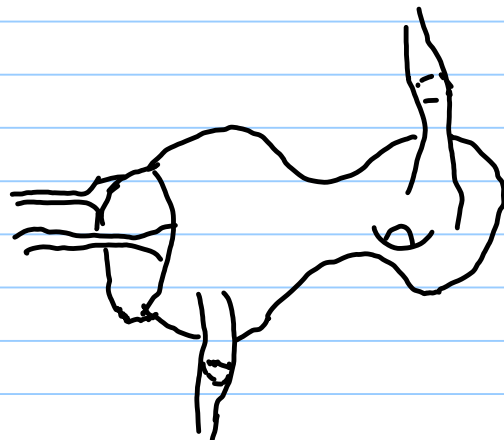
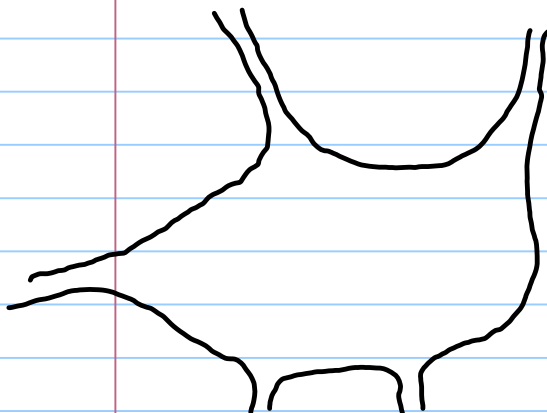
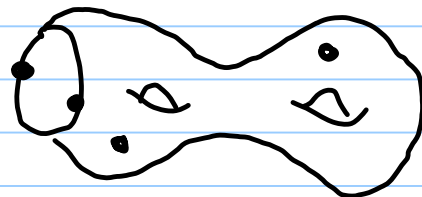
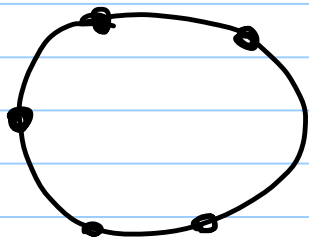


⑩ We now switch gears and talk about Teichmüller theory and how to use hyperbolic geometry to interpret cluster variables and coefficients another way.

Given a marked surface  $(S, M)$   $\mathcal{T}(S, M)$  Teichmüller Space is defined to be the space of metrics on  $(S, M)$  satisfying the following properties:

- hyperbolic (constant curvature  $-1$ )
- has geodesic boundary on boundary of  $S$   
i.e. nbhd of  $p \in \partial S$  looks like nbhd of one side of a geodesic in  $\mathbb{H}^2$
- has cusps at marked points  $M$ .  
(go off to  $\infty$  while area stays bounded.  $(S, M)$  is complete.)

Examples:



①① These metrics are considered up to diffeomorphisms homotopic to identity.

Facts: ①  $\mathcal{T}(S, M)$  is a  $\mathbb{R}$ -manifold of dimension  $6g - 6 + 2p + 3b + c$   
②  $\mathcal{T}(S, M)$  is contractible.

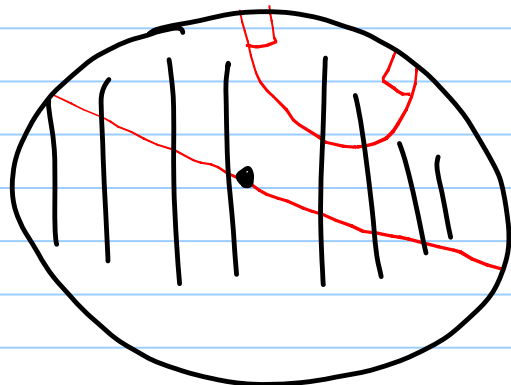
Coordinates on  $\mathcal{T}(S, M)$  by Shear Coordinates (tropical versions defined above)

Thm (W. Thurston) Given  $(S, M)$ , and a triangulation  $T = \{\tau_1, \dots, \tau_n\}$ , then

$L \rightarrow (b_\tau(T, L))_{\tau \in T}$  is a

bijection between the set of (integral unbounded measured) laminations and  $\mathbb{Z}^n$ .

