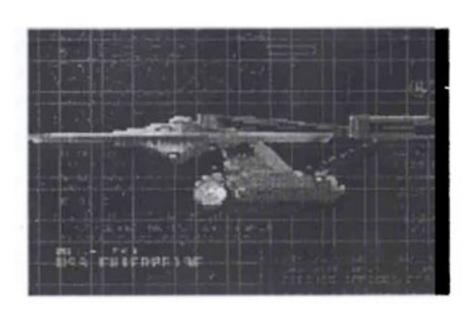
# Quantum Celeportation and Entanglement Rurification

**Brassard**Université de Montréal

# Additional Credit

- Charles H. Bennett Samuel L. Braunstein Richard Cleve
- Claude Crépeau
   David Deutsch
   David P. DiVincenzo
   Artur K. Ekert
- Richard Jozsa
   Chiara Macchiavello
- Asher Peres
- Sandu Popescu Anna Sanpera
- Ben Schumacher
- John A. Smolin
- William K. Wootters







For years, she shared recipes with her friend in Osaka.

She showed him hundreds of ways to use paprika.

He shared his secret recipe for sukiyaki.

One day Margit e-mailed Seiji,

# telepo

Margit is a little premature, but we are working on it.

An IBM scientist and his colleagues have discovered a way to make an object disintegrate in one place and reappear intact in another.

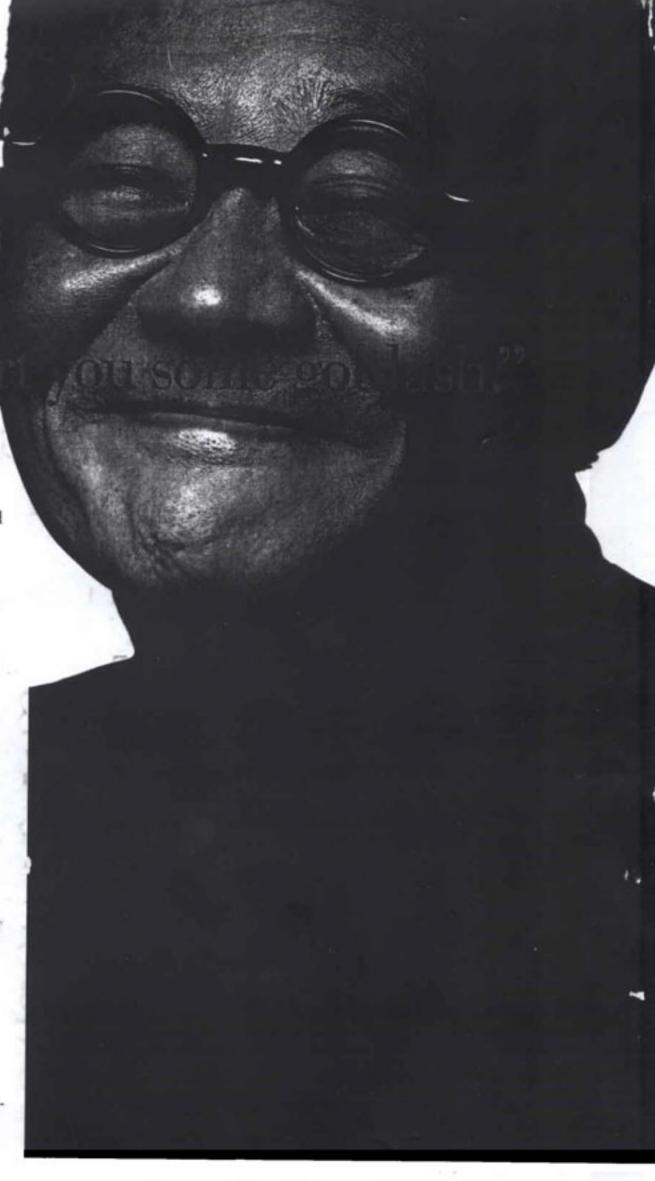
It sounds like magic.
Out their breakthrough could affect everything from the future of computers to our knowledg, of the cos. .os.

Smart guys. But none of them can stuff a cabbage.

Yet.

IBM

Solutions for a small planet



# IBM Advertisements Boldly Go Where No Study Has Gone Before

By THOMAS E. WEBER

Staff Reporter of THE WALL STREET JOURNAL Somebody at International Business Machines Corp. may have been watching

too much "Star Trek."

varau.

IBM ads are touting an impressive new research project: "An IBM scientist and his colleagues have discovered a way to make an object disintegrate in one place and reappear intact in another," the ads read, evoking the sci-fi "transporter" used by Captain Kirk and his friends.

Then there's the elderly lady pictured in the ad, who proclaims to a gourmet pen pal, "Stand by. I'll teleport you some goulash." Her promise may be "a little premature," the ad says, but IBM is

working on it.

Really? Robert L. Park, a physics professor at the University of Maryland, was so intrigued by the ad, which ran in Rolling Stone, Scientific American and other magazines, that he e-mailed IBM for more information. "This is still under development and no further information is currently available," came the curt reply.

