

CHAOTIC FIELD THEORY

$$\phi = [\bar{\psi}_{\alpha\alpha}(x), \bar{\psi}^{\alpha\alpha}(x), A_{i\mu}(x)]$$

$$S[\phi] = \bar{\psi}(\gamma + \not{A} + m)\psi + \frac{1}{4}F^2$$

have:

$$Z[J] = \int [d\phi] e^{\frac{i}{\hbar}[S + \phi \cdot J]}$$

proposal:

$$e^{\Gamma[\phi] + \phi \cdot J} \approx \sum_{sc} e^{\frac{i}{\hbar}[S_{sc} + \frac{\pi}{4}m + \phi_{ii} \cdot J + \sum \hbar^n \Gamma^{(n)}]} \frac{1}{|\det S''|^{1/2}} +$$

$$\approx \sum_{sc} \left(\begin{array}{l} \text{turbulent} \\ \text{classical} \\ \text{solutions} \end{array} \right) + \begin{array}{l} \text{(tunneling)} \\ \text{(diffraction)} \end{array}$$

NOT: free \rightarrow , \sim

instantons, ...

"e^T" unstable patterns: numerical

- * 1-d field theories, dissipative
- * search for spatiotemporal recurrences
- * high d Hamiltonian flows ?

$$\sum_P^{\text{saddles}} \frac{e^{\frac{i}{\hbar} S_p + \frac{i\pi}{4} m_p + J\phi}}{|\det S''|^{1/2}} \left(1 + \sum^{\text{Feynman}} \hbar^n \Gamma^{(n)} + \dots \right)$$

- * 1-d map + noise
 - Feynman diagrams
 - Feynman + Poincaré
 - Transfer operator
- * gauge invariance ?
- * renormalization ?

Go with the flow

desiderata for a "simplest" field theory:

- 1 space, 1 time dim.
- scalar field
- 1 "Raynold's" parameter
"laminar" \leftrightarrow turbulent
- UV finite (smooth solutions)
- IR finite (compact support)