Need hyperbolic geometry to define (non-tropicalized) Shear Coords.



Hyperbolic  $\mathbb{D}$   
Poincaré disk  
model

geodesics

metric =  $ds$  satisfying  $ds^2 = \frac{dx^2 + dy^2}{(1-r^2)^2}$   
 $r = \sqrt{x^2 + y^2}$

Points on Boundary are infinitely far away from the center.

(12) Example: Distance from  $(1,0)$  to  $(0,0)$

$$\text{is } \int_0^1 \frac{dx}{1-x^2} = \operatorname{arctanh}(x) \Big|_0^1 \\ = \frac{1}{2} \log \left( \frac{1+x}{1-x} \right) \Big|_0^1 \rightarrow \infty.$$

Related Model (Upper Halfplane  $\mathbb{H}^2$ )

$$\text{Metric } ds^2 = \frac{dx^2 + dy^2}{y^2}$$

points on real line are infinitely far away.

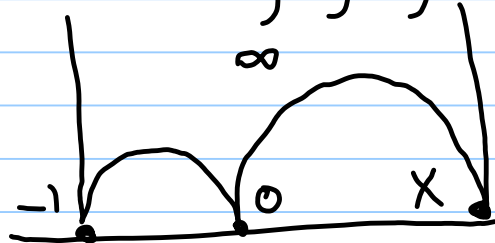
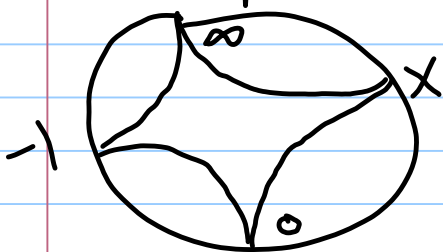


Symmetries by  $PSL_2(\mathbb{R})$  acting by

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} \cdot z = \frac{az+b}{cz+d} \quad \begin{matrix} \text{(fractional)} \\ \text{(linear transf.)} \end{matrix}$$

Action of  $PSL_2(\mathbb{R})$  fixes 3 pts on the boundary and allows freedom to choose others.

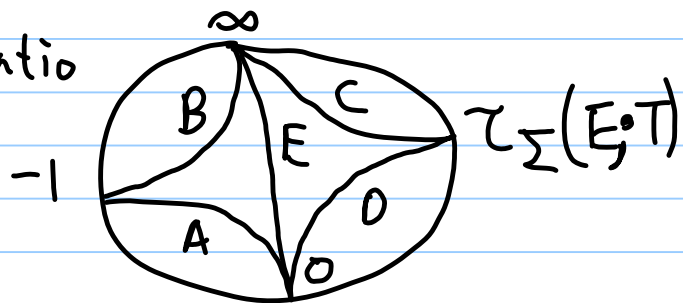
Fig. For quadrilateral on  $\mathbb{D}$  or  $\mathbb{H}^2$  can place points at  $-1, 0, \infty, X$



(13) Remaining corner of a quadrilateral is free  $X \in \mathbb{R}_{>0}$  but is an invariant of the quadrilateral, cross-ratio shear coordinates.

Given a hyperbolic structure  $\Sigma \in \mathcal{J}(S, M)$  and a triangulation  $T = \{E_i\}_{i=1}^n$ , the shear coord.