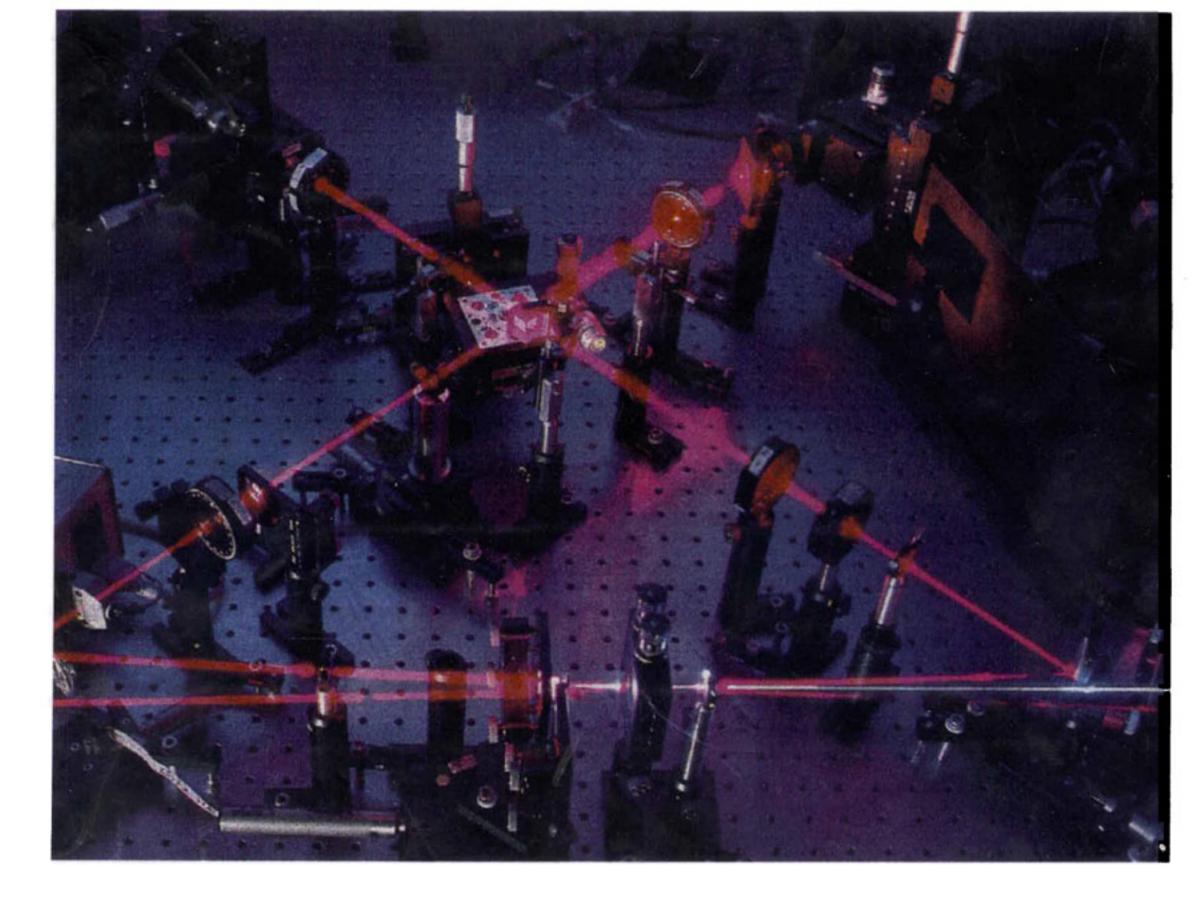
It turns out that an IBM scientist hasn't really "discovered" anything of the sort. Mr. Park is still stewing about it. "It's sort of a bait-and-switch," he complains. "It

leaves a taste of bad goulash in my mouth." The IBM scientist in question, one Charles H. Bennett, is a physicist at IBM's vaunted research center in Yorktown Heights, N.Y. But he concentrates more on photons — "a quantum of electromagnetic energy having both particle and wave properties," the dictionary says — than on Hungarian dishes.

"This doesn't really have anything to do with teleporting goulash," he says. It does have to do with an area of physics known as quantum teleportation, on which Mr. Bennett is happy to expound at length. Never mind what he said, exactly; Mr. Spock would have trouble following it.

At ad agency Ogilvy & Mather, executive group director Matt Ross describes the goulash gimmick as "a kind of dramatization" meant to emphasize IBM's research and development capabilities. "It's really just a fun way to evoke a sense of what this could mean."

But it hasn't been all fun for Mr. Bennett. He says he's more than a little embarrassed about the ad's unequivocal claims. "In any organization there's a certain tension between the research end and the advertising end," he says. "I struggled with them over it, but perhaps I didn't struggle hard enough."



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# QUANTUM Teleportation

The Future of Travel?
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### Of Hice and Mensa

