

Lecture 8
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Normal forms (enhanced lifespan of solutions)

Eg 1 If one has an eq that quadratic nonlinearities

$$\frac{d}{dt} E(u) \lesssim \|u\| E(u) \rightarrow \text{quadratic energy estimate}$$

control norm. $\dot{x} = x^2$

$$\|u(0)\| < \varepsilon \quad \downarrow \text{ by Gronwall. } T_\varepsilon \approx \frac{1}{\varepsilon} \quad \text{quadratic lifespan}$$

Eg 2 If one performs energy estimates for an equation that only cubic nonlinearities, and we look at small initial data,

$$\frac{d}{dt} E(u) \lesssim \|u\|^2 E(u) \rightarrow \text{cubic energy estimate}$$

$$\|u(0)\| < \varepsilon \rightarrow \text{ by Gronwall } T_\varepsilon \approx \frac{1}{\varepsilon^2} \quad \text{cubic lifespan.}$$

$\dot{x} = x^3$

Obs: If one gets from eg 1 \rightarrow to eg 2 then you improve the lifespan of your solutions.

obs: this idea applies to any type of nonlinearity
cubic nonlinearity \rightarrow quartic or higher
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Example 1.

$$\begin{cases} i u_t - a u_{xx} + b u^2 = 0 & x \in \mathbb{R} \\ u(a, x) = u_0 \end{cases}$$

Use Hopf-Cole transformation. (Evans, chp 4)

$w := \phi(u)$, $\phi: \mathbb{R} \rightarrow \mathbb{R}$ a smooth function.

$$\begin{array}{l|l}
 w_t = \phi' u_t & i w_t = \phi' u_t \\
 w_x = \phi' u_x & = \phi' [a u_{xx} - b u_x^2] \\
 w_{xx} = \phi'' u_x^2 + \phi' u_{xx} & = \phi' a u_{xx} - b \phi' u_x^2 \\
 & = a [w_{xx} - \phi'' u_x^2] - b \phi' u_x^2 \\
 & = a w_{xx} - a \phi'' u_x^2 - b \phi' u_x^2 \\
 \hline
 i w_t & = a w_{xx} - u_x^2 [a \phi'' + b \phi']
 \end{array}$$

$a \phi'' + b \phi' = 0$

$a k^2 + b k = 0 \quad r_1 = 0, r_2 = -\frac{b}{a}$

$w := \phi(u) = c_1 e^{u \cdot 0} + c_2 e^{u \cdot (-\frac{b}{a})} \quad c_1 = 1, c_2 = -1 \rightarrow \text{initial cond.}$

$w := e^{-\frac{b}{a} u} - 1$

= Hopf-Cole transformation.

$\begin{cases} i w_t - a w_{xx} = 0 \\ w_0 = e^{-\frac{b u_0}{a}} - 1. \end{cases}$

→ solve it explicitly.

$w(t, x) = \frac{1}{(\pi a t)^{n/2}} \int_{\mathbb{R}^n} e^{-\frac{|x-y|^2}{4 a t} - \frac{b u_0(y)}{a}} dy$

bounded.

→ bounded transformation from $H^{\frac{n}{2} + \epsilon} \rightarrow H^{\frac{n}{2} + \epsilon}$.

orig eq was quadratic $\xrightarrow{\text{bounded transf}}$ to a linear eq which has global sols



Question: Do I get to always find such transformations that remove nonlinearities?

Answer: No!

Conclusion: Before computing a normal form to remove a quadratic nonlinearity for example, one needs to solve a particular system.