$\tau_\Sigma(E_j; T)$  of edge  $E \in T$  is the cross-ratio

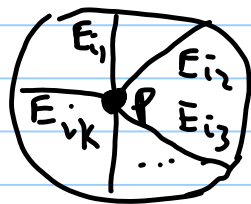


Theorem: The map  $\mathcal{J}(S, M) \rightarrow \mathbb{R}^n$

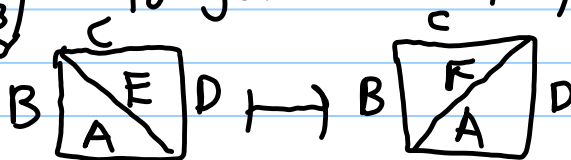
$$\Sigma \mapsto \left\{ \tau_\Sigma(E_{ij}; T) \right\}_{i=1}^n$$

is a homeomorphism onto the subset of  $\mathbb{R}^n$  where for each puncture  $p_j$  and incident arcs  $E_{i_1}, \dots, E_{i_k}$ , we have

$$\prod_{j=1}^k \tau_\Sigma(E_{i_j}; T) = 1.$$



When we Flip quads to get from  $T \rightarrow T'$ ,



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shear coordinates change in a predictable way:

$$\tau_{\Sigma}(F; T') = \tau_{\Sigma}(E; T)^{-1}$$

$$\tau_{\Sigma}(A; T') = \tau_{\Sigma}(A; T) \left(1 + \tau_{\Sigma}(E; T)^{-1}\right)^{-1}$$

$$\tau_{\Sigma}(B; T') = \tau_{\Sigma}(B; T) \left(1 + \tau_{\Sigma}(E; T)\right)$$

$$\tau_{\Sigma}(C; T') = \tau_{\Sigma}(C; T) \left(1 + \tau_{\Sigma}(E; T)^{-1}\right)^{-1}$$

$$\tau_{\Sigma}(D; T') = \tau_{\Sigma}(D; T) \left(1 + \tau_{\Sigma}(E; T)\right)$$

As we will see shortly, these mutations also appear in cluster alg's!

Recall, the original definition

$(\underline{X}, \underline{Y}, B)$   $\underline{X} = \{x_1, \dots, x_n\}$  initial cluster,

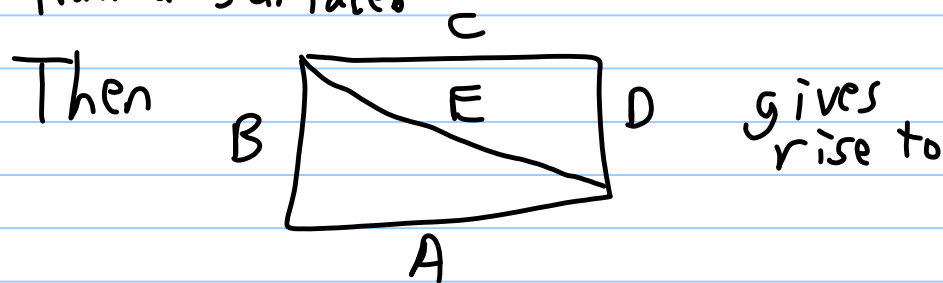
$\underline{Y} = \{y_1, \dots, y_n\}$  initial coeffs,  $B = n \times n$  exchange matrix

$$x'_k x_k = y_k \prod_{b_{ik} > 0} x_i^{b_{ik}} + \prod_{b_{ik} < 0} x_i^{-b_{ik}}$$

$(y_k \oplus 1)$

$$y'_j = \begin{cases} y_k^{-1} & \text{if } j=k, \\ y_j (y_k / (y_k \oplus 1))^{b_{kj}} & \text{if } j \neq k, b_{kj} > 0, \\ y_j (y_k \oplus 1)^{-b_{kj}} & \text{if } j \neq k \& b_{kj} \leq 0. \end{cases}$$

(15) Let us focus on case of tropical semifield and cluster algebra from a surface.



portion of exchange matrix :

$$\begin{array}{c}
 E \ A \ B \ C \ D \\
 \begin{bmatrix}
 E & 0 & 1 & -1 & 1 & -1 \\
 A & -1 & 0 & & & \\
 B & 1 & & 0 & & * \\
 C & -1 & * & 0 & & \\
 D & 1 & & & 0 & 
 \end{bmatrix}
 \end{array}$$

