

Summer School Lecture 7

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Q: Given a quasilinear PDE,
what can we say about the lifespan
of the solutions?

Q1. Suppose the initial data is
small,

$$\|u_0\| \leq \varepsilon$$

$$T(u_0) \geq T_\varepsilon$$

(a) $u_0 \in H^s$, small \rightarrow Michael

(b) $u_0 \in H^s$, small + localized \rightarrow David

The vector field method.

Energy estimates :

$$E^S(u) \Rightarrow H^S \text{ energy}$$

$$\frac{d}{dt} E^S(u) \in \mathcal{B} E^S(u)$$

↓
 L^∞ type control norm.
[may decay]

- prior discussion :

L^∞ norms have dispersive decay in linear flows

→ via fundamental solutions

→ stationary phase method

Nonlinear flow: No access to such methods

→ we control energies $\Rightarrow H^S$ norms

→ useless for pointwise decay

→ can we get access to more energies which can give pointwise decay?

Vector field method

A brief history:

- nonlinear wave equation

60's → Marwitz

actual vector fields
from wave eq. symmetries

80's → John

90's → Klainerman, Christodoulou

90's → Alinhac

90's → Hörmander, Sogge books.

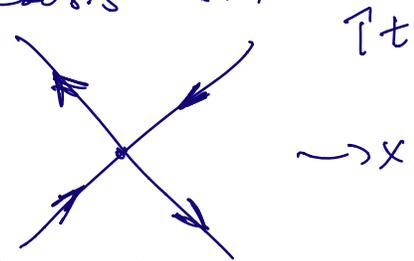
Wave equation

$$\square u = 0, \quad \square = \partial_t^2 - \Delta_x$$

Lorentz group

→ spatial rotations

→ Lorentz boosts (+)



Generators for the Lorentz group:

$$\begin{aligned} \mathcal{L}_{ij} &= x_i \partial_j - x_j \partial_i \\ \mathcal{L}_{0j} &= x_i \partial_t + t \partial_{x_i} \end{aligned} \quad \Bigg| \quad S = t \partial_t + x \partial_x$$

$$\{\Omega_{ij}, \square\} = 0$$

$$\text{If } \square u = 0 \Rightarrow \square \Omega u = 0$$

$E(u) = \text{conserved} \rightarrow E(\Omega u) \text{ conserved.}$

$$Q: \quad \|\nabla u\|_{L^\infty}^2 \leq t^{-\frac{d-1}{2}} \sum_{\alpha} E(\Omega^\alpha u)$$

Klainerman - Sobolev inequalities

\rightarrow They are Sobolev embeddings
in a good frame of reference.

\rightarrow in our work $\left\{ \begin{array}{l} \text{Metcalfe - Tohaneanu - T.} \\ \text{Ifrim - Stingo} \end{array} \right.$

What about nonlinear waves?

Suppose Ω generates a symmetry for nonlinear wave eqn. for u .

$\Omega u = s u$. for linearized eqn.

$\Omega^\alpha u = \text{linearized eqn. with source terms}$

Linear model equations

Example 1: The Schrödinger equation.

$$(i \partial_t - \Delta) u = 0, \quad u(0) = u_0$$

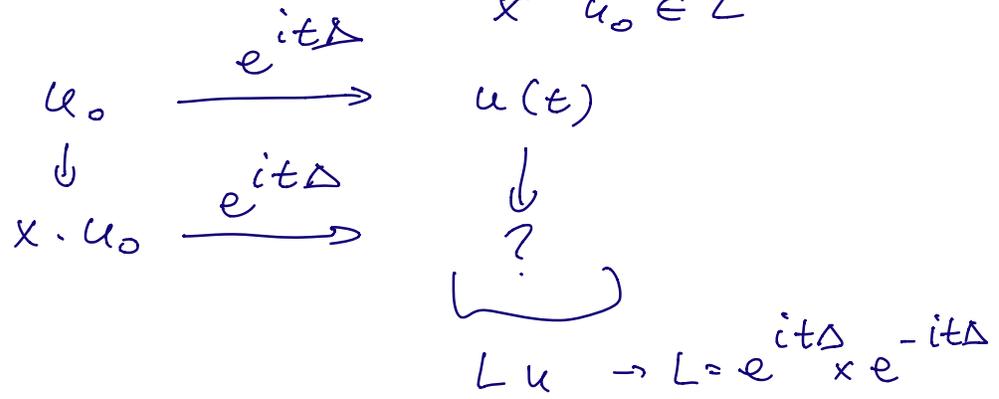
Dispersive decay:

$$\|u(t)\|_{L^\infty} \leq t^{-\frac{n}{2}} \|u(0)\|_{L^1}$$

Challenge: Prove $t^{-\frac{n}{2}}$ decay if $u(0) \in L^2$ is localized.

Vector fields \Rightarrow Galilean symmetry

$$u_0 \in L^2 \text{ localized: } \begin{aligned} u_0 &\in L^2 \\ x \cdot u_0 &\in L^2 \\ \vdots \\ x^\alpha u_0 &\in L^2 \end{aligned}$$



Egorov theorem in microlocal analysis

$$L_j = x_j - 2it \partial_j \Rightarrow [L_j, i\partial_t - \Delta] = 0$$

$u_0 \in L^2$, localized \Rightarrow

$$\|L^\alpha u(t)\|_{L^2} = \|x^\alpha u_0\|_{L^2} < \infty.$$

Claim:

$$\|u(t)\|_{L^\infty} \leq t^{-\frac{n}{2}} \cdot \sum_{\alpha \in \mathbb{N}^n} \|L^\alpha u\|_{L^2}$$

Interpret as a Sobolev embedding:

$$v = e^{i\frac{x^2}{4t}} u$$

$$e^{i\frac{x^2}{4t}} Lu = 2t \cdot \partial v$$

$$\|v\|_{L^\infty} \leq t^{-\frac{n}{2}} \sum_{|\alpha| \leq N} t^{|\alpha|} \|\partial^\alpha v\|_{L^2}$$

$$\|v\|_{L^\infty} \leq \|v\|_{\dot{H}^{\frac{n}{2}}}$$

False but almost true!

$$\|v\|_{L^\infty}^2 \leq \|v\|_{\dot{H}^{\frac{n}{2}-1}} \cdot \|v\|_{\dot{H}^{\frac{n}{2}+1}}$$

True!

Interpolation inequality

Example 2 KdV equation, linear

$$u_t + u_{xxx} = 0$$

Vector field: $L = x - 3t\partial_x^2$ // better

Scaling symmetry $\tilde{L} = \partial_x - 3t\partial_x^3$

$$\|u\|_{L^2} + \|Lu\|_{L^2} \geq L^\infty \text{ bounds for } u.$$

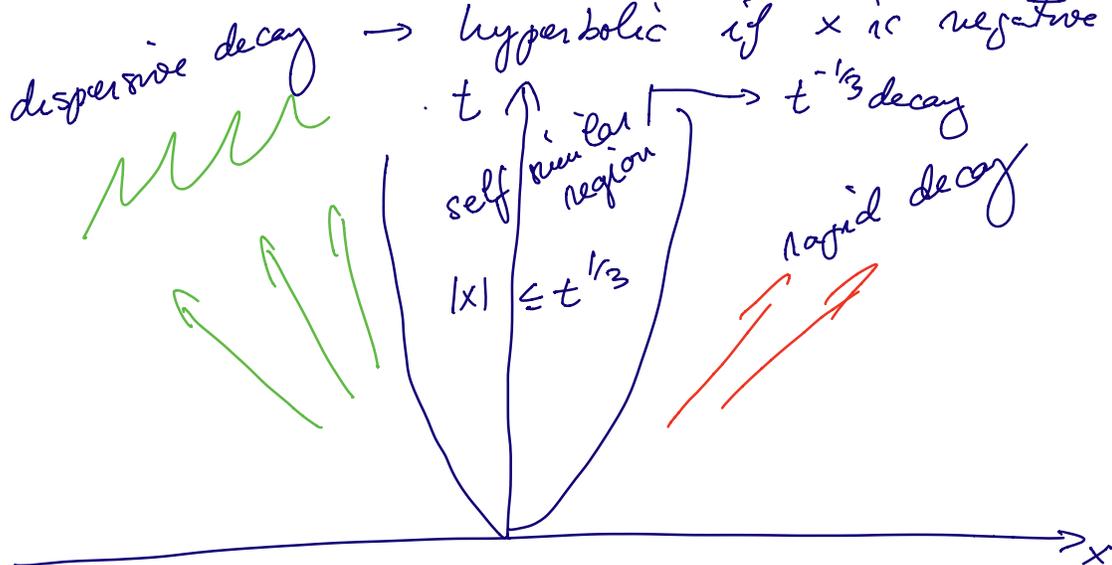
or any other Sobolev norm.