Genetic formula for a smarter mouse

### **Brown Dwarfs**

Stars that fizzled fill the galaxy

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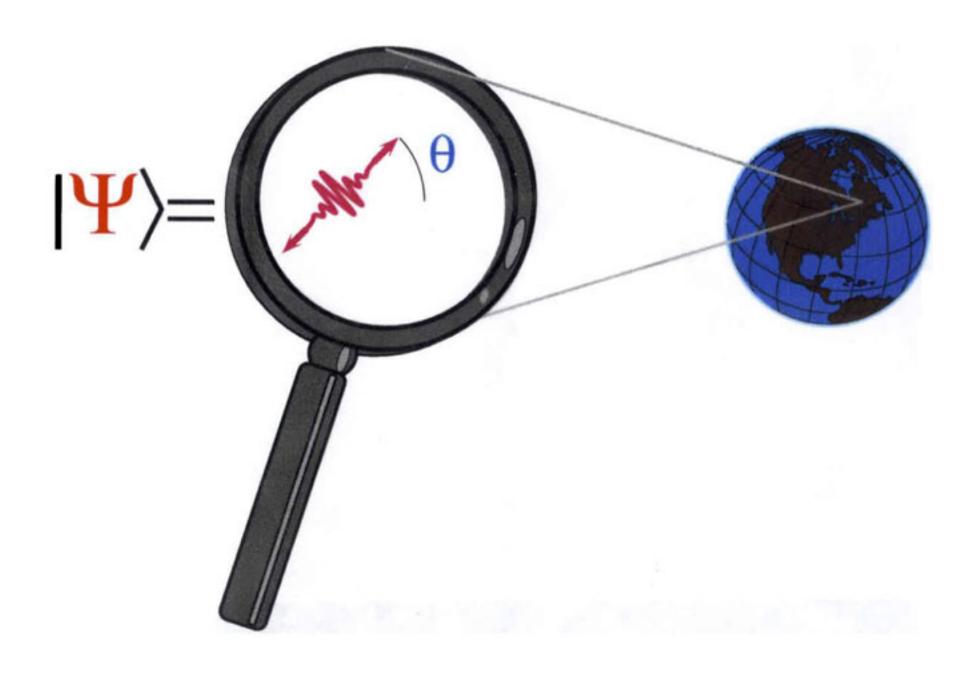
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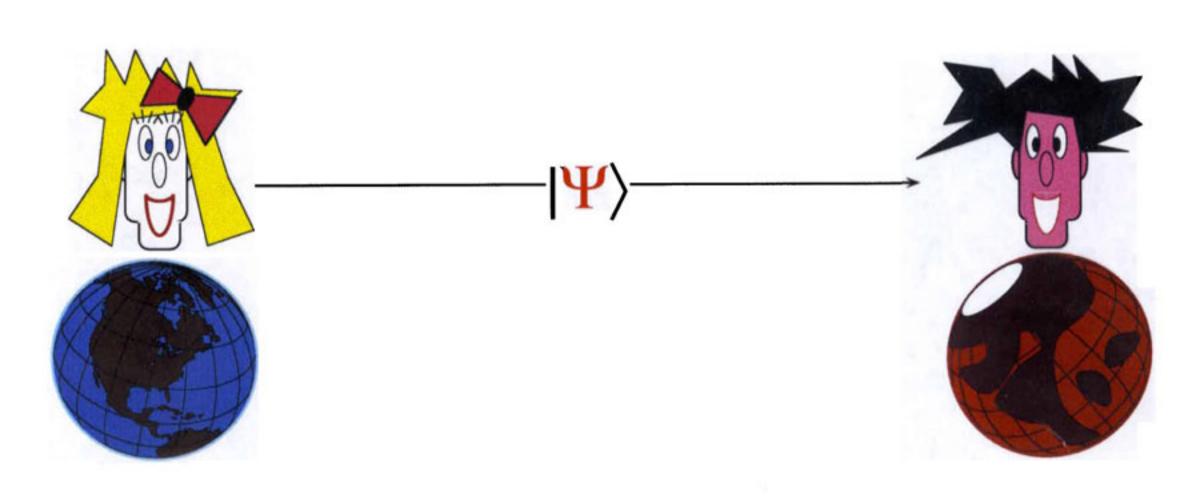
### **Quantum Information**

- Cannot be cloned or copied
- Cannot be broadcast
- Cannot be measured reliably
- Is disturbed by observation
- Sometimes appears to propagate instantaneously
- Can exist in superposition of classical states