Thus, if we mutate  $\mu_E$   
 $X_{E'} = X_F$  and

$$Y_F = Y_{E'} = Y_E^{-1}$$

$$Y_A' = Y_A \left( \frac{Y_E}{Y_E \oplus 1} \right) = Y_A \frac{Y_E}{Y_E} \left( \frac{1}{1 \oplus Y_E^{-1}} \right)$$

$$Y_B' = Y_B (Y_E \oplus 1)$$

$$Y_C' = Y_C \left( \frac{Y_E}{Y_E \oplus 1} \right) = Y_C \frac{Y_E}{Y_E} \left( \frac{1}{1 \oplus Y_E^{-1}} \right)$$

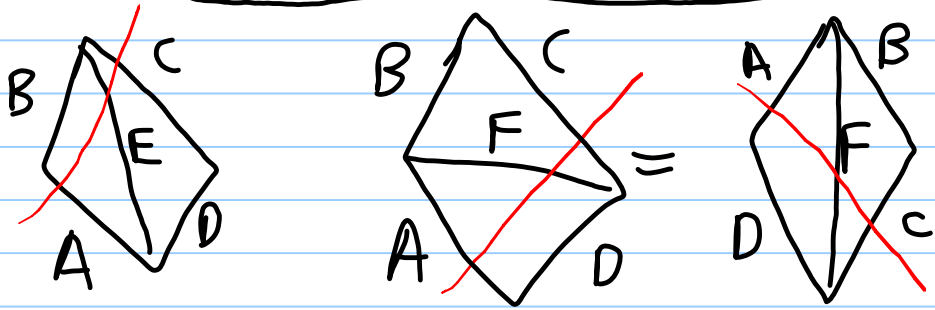
$$Y_D' = Y_D (Y_E \oplus 1)$$

Moral: Coeff Dynamics behave like Tropical Shear coordinate dynamics.

Exactly the Tropical Shear Coordinates or (T, S)'s defined above.

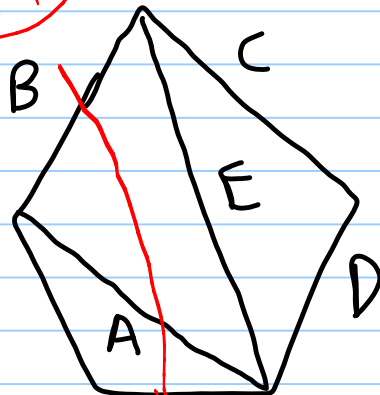
(1b)

More examples comparing how  $b_{\mathcal{L}}(T, L) \rightarrow b_{\mathcal{L}}(T', L)$  changes



sign reversal

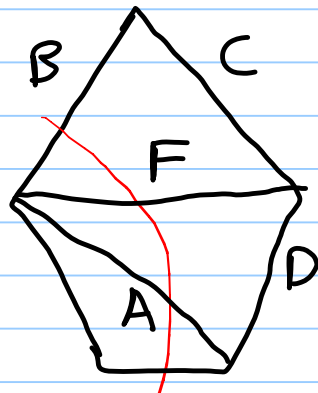
(+1)



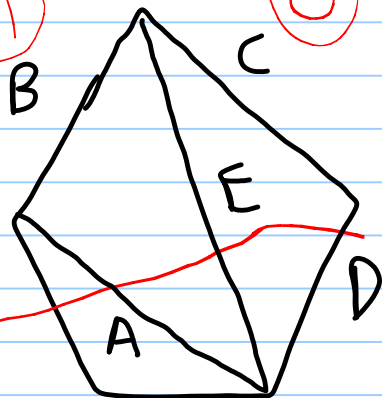
(+1)

still

versus



(-1)

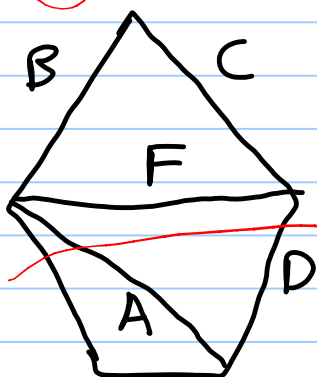


(0)