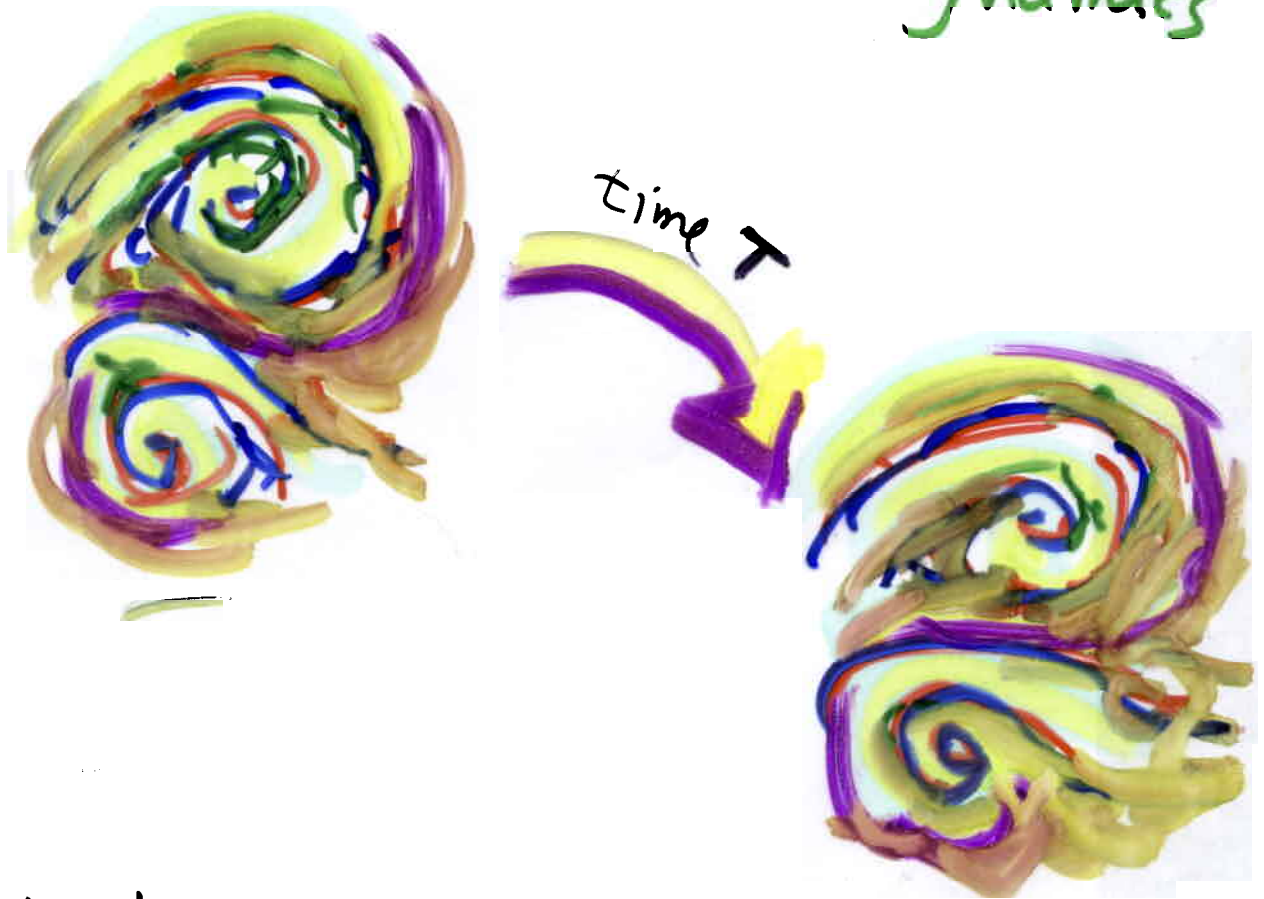
most popular

$$\dot{\phi} = L(\phi, \partial\phi, \dots) + N(\phi, \partial\phi, \dots)$$

- Nonlin Schrödinger
- Burgers
- complex Ginzburg-Landau
- Kuramoto-Sivashinski

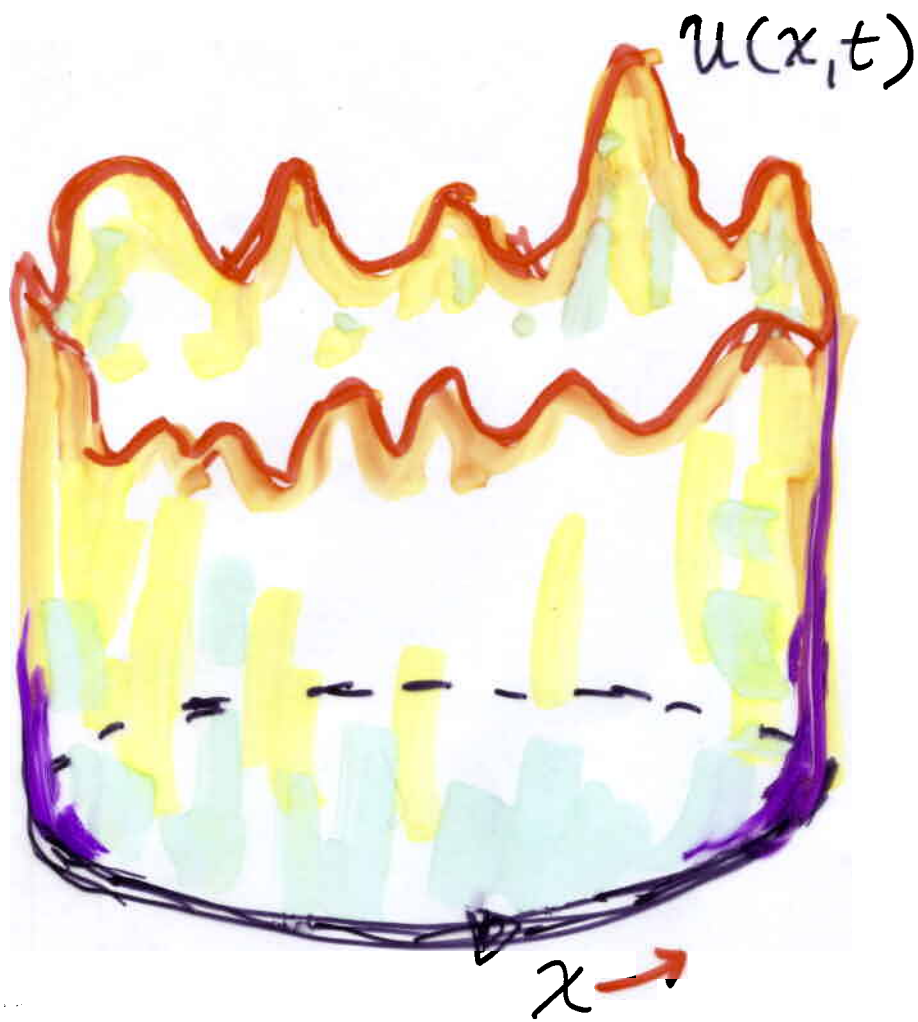
Hopf's last hope

classical theory of turbulent
dynamics



in terms of spatiotemporally
recurrent patterns

Burning flame front



$$u(x,t) = u(x+2\pi,t)$$

STRATEGY

(V. Putkaradze, F. Christiansen, P.C.)