## Unknown Quantum State | 4>

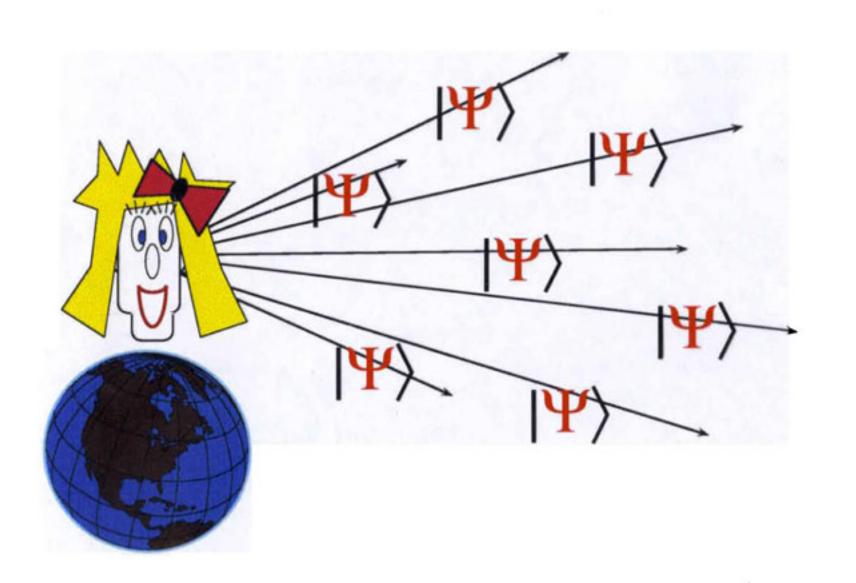


# Unknown Quantum State | 4>



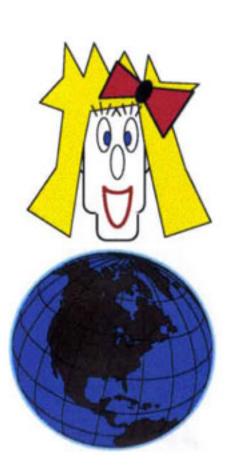
# Impossibility of Measuring \P

### Impossibility of Cloning/Broadcasting |Ψ⟩





## Transmitting |Ψ⟩ Physically







# Noisy Channel

 ${\mathscr{A}}$ lice

$$|\Psi\rangle$$
 —  $|\Psi\rangle$  —  $|\Psi\rangle$ 

### The Quantum Bit: Qubit

Quantum bit: 
$$|\Psi\rangle = \alpha |0\rangle + \beta |1\rangle$$

$$\alpha$$
 and  $\beta$  are called amplitudes

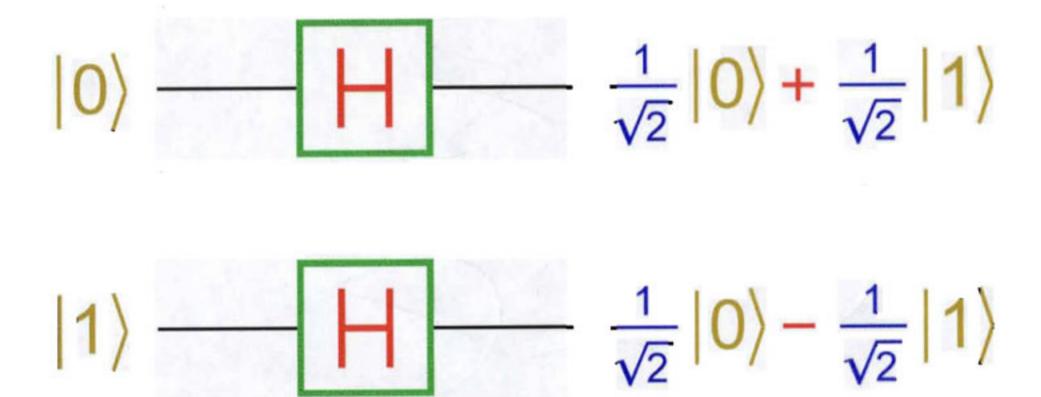
- + denotes the quantum superposition
- α and β are complex numbers

$$|\alpha|^2 + |\beta|^2 = 1$$

# Measurement Gate

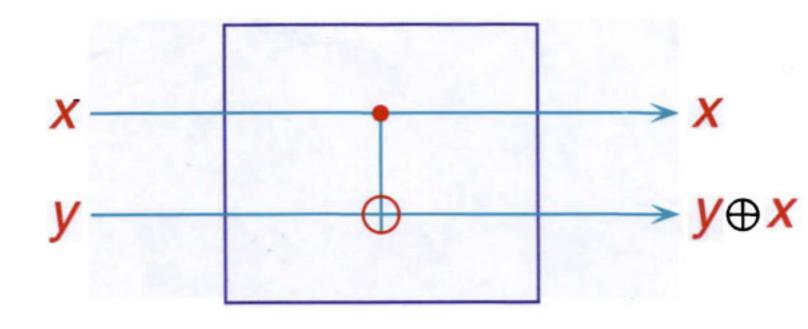
$$\alpha |0\rangle + \beta |1\rangle - M - \{0\}$$
 with prob  $|\alpha|^2$  |1 $\rangle$  with prob  $|\beta|^2$ 

### Walsh-Hadamard Transform



# Quantum XOR

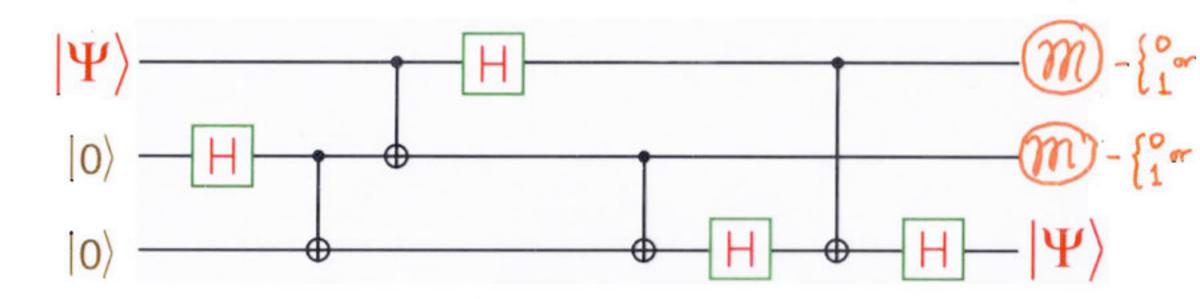
### (Controlled NOT)



$$y \oplus x = \begin{cases} y & \text{if } x = 0 \\ 1 - y & \text{if } x = 1 \end{cases}$$

$$|00\rangle \longrightarrow |00\rangle \qquad |10\rangle \longrightarrow |11\rangle |01\rangle \longrightarrow |01\rangle \qquad |11\rangle \longrightarrow |10\rangle$$