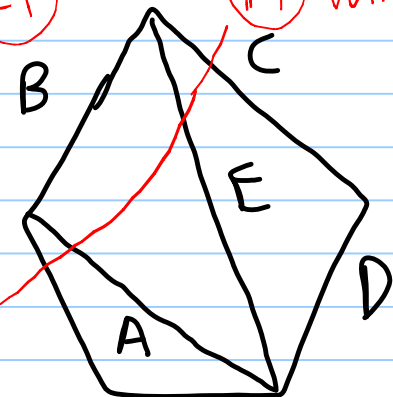
w.r.t. E

(-1) still

versus



(-1)



(+1)

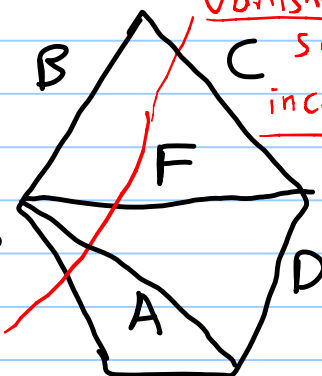
w.r.t. E

(0)

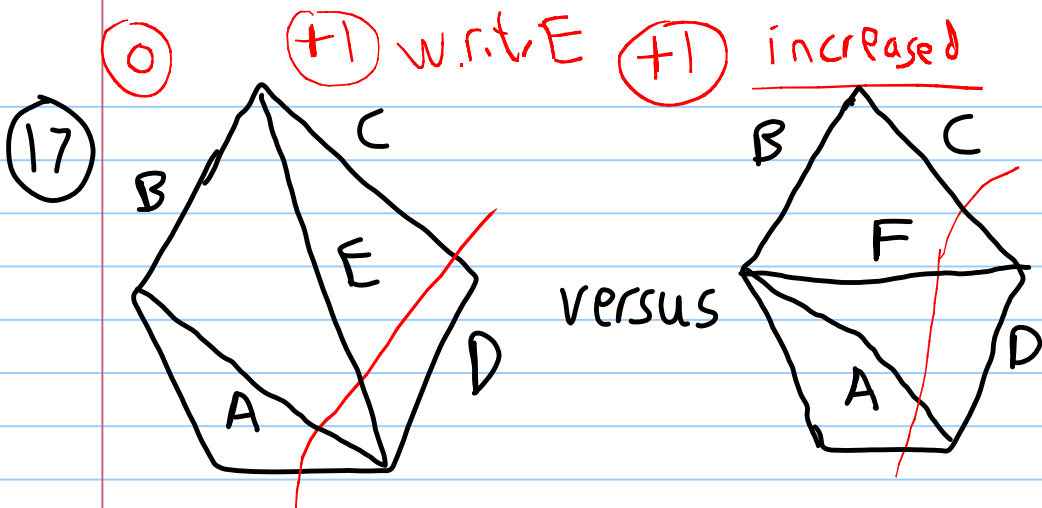
contribution vanishes

so increased

versus







$$b_A(T', L) = b_A(T, L) + \underbrace{b_E(T, L)}_{\text{if } > 0}$$

$$= b_A(T, L) + \max(b_E(T, L), 0).$$

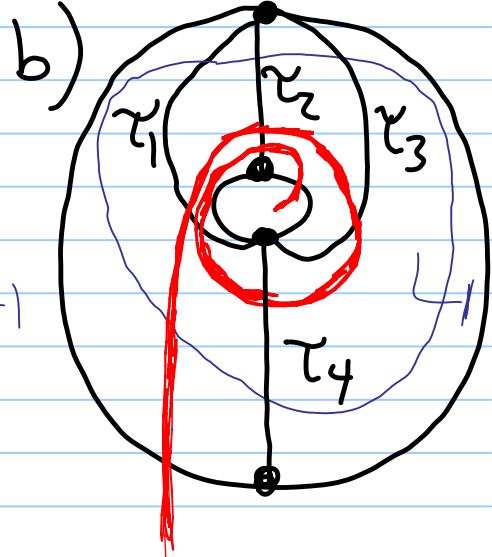
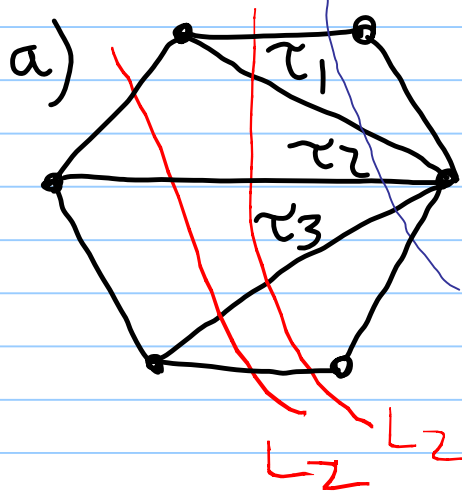
Moral: if we let  $\oplus = \max(-, -)$ ,  