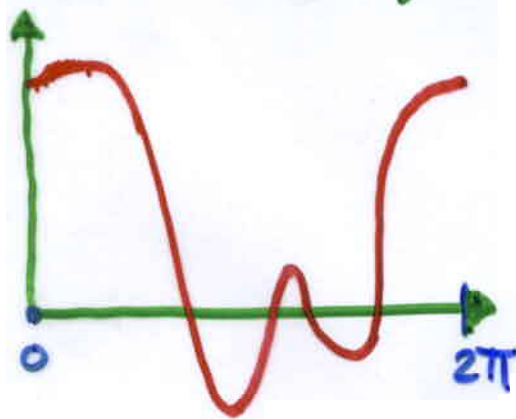
Nonlinearity 10, 1 (1997).

take "simplest" spatiotemporally
 "turbulent" dynamical system of
 physical interest ("parabola" of PDE's):

Kuramoto-Sivashinsky (1976):

$$u_t = \underbrace{(u^2)_x}_{\text{nonlinear}} - u_{xx} - \underbrace{\gamma u_{xxxx}}_{\text{damped}} \quad x \in [-\pi, \pi]$$

$u(x,t)$ = velocity of a 1d flame front



at given instant t

(slide)

Theorem (Temam, ...)

$\nu \neq 0$ attractor finite dimensional

viscosity

$\nu > 1$ death

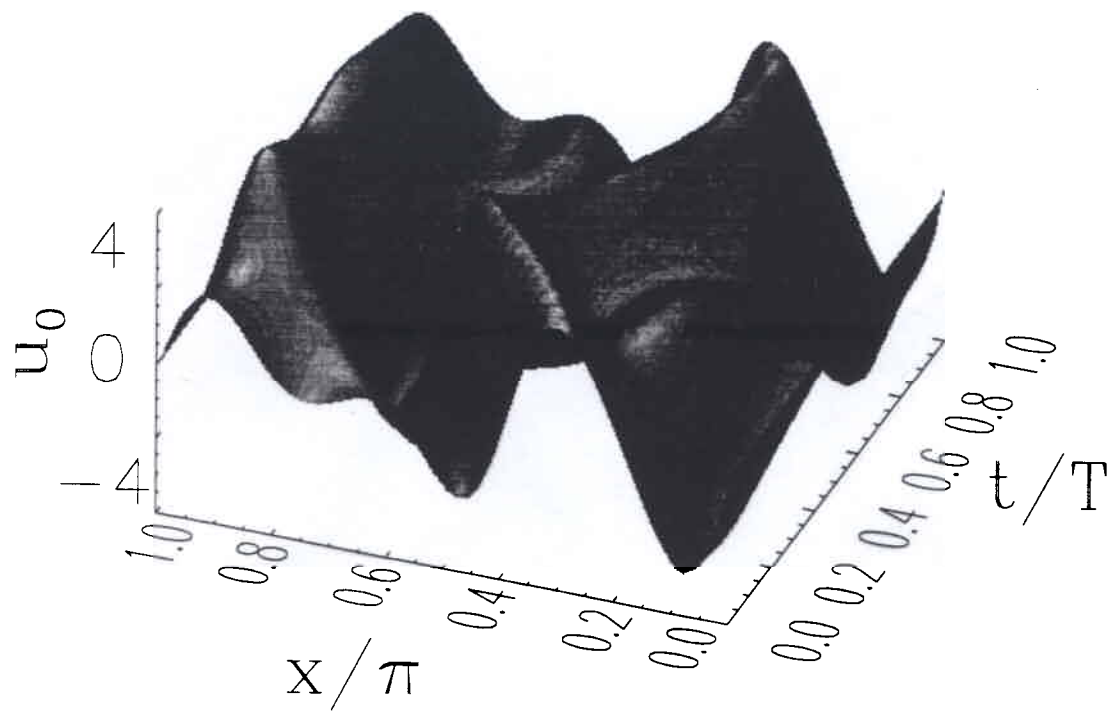
$\nu \leq 1$ spatio-temporal chaos