Three wave resonances system

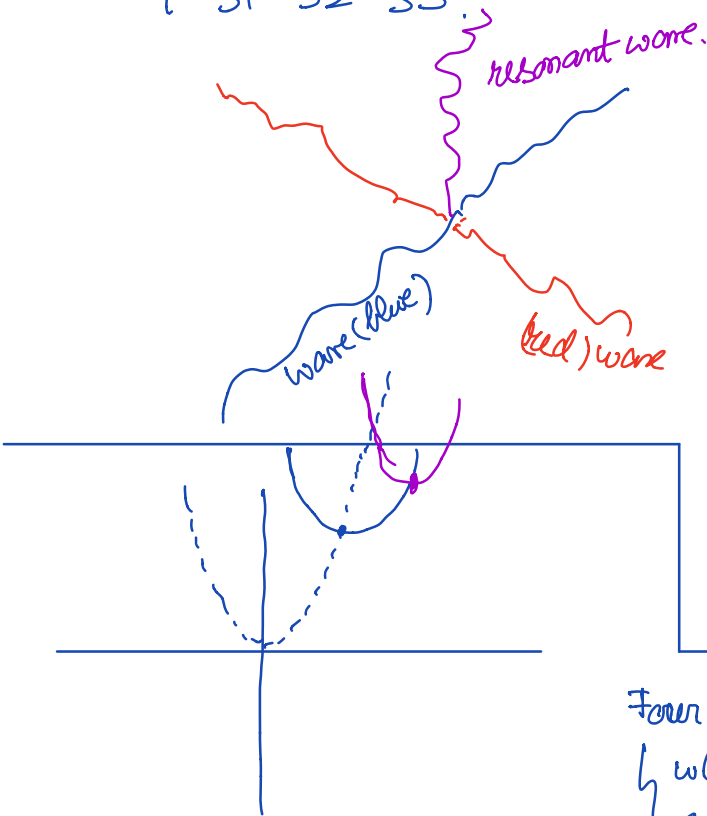
$\omega(\xi) \rightarrow$ dispersion relation of the linear part of a given PDE

$$\begin{cases} \omega(\xi_1) + \omega(\xi_2) = \omega(\xi_3) \\ \xi_1 + \xi_2 = \xi_3 \end{cases}$$

• solutions means you have resonances

• no solution, no resonances (quadratic resonances) \Rightarrow removable quadratic nonlinearity.

\rightarrow if the of the quadratic nonlinearity vanishes on the set of resonances (pairs of resonant ξ) then the quadratic nonlinearity has a null form structure



Four wave resonances.

$$\begin{cases} \omega(\xi_1) + \omega(\xi_2) \pm \omega(\xi_3) = \omega(\xi_4) \\ \xi_1 + \xi_2 \pm \xi_3 = \xi_4 \end{cases}$$

- History:
- ① Normal form theory (NFT) Birkhoff in early 1920s ODE.
 - ② Shatake 1983 Klein-Gordon eq \rightarrow GKP.
 - ③ Simon 1983 Klein-Gordon eq \rightarrow asymptotic completeness.
 - ④ There are approaches nowadays. (see it next lectures)

Example 2: $i u_t + \sqrt{\Delta^2 + 1} u = u^2$ (half-wave for the K-G eq)

$$\omega(\xi) = \sqrt{\xi^2 + 1} \quad x \in \mathbb{R}$$

Three wave resonances system $\rightarrow \omega(\xi) + \omega(\eta) = \omega(\xi + \eta)$

$$\sqrt{\xi^2 + 1} + \sqrt{\eta^2 + 1} = \sqrt{(\xi + \eta)^2 + 1}$$

$$4(\xi^2 + \eta^2 - \xi\eta) = -3$$

needs to be odd

$\tilde{u} := u + B(u, u)$ $B =$ bilinear symmetric form in u .

$$u = \tilde{u} - B(u, u)$$

$$i u_t = i \tilde{u}_t - \cancel{B(\tilde{u}_t, u)} - i B(u, \tilde{u}_t)$$

$$= i \tilde{u}_t - B(u^2 - \sqrt{\Delta^2 + 1} u, u) - B(u, u^2 - \sqrt{\Delta^2 + 1} u)$$

$$= i \tilde{u}_t - B(u^2, u) + B(\sqrt{\Delta^2 + 1} u, u)$$

$$- B(u, u^2) + B(u, \sqrt{\Delta^2 + 1} u)$$

$$\sqrt{\Delta^2 + 1} u = \sqrt{\Delta^2 + 1} \tilde{u} - \sqrt{\Delta^2 + 1} B(u, u)$$

