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Lecture 1

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Plan: Eulerian set-up & Solvability issues
Lagrangian set-up
Regularity in Lagrangian vs. Eulerian setting
Local existence for free-surface Euler

Euler equations

$$\begin{cases} \partial_t u + u \cdot \nabla u + \nabla p = 0 \\ \nabla \cdot u = 0 \end{cases}$$

$$u \cdot \nabla u = u_j \partial_j u$$

$$(j = 1, \dots, d)$$

Initial condition

$$u|_{t=0} = u_0$$

EE incompress. inv. fluid

1) \mathbb{R}^N

2) \mathbb{T}^N

3) $\Omega \subseteq \mathbb{R}^N$ bdd

with $u \cdot n = 0$ on $\partial\Omega$

Also interesting: $\Omega = \mathbb{R}_+^d$ in channel

One first eq. comes from the Newton's limit
mat. derivative of $u = \text{force}$

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Div. free free cond. comes from

$$S_t + \operatorname{div}(u \cdot f) = 0$$

by setting $f \equiv 1$.

Great book: Majda-Bertozzi.

Existence & uniqueness: $W^{m,p}$, $m > \frac{d}{p} + 1$

loc. Well-posedness: $u \in C^{1,\alpha}$ $\alpha \in (0,1)$

Particle Trajectory

Let u be a v.f. Consider particle traj.

$$\eta(\cdot, t) : y \mapsto \eta(y, t) \quad (\Sigma)$$

by for y fixed

$$\begin{cases} \frac{d\eta}{dt}(y, t) = u(\eta(y, t), t) \\ \eta(y, 0) = y \end{cases}$$



y Lagrangian label. Denote

$$J = \det(\nabla \eta)$$

Proposition u is a smooth v.f. Then
 $\partial_t J = \operatorname{div} u(\eta(y, t), t) J$

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$$\operatorname{div} u = 0 : \quad J \equiv 1$$

Proof : $\partial_t (\nabla_y \eta) = \nabla u \nabla \eta$

Proposition (The Transport Theorem)

$\Omega \subset \mathbb{R}^d$ smooth domain, u smooth v.f.,

η is a particle traj map. For any smooth f , we have

$$\frac{d}{dt} \int_{\eta(\Omega, t)} f \, dx = \int_{\eta(\Omega, t)} (\partial_t f + \operatorname{div}(f u)) \, dx$$

$f = \rho$ density : $\partial_t \rho + \operatorname{div}(u \rho) = 0$

$\rho \equiv 1$: $\operatorname{div} u = 0$

Proof $\frac{d}{dt} \int_{\Omega} f(\eta, t) J \, dy$

$$= \int_{\Omega} (\partial_t f + \nabla f \cdot \partial_t \eta + f \underbrace{\partial_t J}_{\operatorname{div} u J}) \, dy$$

Lagrangian form : $v(y, t) = u(\eta(y, t), t)$

Local Existence

(Sketch) The Euler eq. keep L^2 norm of u constant

$$\partial_t u + u \cdot \nabla u + \nabla p = 0 \quad | \cdot u$$

$$\int u_j \underbrace{\partial_j u_k u_k}_{\frac{1}{2} \partial_j (u_k u_k)} = 0 \quad (\text{in } \Omega : u \cdot u = 0)$$

$$\int \partial_k p u_k = 0 \quad (\text{by } \nabla \cdot u = 0)$$

Well-posedness in $C^{1,\alpha}$ $0 < \alpha < 1$: Linchensterm '25, Guntner '27

H^m , $m > \frac{d}{2} + 1$: Ebin-Marsden '70,

Bourgoinien-Brezis '84

Temam '86

A-priori estimates showing existence for $m=3, d=3$

(assume sol'n exists, smooth, w. suff. decay; find bounds based on initial data)

For $|k| \leq 3$:

$$\partial_t \partial^\alpha u_k + \partial^\alpha (u_j \partial u_k) + \partial_k \partial^\alpha p = 0 \quad | \cdot \partial^\alpha u_k, \int_\Sigma$$

$$\frac{1}{2} \frac{d}{dt} \sum_{|k| \leq 3} \int \partial^\alpha u_k \partial^\alpha u_k + \sum_{\substack{|k| \leq 3 \\ 0 < \beta \leq \alpha}} \int \partial^\beta u_j \partial_j \partial^{\alpha-\beta} u_k \partial^\alpha u_k = 0$$

Since $\int u_j \partial_j \partial^\alpha u_n \partial^\alpha u_n = 0$ by $\nabla \cdot u = 0$