### **Quantum Circuit**

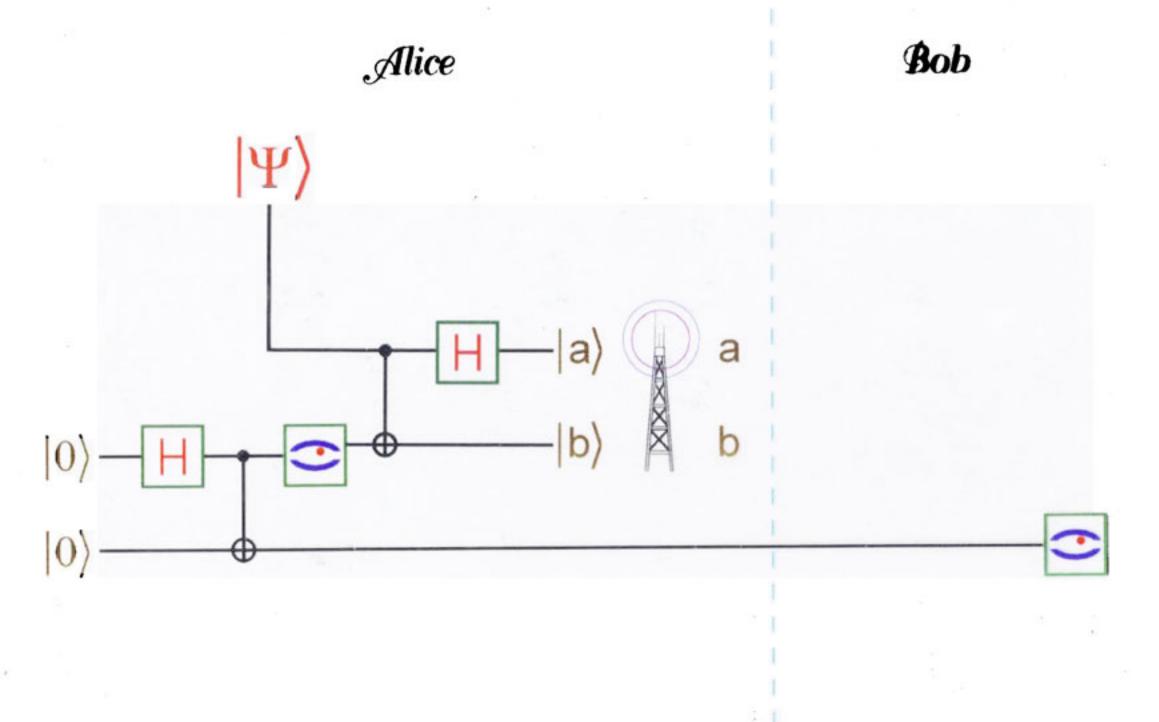


### To Measure or not to Measure

Robert Griffiths and Chi-Sheng Niu

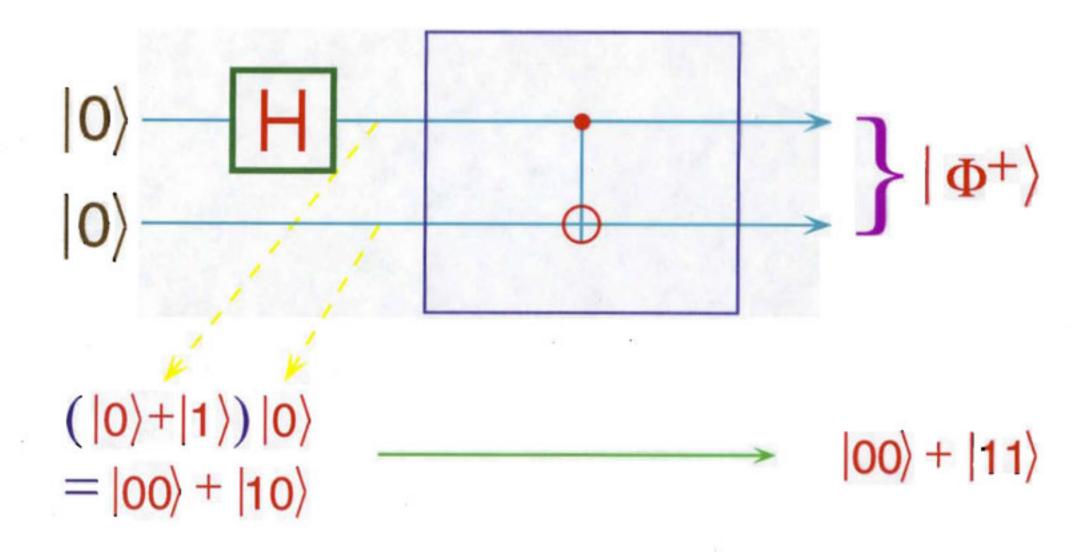
$$\alpha |0\rangle + \beta |1\rangle$$
  $M$  with prob  $|\alpha|^2$   $|1\rangle$  with prob  $|\beta|^2$ 



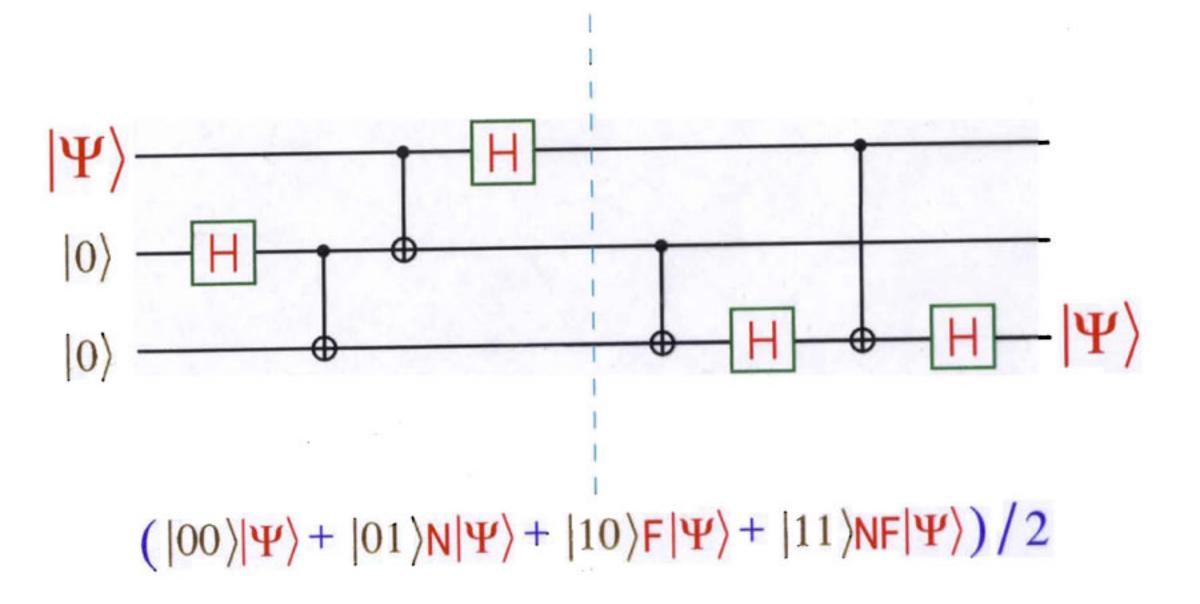


 ${\mathcal{A}\!\mathit{lice}}$  $\mathcal{B}ob$  $|a\rangle$ a a

# Creation of Entanglement



(up to normalization)