 coefficient dynamics agree  
 with computing trop. shear coordinates with  
 laminations.

Next time: More on  
 Hyperbolic geometry  
 and Lambda Lengths.

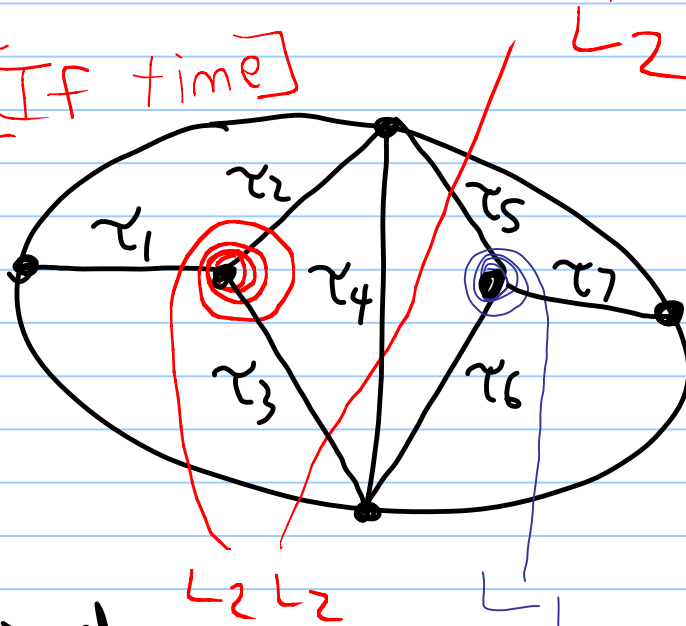
Will also discuss related  $2 \times 2$  matrix  
 formulas (time permitting).

# Lecture 4 Exercises

4-1) Compute  $\tilde{B}$  for some of the following triangulation/multi-lamination pairs:

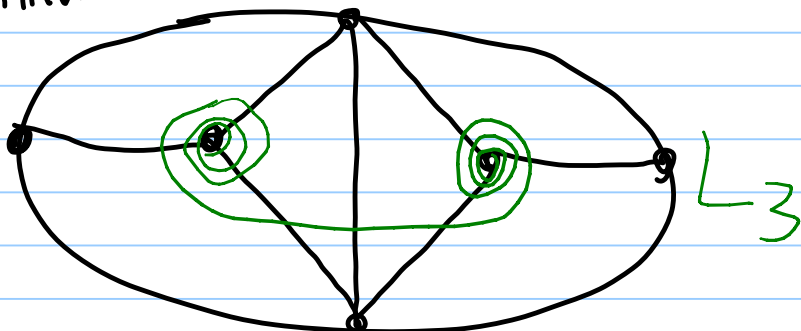


\* [If time]