$i \tilde{u}_t \neq \sqrt{\Delta^2 + 1} \tilde{u}$
 $= 2B(u^2, u)$

$$i u_t + \sqrt{\Delta^2 + 1} u = i \tilde{u}_t + \sqrt{\Delta^2 + 1} \tilde{u} +$$

$$+ B(\sqrt{\Delta^2 + 1} u, u) + B(u, \sqrt{\Delta^2 + 1} u) - \sqrt{\Delta^2 + 1} B(u, u)$$

$$+ 2B(u^2, u)$$

I need to find $B(\cdot, \cdot)$ s.t.

$$u^2 = B(\sqrt{\Delta^2 + 1} u, u) + B(u, \sqrt{\Delta^2 + 1} u) - \sqrt{\Delta^2 + 1} B(u, u)$$

$$1 = \hat{B}(\xi, \eta) \sqrt{\xi^2+1} + \hat{B}(\xi, \eta) \sqrt{\eta^2+1} - \sqrt{(\xi+\eta)^2+1} \hat{B}(\xi, \eta)$$

$$\hat{B}(\xi, \eta) = \frac{1}{\sqrt{\xi^2+1} + \sqrt{\eta^2+1} - \sqrt{(\xi+\eta)^2+1}} = \frac{1}{\omega(\xi) + \omega(\eta) - \omega(\xi+\eta)}$$

$$B(u, u) \approx 1 + \min\{|\xi|, |\eta|\}$$

$$B(u, u) = \underbrace{\omega \cdot Du}_{\text{low freq}} \rightarrow B: H^s \rightarrow H^s$$

$$H^s, s \geq \frac{3}{2}$$

$$\partial_t \tilde{u} + \sqrt{S^2+1} \tilde{u} = 2B(u^2, u)$$

$$\|B(u^2, u)\|_{H^s} \lesssim \|u^2\|_{H^s} \|u\|_{H^s} \lesssim \|\tilde{u}\|_{H^s}^3 \checkmark$$

Ans: Both examples worked, semilinear eq!

For quasi-linear problems things are harder:

Burgers Hilbert: $u_t + (u u_x) = \mu u \rightarrow \frac{1}{2} u u_x + \frac{1}{2} u_x u$
 $\tilde{u} = u + B(u, u) \quad \frac{1}{2} i(\eta) + \frac{1}{2} i(\xi)$

$$\hat{B}(\xi, \eta) = \frac{1}{2} \frac{\xi + \eta}{\text{sgn } \xi + \text{sgn } \eta - \text{sgn } (\xi + \eta)}$$

$$\omega(\xi) = -\text{sgn } \xi$$

$$\frac{1+1=1}{-1-1=-1}$$

$$\frac{-1-1=-1}{-1+1=+1}$$

$$\frac{-1+1=+1}{-1-1=-1}$$

$$B(u, u) \uparrow$$

$$\text{I } \xi, \eta > 0$$

$$\text{II } \xi, \eta < 0$$

$$\text{III } \xi > 0, \eta < 0, |\xi| > |\eta|$$

$$\vdots$$

$\text{sgn } (\xi) \text{sgn } (\eta) \text{sgn } (\xi+\eta)$
 We need to see the spatial formulation of B is (when possible) so the idea is to get rid of the denominator. Obviously, not in all cases this is possible. Here it is.

$$B(\xi, \eta) = \frac{1}{2} (\xi + \eta) \operatorname{sgn}(\xi) \operatorname{sgn}(\eta) \operatorname{sgn}(\xi + \eta)$$

$$B(u, u) = \frac{1}{2} H [\underbrace{Hu}_{\text{low freq}} \cdot \underbrace{Hu_x}_{\text{high freq}}]$$

issues !!

1. HFT is unbounded
2. the transf ϵ_2 does not have energy estimate

$$\tilde{u}_t = H \tilde{u} + \text{Cubic}(u, u, u)$$

$$\frac{d}{dt} \|\partial_x^k \tilde{u}\|_{L^2}^2 \lesssim \|u_x\|_{L^\infty}^2 \cdot \|u\|_{H^{k+1}}^2$$