$$\mathsf{N} \left\{ \begin{array}{l} |0\rangle \longrightarrow |1\rangle \\ |1\rangle \longrightarrow |0\rangle \end{array} \right. \qquad \mathsf{F} \left\{ \begin{array}{l} |0\rangle \longrightarrow |0\rangle \\ |1\rangle \longrightarrow -|1\rangle \end{array} \right.$$

# Measurement Gate

$$\alpha |0\rangle + \beta |1\rangle$$
 —  $M$  with prob  $|\alpha|^2$   $|1\rangle$  with prob  $|\beta|^2$ 

Same if input is entangled within a system in state

$$\alpha |0\rangle |\Phi\rangle + \beta |1\rangle |\Psi\rangle$$

for arbitrary normalized  $|\Phi\rangle$  and  $|\Psi\rangle$ 

Bob Alice

# Quantum Entanglement Zurification

Gilles Brassard

Université de Montréal

### Purification of Noisy Entanglement and Faithful Teleportation via Noisy Channels

Charles H. Bennett, 1,\* Gilles Brassard, 2,† Sandu Popescu, 3,‡ Benjamin Schumacher, 4,§ John A. Smolin, 5,|| and William K. Wootters 6,9

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<sup>3</sup>Physics Department, Tel Aviv University, Tel Aviv, Israel

<sup>4</sup>Physics Department, Kenyon College, Gambier, Ohio 43022

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<sup>6</sup>Physics Department, Williams College, Williamstown, Massachusetts 01267

(Received 24 April 1995)

Two separated observers, by applying local operations to a supply of not-too-impure entangled states (e.g., singlets shared through a noisy channel), can prepare a smaller number of entangled pairs of arbitrarily high purity (e.g., near-perfect singlets). These can then be used to faithfully teleport unknown quantum states from one observer to the other, thereby achieving faithful transmission of quantum information through a noisy channel. We give upper and lower bounds on the yield D(M) of pure singlets  $(|\Psi^-\rangle)$  distillable from mixed states M, showing D(M) > 0 if  $\langle \Psi^-|M|\Psi^-\rangle > \frac{1}{2}$ .

PACS numbers: 03.65.Bz, 42.50.Dv, 89.70.+c



Perfect with probability 1/2
Randomized with prob 1/2

# **Bell States**

$$|\Psi^{-}\rangle = |10\rangle - |01\rangle$$

$$|\Psi^{+}\rangle = |10\rangle + |01\rangle$$

$$|\Phi^{-}\rangle = |11\rangle - |00\rangle$$

$$|\Phi^{+}\rangle = |11\rangle + |00\rangle$$

(up to normalization)

**Note:** only  $|\Psi^-\rangle$  is basis independent

# Werner State

$$\frac{1}{2} | \Phi^{+} \rangle \langle \Phi^{+} | + \frac{1}{2} \left[ \frac{1}{4} | 00 \rangle \langle 00 | + \frac{1}{4} | 01 \rangle \langle 10 | + \frac{1}{4} | 11 \rangle \langle 11 | \right] + \frac{1}{4} | 10 \rangle \langle 10 | + \frac{1}{4} | 11 \rangle \langle 11 | \right]$$

$$\frac{1}{4} |00\rangle\langle 00| + \frac{1}{4} |01\rangle\langle 10| + \frac{1}{4} |10\rangle\langle 10| + \frac{1}{4} |11\rangle\langle 11|$$

$$\frac{1}{4} | \Phi^{+} \rangle \langle \Phi^{+} | + \frac{1}{4} | \Phi^{-} \rangle \langle \Phi^{-} |$$

$$+ \frac{1}{4} | \Psi^{-} \rangle \langle \Psi^{-} | + \frac{1}{4} | \Psi^{+} \rangle \langle \Psi^{+} |$$

## **Density Matrices**

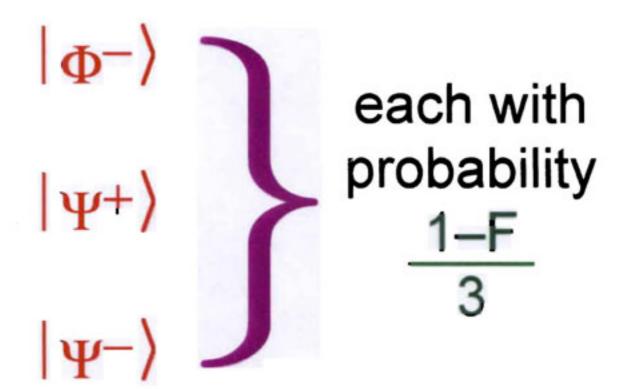
$$|W_{5/8}\rangle = \frac{5}{8} |\Phi^{+}\rangle\langle \Phi^{+}| + \frac{1}{8} |\Phi^{-}\rangle\langle \Phi^{-}| + \frac{1}{8} |\Psi^{-}\rangle\langle \Psi^{-}| + \frac{1}{8} |\Psi^{+}\rangle\langle \Psi^{+}| =$$

$$\frac{1}{2} |\Phi^{+}\rangle \langle \Phi^{+}| + \frac{1}{2} \left[ \frac{1}{4} |\Phi^{+}\rangle \langle \Phi^{+}| + \frac{1}{4} |\Phi^{-}\rangle \langle \Phi^{-}| + \frac{1}{4} |\Psi^{-}\rangle \langle \Psi^{-}| + \frac{1}{4} |\Psi^{+}\rangle \langle \Psi^{+}| \right]$$