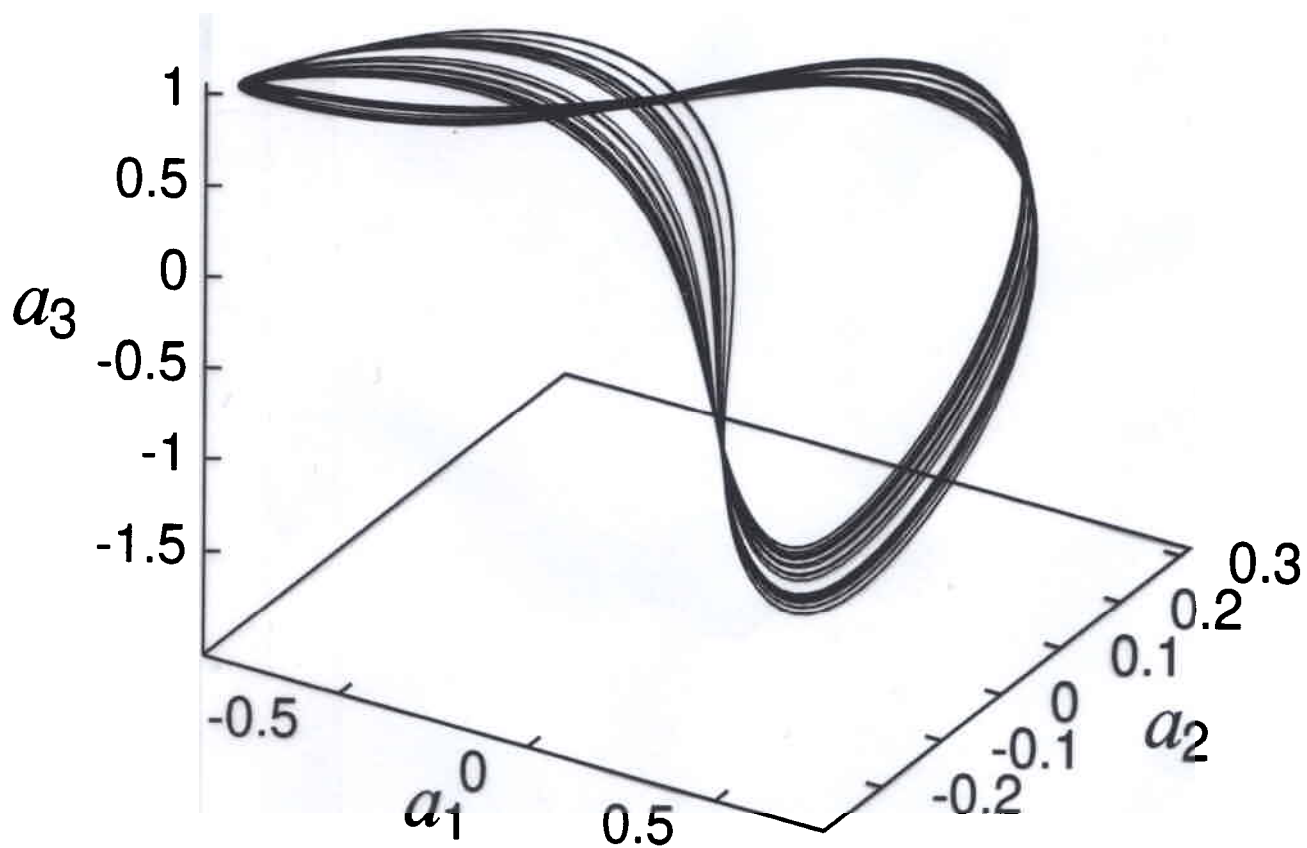
$\nu \ll 1$ turbulence

difficult because:

$$u(x,t) \rightarrow (a_1, \dots, a_N)$$

$$N = 15 - 1000$$





accomplished so far:

1. determined "all" unstable spatiotemporally periodic solutions $T_p < T_{\text{cutoff}}$

$$u_p(x, t) = u_p(x, t + T_p) \quad p = 1, 2, \dots, "1000"$$

2. described how $\langle \text{average} \rangle$ s are computed (not linear superpositions of p 's)

open questions

3. unclear if useful as viscosity $\nu \rightarrow 0$
("strong turbulence")

hope: periodic orbit skeleton remains sufficiently dense to explore the ∞ dim space of patterns?

