

Lecture 3

Aug 4, 2020

Last time:

Theorem (CKV) Assume $v_0 \in H^{r+1}$ where $r > \frac{d}{2}$, and assume $\nabla v_0 \in G_{s,r}$ for some $\delta > 0$ and Gevrey index $s > 1$. Then $\exists T > 0$ and a unique sol'n $(v, a) \in C([0, T], H^{r+1}(\mathbb{R}^d)) \times C([0, T], H^r(\mathbb{R}^d))$ s.t. $(\nabla v, a) \in L^\infty([0, T], G_{s, \delta})$

$$\|f\|_{G_{s, \delta}} = \sum_{m \geq 0} \frac{\delta^m}{m! s^m} \| \partial_1^m f \|_{H^r}$$

$s=1$: analytic
 $s > 1$: Gevrey

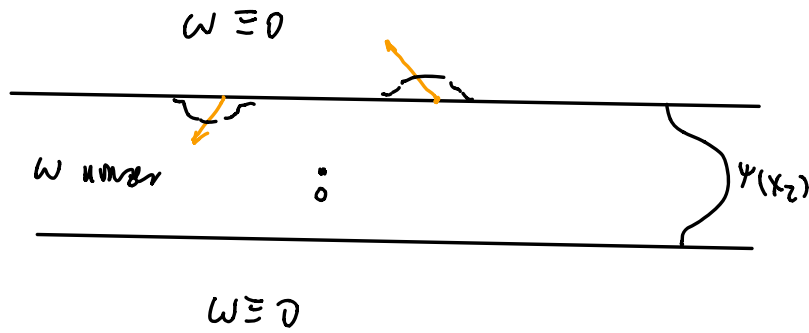
1) Anisotropic analytic regularity fails in the Eulerian framework:

Choose

$$\omega^{(k)}(x_1, x_2) = \tau_0 k e^{-k^2(x_1^2 + x_2^2)} \Psi(x_2), \quad k=1, 2, \dots$$

As $k \rightarrow \infty$ this converges to δ_0 as $k \rightarrow \infty$

Ψ : compactly supported in x_2 .



2) Ill-posedness in the analytic space $G_{1,0}$

Diperna-Majda: For any f, g , the fn

$$u(x_1, x_2, x_3, t) = (f(x_2), 0, g(x_1 - t f(x_2)))$$

solves the 3D Euler $\forall f, g$. Initial data

$$u_0(x_1, x_2, x_3) = (f(x_2), 0, g(x_1))$$

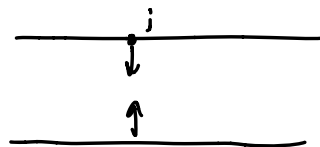
(Bardos-Titi: ill-posedness of the Euler in $C^{0,\alpha}$)

Take

$$g \propto \frac{1}{1+x^2}$$

complex singularity at i

$$f = \sin x$$



To get the correct g , we integrate four times and remove the low moments & periodize.

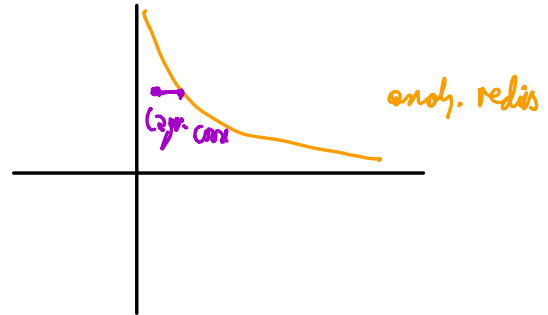
(2)

Ex (Shnirelman) : Cellular flow

$$u(x_1, x_2) = (\sin x_1, \cos x_2, -\cos x_1, \sin x_2)$$

stationary sol'n of the EE which is outside.

In Lagrangian coordinates :



Euler Equations with Free Interface

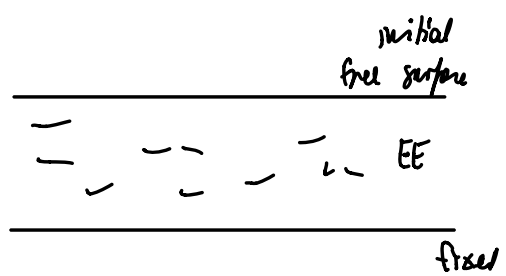
curl u_0 need not be zero.

Consider the EE in

$$\Omega = \mathbb{T}^2 \times (0, 1) \subseteq \mathbb{R}^3$$

$$\Gamma_1 = \mathbb{T} \times \{1\}$$

$$\Gamma_0 = \mathbb{T} \times \{0\}$$



Initially assume $h = 0$ (height) for simplicity.