# Werner State

# 2 qubits jointly in state

Φ<sup>+</sup>) with probability F



$$(F = 5/8)$$

## Distinguishing Bell States

$$|\Psi^{\pm}\rangle$$
  $\left\{\begin{array}{c} M & 0 & 1 \\ M & 1 & 0 \end{array}\right.$  or  $\left.\begin{array}{c} M \\ \Phi^{\pm} \end{array}\right\}$   $\left\{\begin{array}{c} M & 0 \\ M & 0 \end{array}\right.$  or  $\left.\begin{array}{c} 1 \\ 0 \\ M & 0 \end{array}\right.$ 

### Manipulating Bell States

$$|\Psi^{-}\rangle$$
  $\left\{\begin{array}{c} |0\rangle \rightarrow |1\rangle \\ |1\rangle \rightarrow -|0\rangle \end{array}\right\} |\Phi^{+}\rangle$ 

$$|01\rangle - |10\rangle$$
  $|11\rangle + |00\rangle$ 

$$|\Psi^{+}\rangle \longleftrightarrow |\Phi^{+}\rangle$$

### Manipulating Bell States

$$|\Psi^{-}\rangle\left\{egin{array}{c} |0
angle
ightarrow|0
angle-i|1
angle \\ |1
angle
ightarrow|1
angle-i|0
angle \\ |0
angle
ightarrow|1
angle-i|0
angle \\ |1
angle
ightarrow|0
angle-i|1
angle \end{array}
ight\}|\Phi^{-}
angle$$

$$|\Psi^{-}\rangle \longleftrightarrow |\Phi^{-}\rangle$$

 $|\Psi^{+}\rangle$  and  $|\Phi^{+}\rangle$  are unaffected

(up to phase)

#### Bilateral XOR

$$|\Psi^{-}\rangle \left\{ \begin{array}{c} -A_1 \\ -B_1 \end{array} \right\} |\Psi^{+}\rangle$$

$$|\Psi^{-}\rangle \left\{ \begin{array}{c} -A_2 \\ -B_2 \end{array} \right\} |\Phi^{-}\rangle$$

$$(|01\rangle - |10\rangle) (|01\rangle - |10\rangle) = |0100\rangle - |0111\rangle - |1011\rangle + |1000\rangle$$
  
 $|0101\rangle - |0110\rangle - |1001\rangle + |1010\rangle = (|01\rangle + |10\rangle) (|00\rangle - |11\rangle)$ 

#### Bilateral XOR

$$\begin{vmatrix} \Phi^{+} \rangle \left\{ \begin{array}{c} -A_{1} \\ -B_{1} \end{array} \right\} \begin{vmatrix} \Phi^{+} \rangle \\ |\Phi^{+} \rangle \left\{ \begin{array}{c} -A_{2} \\ -B_{2} \end{array} \right\} \begin{vmatrix} \Phi^{+} \rangle \\ -B_{2} \end{vmatrix}$$

$$(|00\rangle + |11\rangle) (|00\rangle + |11\rangle) = |0000\rangle + |0011\rangle + |1111\rangle + |1100\rangle = (|00\rangle + |0011\rangle + |1111\rangle) = (|00\rangle + |11\rangle) (|00\rangle + |11\rangle)$$

before		auahahilit.	after		target	example
source	target	probability	source	target	$= \Phi\rangle$ ?	F = 5/8
$ \Phi^+\rangle$	$ \Phi^{+}\rangle$	$F^2$	$ \Phi^{+}\rangle$	$ \Phi^+\rangle$	YES	25/64
$ \Phi^+\rangle$	$ \Phi^{-}\rangle$	F(1-F)/3	$ \Phi^{-}\rangle$	$ \Phi^-\rangle$	YES	5/64
$ \Phi^+\rangle$	$ \Psi^{+}\rangle$	F(1-F)/3	$ \Phi^{+}\rangle$	$ \Psi^+\rangle$		5/64
$ \Phi^+\rangle$	$ \Psi^{-}\rangle$	F(1-F)/3	$ \Phi^{-}\rangle$	$ \Psi^{-}\rangle$		5/64
$ \Phi^-\rangle$	$ \Phi^{+}\rangle$	F(1-F)/3	$ \Phi^{-}\rangle$	$ \Phi^+\rangle$	YES	5/64
$ \Phi^-\rangle$	$ \Phi^-\rangle$	$(1-F)^2/9$	$ \Phi^{+}\rangle$	$ \Phi^-\rangle$	YES	1/64
$ \Phi^-\rangle$	$ \Psi^{+}\rangle$	$(1-F)^2/9$	$ \Phi^{-}\rangle$	$ \Psi^+\rangle$		1/64
$ \Phi^-\rangle$	$ \Psi^{-}\rangle$	$(1-F)^2/9$	$ \Phi^{+}\rangle$	$ \Psi^{-}\rangle$		1/64
$ \Psi^+\rangle$	$ \Phi^+\rangle$	F(1-F)/3	$ \Psi^{+}\rangle$	$ \Psi^+\rangle$		5/64
$ \Psi^+\rangle$	$ \Phi^{-}\rangle$	$(1-F)^2/9$	$ \Psi^{-}\rangle$	$ \Psi^{-}\rangle$		1/64
$ \Psi^{+}\rangle$	$ \Psi^{+}\rangle$	$(1-F)^2/9$	$ \Psi^{+}\rangle$	$ \Phi^+\rangle$	YES	1/64
$ \Psi^+\rangle$	$ \Psi^{-}\rangle$	$(1-F)^2/9$	$ \Psi^{-}\rangle$	$ \Phi^-\rangle$	YES	1/64
$ \Psi^{-}\rangle$	$ \Phi^{+}\rangle$	F(1-F)/3	$ \Psi^{-}\rangle$	$ \Psi^+ angle$		5/64
$ \Psi^{-}\rangle$	$ \Phi^{-}\rangle$	$(1-F)^2/9$	$ \Psi^{+}\rangle$	$ \Psi^- angle$		1/64
$ \Psi^{-}\rangle$	$ \Psi^{+}\rangle$	$(1-F)^2/9$	$ \Psi^{-}\rangle$	$ \Phi^+\rangle$	YES	1/64
$ \Psi^-\rangle$	$ \Psi^{-}\rangle$	$(1-F)^2/9$	$ \Psi^{+}\rangle$	$ \Phi^-\rangle$	YES	1/64
rob kee		P				40/64
8.00	$d\ket{\Phi^+}$	Q				26/64