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By Hölder and Sobolev

$$\frac{1}{2} \frac{d}{dt} \|u\|_{H^3}^2 \leq C \|u\|_{H^3}^3$$

A-priori estimates for local existence.

$$\|\partial^\alpha (uv) - u \partial^\alpha v\|_{L^2} \lesssim \|u\|_{H^3} \|v\|_{L^\infty} + \|\nabla u\|_{L^\infty} \|v\|_{H^3}$$

α integer: Leibnitz rule Kato - Ponce ineq.
($m > \frac{d}{p} + 1$)

$$T = (I - \Delta)^m, \quad m > \frac{d}{2} + 1$$

Kato - Ponce ineq. does hold in L^p

(Bourgain - Li)

To construct sol'n's, we can use Galerkin (\mathbb{R}^n)

$$u_t + \mathbb{P}_m (u \cdot \nabla u) + \mathbb{P}_m \nabla p = 0$$

where \mathbb{P}_m is proj. on spectral space.

More common:

$$u_t + (u_j)_\varepsilon \cdot \nabla u + \nabla p = 0$$

where $v_\varepsilon = v * \beta_\varepsilon$, $\beta_\varepsilon = \frac{1}{\varepsilon^n} \beta(\frac{\cdot}{\varepsilon})$.

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2D Yudovich data : Global & uniq.

$$u_0 = \text{curl } u_0 \in L^1 \cap L^\infty$$

M. Vishik : fails in $L^1 \cap L^p$, $p < \infty$

3D global existence : unknown

$$\text{BKM: } \text{if } T_{\max} < \infty : \int_0^{T_{\max}} \|u\|_{L^\infty} = \infty$$

Elpridi : Global existence may fail in $C^{1,\alpha}$

NSE: C^2 data : global existence of Leray weak sol's

EP: this is unknown

\exists measure-valued sol's

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Vorticity Transport Formula

Introduce

$$\omega = \text{curl } u$$

(motivation: $\text{curl } \nabla p = 0$)

In 2D:

$$\omega = \partial_1 u_2 - \partial_2 u_1$$

In 3D:

$$\omega = \begin{bmatrix} \partial_2 u_3 - \partial_3 u_2 \\ \partial_3 u_1 - \partial_1 u_3 \\ \partial_1 u_2 - \partial_2 u_1 \end{bmatrix}$$

$$(\Omega = \frac{1}{2}(\nabla u - (\nabla u)^T))$$

To get an equation for ω in 3D, use

$$u \cdot \nabla u = \frac{1}{2} \nabla(|u|^2) - u \times \text{curl } u$$

We get in 2D:

$$\omega_t + u \cdot \nabla \omega = 0$$

In 3D

$$\omega_t + u \cdot \nabla \omega = \omega \cdot \nabla u.$$

To compute ω :

$$\begin{aligned} \text{div } u &= 0 \\ \text{curl } u &= \omega \\ (u \cdot n &= 0) \end{aligned}$$

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To solve for u : curl of 2nd eq.

$$-\Delta u = \text{curl } \omega$$

Invert: $u = K(\omega)$

K smoothing of degree 1.

We write GF :

$$\begin{aligned} \omega_t + K(\omega) \cdot \nabla \omega &= 0 && \text{in } 2D \\ &= \omega \cdot \nabla u && \text{in } 3D \end{aligned}$$

$$\omega \cdot \nabla u = \omega_j \partial_j u = \nabla u \cdot \omega$$

$$(\nabla u)_{i,j} = u_{i,j} = \partial_j u_i$$

Theorem (3D vorticity transport formula)

$$\omega(\eta(y,t), t) = \underbrace{\nabla_y \eta(y,t)}_{\text{matrix}} \omega_0(y)$$

$$\xi(y,t) = \omega(\eta(y,t), t) \text{ --- Lagrangian vorticity}$$

Proof $(D_t = \partial_t + u \cdot \nabla)$

For

$$D_t \omega = \nabla u \cdot \omega$$

\therefore

$$\frac{\partial}{\partial t} \xi(y,t) = \nabla u |_{(\eta(y,t), t)} \cdot \omega(\eta(y,t), t)$$

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On the other hand, $\partial_t \eta = u(\eta, t)$ implies

\forall fixed vector h :

$$\frac{\partial}{\partial t} (\nabla_{\eta} h) = \nabla_u \nabla_{\eta} h$$

If we choose $h = \omega_0$:

$$\nabla_{\eta} h \Big|_{t=0} = \nabla_{\eta} \Big|_{t=0} \omega_0 = \omega_0$$

By the uniqueness for ODEs:

$$\xi(\eta, t) = \nabla_{\eta} h = \nabla_{\eta} \omega_0. \quad \square$$