Christodoulou - Lindblad '00

Coatland - Shkoller '07 (loc existence in H^3)

Shatah - Zeng '08 (loc existence in H^3)

Zhang - Zhong

K-Tuffaha ~ 14, K-Tuffaha - Vicol 17
 $u_0 \in H^{2.5+\delta}$, curl $u_0 \in H^3$

Wong et al (to appear) 2015 preprint
 $u_0 \in H^{2.5+\delta}$ ($\sigma = 0$)

Disconzi - K-Tuffaha (19) $u_0 \in H^{2.5+\delta}$, $\sigma > 0$,
& a cond. on free surface

(long time behavior: S.Wu, Alazard - Burq - Zuily, Ifrim - Tataru, ...)

let

$$\eta_t(y, t) = u(\eta(y, t), t)$$

$$\eta(y, 0) = y$$

and

$$a = (\nabla \eta)^{-1}$$

Also, denote

$$v(y, t) = u(\eta(y, t), t)$$

$$q(y, t) = p(\eta(y, t), t).$$

Then

$$\begin{cases} \partial_t v_i + a_{ki} \partial_k q = 0 \\ \partial_{ki} \partial_k v_i = 0 \end{cases}$$

$$\left(\begin{array}{l} \text{compare } \partial_t u + u \cdot \nabla u + \nabla p = 0 \\ \nabla \cdot u = 0 \end{array} \right)$$

The particle map η satisfies

$$\eta_t = v$$

$$\eta(y, 0) = y$$

The cofactor matrix $a = (\nabla \eta)^{-1}$ satisfies

$$\partial_t a = -a \nabla v a$$

$$a(y, 0) = I$$

$$\left(\begin{array}{l} a \nabla \eta = I \\ \partial_t \nabla \eta = -a \nabla \eta_t \\ \quad \quad \quad = -a \nabla v \\ \therefore \partial_t a = -a \nabla v a \end{array} \right)$$

So the system reads $(y \rightarrow x)$

$$\begin{cases} \partial_t v^i + \partial_{x_i} \partial_x q = 0 \\ \partial_{x_i} \partial_x v_i = 0 \\ \partial_t a = -a \nabla v \cdot a \\ \eta_t = v \end{cases} \begin{matrix} \text{(ODE)} \\ \text{(ODE)} \end{matrix}$$

with the initial conditions

$$\begin{cases} v(0) = v_0 \\ a(0) = I \\ \eta(x, 0) = x \end{cases} \text{i.c.}$$

Boundary condition:

$$\begin{aligned} q = 0 & \text{ on } \Gamma_1 = \mathbb{T}^2 \times \{1\} && \text{(no surface tension)} \\ v \cdot n = 0 & \text{ on } \Gamma_0 = \mathbb{T}^2 \times \{0\} \end{aligned}$$

$$H = \{v \in L^2 : \nabla \cdot v = 0, v \cdot n = 0 \text{ in } \Gamma_0\}$$

Theorem (KTV) Let $\delta > 0$ be arbitrary. Assume

$$v(0) \in H^{2\delta+\delta} \cap H$$

$$\text{Curl } v(0) \in H^{2+\delta}$$

and the Rayleigh-Taylor condition

$$\frac{\partial q}{\partial n}(x, 0) \leq -\frac{1}{C} \text{ on } \Gamma_1.$$

(25)

Then \exists sol'n $(v, \varrho, \eta, \varrho_t)$ s.t.

$$v \in L^\infty H^{2.5+\delta} \cap C H^{2+\delta} \quad (2.5+\delta)$$

$$\eta \in L^\infty H^{3+\delta} \quad (3+\delta)$$

$$v_t \in L^\infty H^{2+\delta} \quad (2+\delta)$$

$$\varrho \in L^\infty H^{3+\delta} \quad (3+\delta)$$

$$\varrho_t \in L^\infty H^{2.5+\delta} \quad (2.5+\delta)$$

in $[0, T]$, where T depends on the initial data.

A priori estimates:

1. Pressure estimates
2. Tangential estimates
3. div-curl estimates on η and v

the boundedness of coefficients.