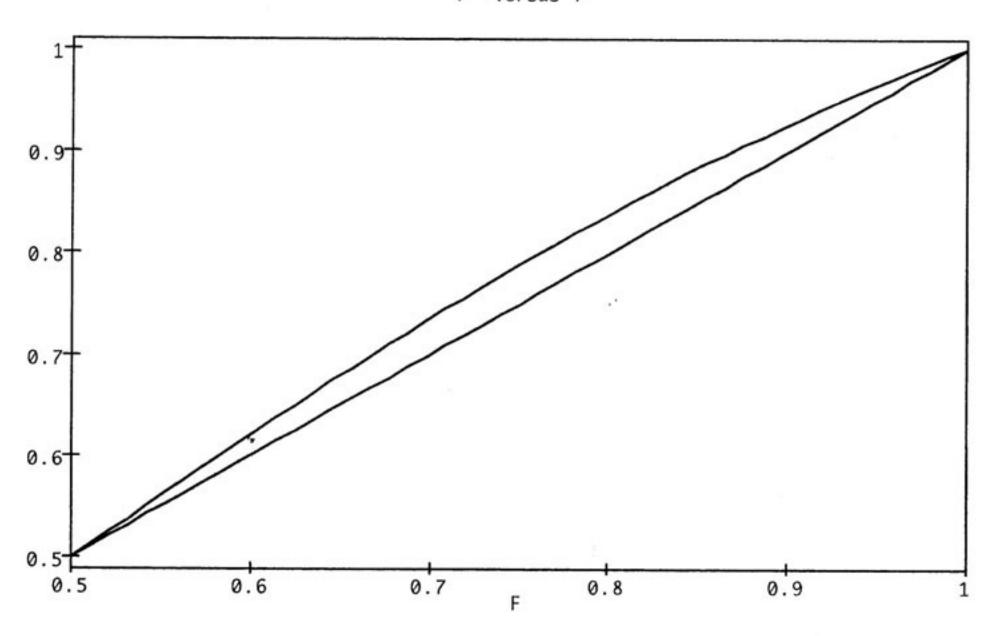
Prob keep 
$$P$$
 40/64 Keep and  $|\Phi^+\rangle$   $Q$  26/64  $|\Phi^+\rangle$  if kept  $F'=Q/P$  13/20

$$P = F^{2} + \frac{2F(1-F)}{3} + \frac{5(1-F)^{2}}{9} = \frac{8F^{2} - 4F + 5}{9}$$

$$Q = \mathbf{F}^4 + \frac{(1-\mathbf{F})^2}{9} = \frac{8\mathbf{F}^2 - 4\mathbf{F} + 5}{9}$$

$$F' = \frac{10F^2 - 2F + 1}{8F^2 - 4F + 5} > F$$
 provided  $F > 1/2$ 

F' versus F



## Repeat Process?

Round	Fideli	ty
0	62.5%	(5/8)
1	65%	(13/20)
2	59.146	3%
3	51.943	0%

## What's Wrong?

#### After First Round:

Φ<sup>+</sup>) with probability 13/20

 $|\Phi^-\rangle$  with probability 1/4

Ψ+) with probability 1/20

Ψ-) with probability 1/20

Not a Werner State!

# Solution (1)

Wernerize between rounds

$$|\Psi^{+}\rangle \longleftrightarrow |\Psi^{-}\rangle$$
 or  $|\Psi^{-}\rangle \longleftrightarrow |\Phi^{-}\rangle$  or  $|\Phi^{-}\rangle \longleftrightarrow |\Psi^{+}\rangle$ 

each with probability 1/3

#### After Wernerizing First Round:

$$|\Phi^+\rangle$$
 with probability 13/20 each with probability  $|\Psi^+\rangle$  probability 7/60

### Result of Wernerization

Round	<b>Fidelity</b>	Rate
0	62.5%	100.00%
1	65.0%	31.25%
2	67.9%	10.03%
3	71.2%	3.33%
4	74.8%	1.15%
•••		•••

15

 $99.1\% 1.29 \times 10^{-6}$ 

# Solution (2)

Macchiavellize between rounds

$$|\Psi^{-}\rangle \longleftrightarrow |\Phi^{-}\rangle$$

Round	<b>Fidelity</b>	Rate
0	62.5%	100.00%
1	65.0%	31.25%
2	73.3%	9.06%
3	82.6%	2.96%
4	89.3%	1.13%
5	96.7%	0.47%
6	99.3%	0.22%
		≈1/455