Symbol for L :

$$\ell(x, \xi) = x + 3t\xi^2$$

→ elliptic if x is positive

→ hyperbolic if x is negative



Example 3

$$(i\partial_t + A(D))u = 0$$

$$Lu = (x + t a_{\frac{1}{2}}(D))u$$

$$[L, i\partial_t + A(D)] = 0$$

$$\|u\|_{L^2} + \|Lu\|_{L^2} \stackrel{?}{\geq} c(v)\|u\|_{L^\infty}$$

Nonlinear problems

Example:

$$(i\partial_t - \Delta)u = u|u|^4 \quad \text{in 1-d.}$$

$$u_0 \in H^1, \quad xu_0 \in L^2$$

$$\|u_0\|_{H^1} + \|xu_0\|_{L^2} \leq \varepsilon$$

Theorem. If $\varepsilon \ll 1$ then solutions are global, and decay like $t^{-1/2}$.

Proof Energy conservation:

$$\|u\|_{L^2} + \|\partial_x u\|_{L^2} \leq \varepsilon$$

$$Lu = (x - 2it\partial_x)u$$

ii.

\checkmark
 $L u$ solves the linearized equation.

$$(i \partial_t - \Delta) v \approx v \cdot u^4$$

$$\frac{d}{dt} \|v\|_{L^2}^2 \leq \|v\|_{L^2}^2 \cdot \|u\|_{L^\infty}^4$$

$$\|L u(t)\|_{L^2} \leq \|L u(0)\|_{L^2} \cdot e^{\int_0^t \|u\|_{L^\infty}^4 dx}$$

Expect: $\|u\|_{L^\infty} \leq \frac{\varepsilon}{\sqrt{t}}$

If this is true, then we get:

$$\textcircled{2} \quad \|L u(t)\|_{L^2} \leq \varepsilon$$

K-S bound:

$$\textcircled{3} \quad \|u\|_{L^\infty} \leq \frac{1}{\sqrt{t}} \|u\|_{L^2} \cdot \|L u\|_{L^2} \\ \leq \frac{\varepsilon}{\sqrt{t}}$$

Make bootstrap assumption:

$$\|u\|_{L^\infty} \leq \frac{1}{\sqrt{t}} \left[\frac{\varepsilon^{1/2}}{\sqrt{t}} \right]$$

This implies $\textcircled{2}$ and then $\textcircled{3}$.

Final comments:

① Sometimes the decay might not be enough to get global solutions.

- Ifim-T. Benjamin Ono

$$T_\varepsilon = e^{-\frac{c}{\varepsilon}} \text{ (almost global)}$$

- Ifim-Kob-T. KdV equation

$$T_\varepsilon = \varepsilon^{-3} \text{ (scattering time)}$$

Related to solitons / soliton resolution conjecture.

② Classical KS ineq. $\Rightarrow L$ is a linear operator.

$$KS = \hat{\text{linear}}$$

BO + KdV:

$$L \rightarrow L^{NL}$$

linear

nonlinear

$$\|u\|_{L^2} + \|L^{NL}u\|_{L^2} \geq \sqrt{c(V)} \|u\|_{L^\infty}$$

Nonlinear K-S inequalities