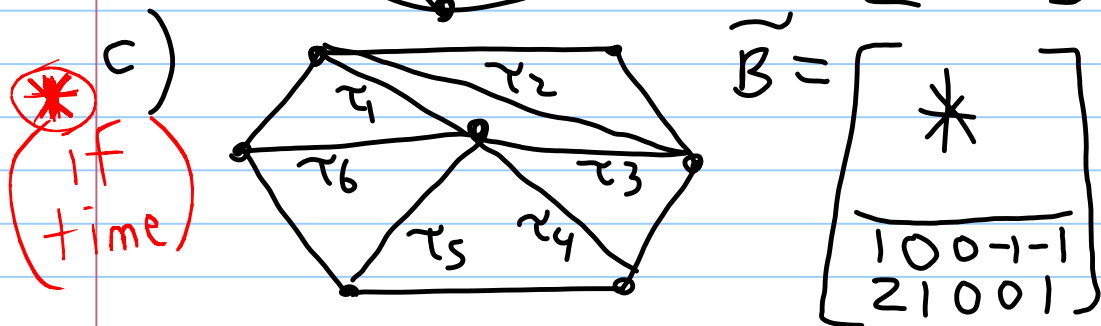
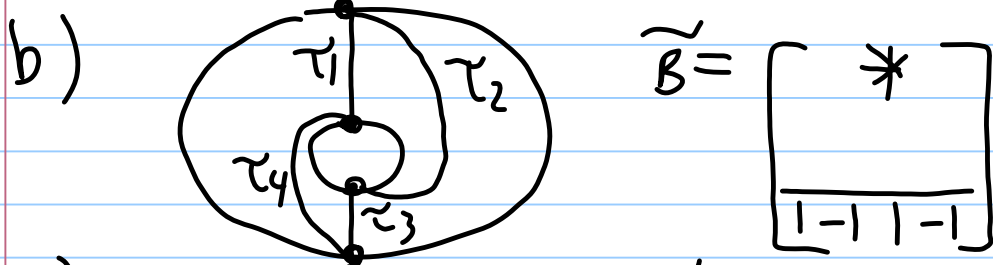
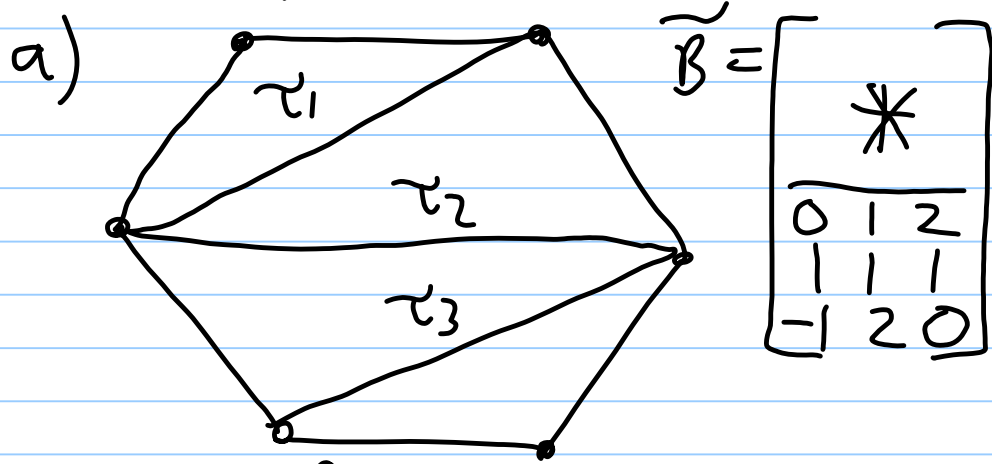


compute bottom of  $\tilde{B}$  only for this example

c) continued



4-2) For the following triangulation  $T$  and  $\tilde{B}$  matrix, compute a corresponding multi-lamination:



4-3) Prove that when we flip triangulation  $T$  to  $T'$ , and let  $b'_\gamma(T, L) := b_\gamma(T', L)$ , ( $\gamma \in T$  or  $T'$ )

we obtain: Exchange relations for  $b_\gamma(T, L)$ 's are tropical version of exchange relations for shear coordinates  $\tau_\gamma(\gamma, T)$ .

*Hint: Look at calculations on pages 14-17 of today's lecture notes*