Lemma Assume that

$$\| \nabla v \|_{L^\infty H^{2.5+\delta}} \leq \Omega \quad \leftarrow \quad \Omega = (2 \times \text{size of the initial data}) = 2 \| u_0 \|_{H^{2.5+\delta}}$$

if

$$T \leq \frac{1}{C\Omega}$$

then

$$\| \nabla \eta \|_{H^{1.5+\delta}} \lesssim 1 \quad \& \quad \| \varrho \|_{H^{1.5+\delta}} \lesssim 1 \quad (1)$$

doesn't depend on Ω

If $T \leq \frac{\varepsilon}{CM}$, then (2)

$$\|a - I\|_{H^{1.5+\delta}} \leq \varepsilon \quad \left(\Rightarrow \|a - I\|_{L^\infty} \leq \varepsilon \right) \quad (2)$$

and $\|a^T a - I\|_{H^{1.5+\delta}} \leq \varepsilon$ needed for the pressure (3)

Recall: $\eta_t = v$
 $a_t = -a \nabla v a$

All the statements are based on these two ODEs.

Proof (1): $(\nabla \eta_t = \nabla v)$

$$\nabla \eta = I + \int_0^t \nabla v \, ds$$

\Rightarrow

$$\|\nabla \eta\|_{H^{1.5+\delta}} \lesssim 1 + Mt \lesssim 1 + MT$$

So if $T \leq \frac{1}{CM} \simeq (1)_1$

For a : Jacobian = 1 $(\det \nabla \eta = 1)$

$$\|a\|_{H^{1.5+\delta}} \lesssim \|\nabla \eta\|_{H^{1.5+\delta}}^2 \lesssim 1$$

(2) We have

$$a - I = \int_0^t a_t \, ds = - \int_0^t a \nabla v a \, ds$$

\therefore

$$\|a - I\|_{H^{1.5+\delta}} \lesssim \int_0^t \|a \nabla v a\|_{H^{1.5+\delta}} \, ds \leq MT$$

$$MT \leq \varepsilon \quad \Rightarrow \quad \|a - I\|_{H^{1.5+\delta}} \lesssim \varepsilon$$

(27)

(3) We have

$$\begin{aligned} \partial \partial^T - I &= \int_0^t \partial_t \partial^T + \int_0^t \partial \partial_t^T \\ &= - \int_0^t \partial \nabla v \partial \partial^T - \int_0^t \partial \partial^T (\nabla v)^T \partial^T \end{aligned}$$

$$\Rightarrow \|\partial \partial^T - I\|_{H^{1.5+\epsilon}} \lesssim \int_0^T \|\nabla v\|_{H^{1.5+\epsilon}} \leq MT. \quad \square$$

Pressure Estimates

Take the div of the E_v and obtain the Laplace eq. for q (or p)

$$(u_t + u \cdot \nabla u + \nabla p = 0$$

$$\Rightarrow \nabla \cdot u_t + \nabla \cdot (u \cdot \nabla u) + \Delta p = 0$$

$$\therefore \Delta p = - \nabla \cdot (u \cdot \nabla u) = - \partial_j (u_i \partial_i u_j)$$

$$\Delta p = - \partial_j (u_i \partial_i u_j)$$

$$\Delta p = - \partial_{ij} (u_i u_j)$$

(useful if no bdry)

$$\Delta p = - \partial_j u_i \partial_i u_j$$

$$(\Delta p = - \nabla (u \cdot \nabla u) , \Delta p = D^2 (uu) , \Delta p = \nabla_u \nabla_u)$$

Apply $\partial_{j_i} \partial_j$ to

$$\begin{cases} \partial_t v_i + \partial_{k_i} \partial_k q = 0 \\ \partial_{k_i} \partial_k v_i = 0 \end{cases}$$

We get

$$\partial_{j_i} \partial_j (\partial_{k_i} \partial_k q) = - \partial_{j_i} \partial_j \partial_t v_i$$

$$\therefore \partial_{j_i} \partial_j (\partial_{k_i} \partial_k q) = \partial_t \partial_{j_i} \partial_j v_i$$

Close to

$$\begin{aligned} & \uparrow \\ & \partial_{j_i} \partial_j (\partial_{k_i} \partial_k q) \\ & \parallel \\ & \Delta q \end{aligned}$$

($\partial_t a$ is "good"
since we eliminate
one time derivative
from the term
 $\|a\| \lesssim 1 + \int_0^t \|v\|$
 $\|\partial_t\| \lesssim \|v\|$)

By the Poincaré identity

$$\partial_j \partial_{j_i} = 0$$

Cofactor matrices: the divergence of every column is 0

(Evans, p. 462).

$$\begin{aligned} \partial_j (\underbrace{\partial_{j_i} \partial_{k_i}}_{\uparrow}) \partial_k q &= \partial_t \partial_{j_i} \partial_j v_i \\ & (\partial a^T)_{j_i} \quad (\approx \delta_{ij} \text{ within } \epsilon \text{ in } H^{1.5+1}) \end{aligned}$$