periodic orbit p

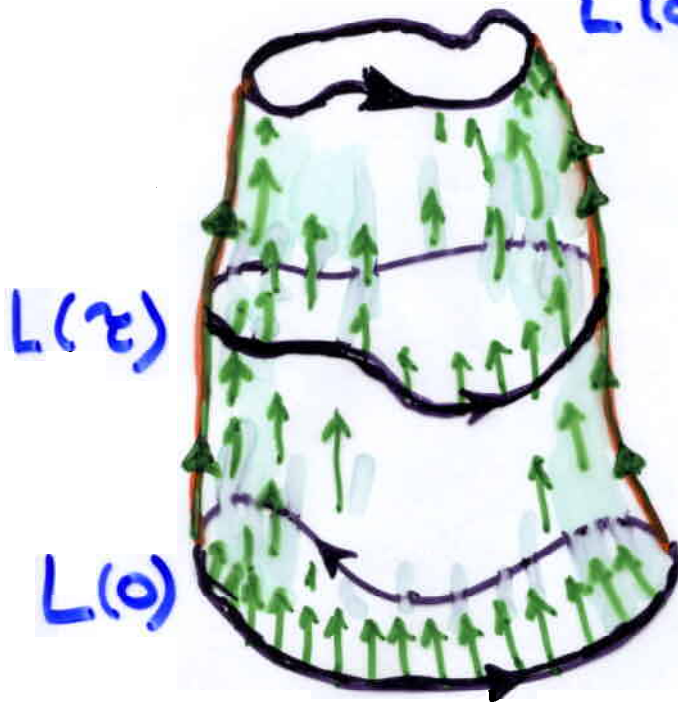
$$x(t) \in p$$

$$x \in \mathbb{R}^{2n}$$

velocity: $\frac{dx}{dt} = v(x)$,
 $t = \text{time}$

$$x(t) = x(t + T_p)$$

$$L(\infty) = p$$

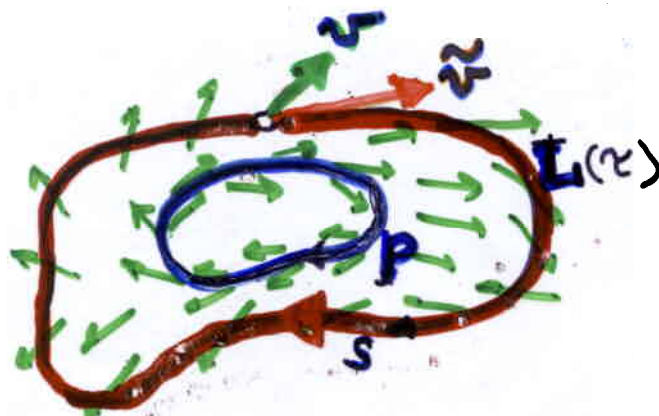


τ -flow in
loop space

$$x(s) \in \mathbb{L} \quad \text{smooth}$$

tangent: $\frac{dx}{ds} = \tilde{v}(x)$, $x(s) = x(s + 2\pi)$

$s = \text{arbitrary loop parameter}$



$v(x(s, \tau)) = \text{velocity}$
 $\tilde{v}(x(s, \tau)) = \text{tangent}$
 $x(s, \tau) \in L(\tau) = \text{loop}$
 $p = \text{periodic orbit}$

$$F^2(\tau) = \oint_L ds (\tilde{v} - \lambda v)^2$$

$$\lambda = |\tilde{v}|/|v|$$

τ -fictitious time flow:

$$\frac{d}{d\tau} (\tilde{v} - \lambda v) = -(\tilde{v} - \lambda v)$$

“Newton descent”

$$\frac{\partial^2 x}{\partial s \partial \tau} - \lambda A \cdot \frac{\partial x}{\partial \tau} - \frac{\partial \lambda}{\partial \tau} v = \lambda v - \tilde{v}$$

drives the initial loop $L_1(0) \rightarrow p$

with $(\tilde{v} - \lambda v)|_{\tau} = e^{-\tau} (\tilde{v} - \lambda v)|_{\tau=0}$

robust p.o. search

✓ SOLVED TURBULENCE

what do I do next?

- Hamiltonian dynamics
in ∞ dimensions
(mhm...)

- "Significant corrections"

$$e^{\Gamma[\phi] + \phi \cdot J} = \sum_{sc} \frac{e^{-\frac{i}{\hbar} S_{sc} + \frac{i\pi}{4} n + \phi \cdot J + \hbar \Gamma_{sc}^{(1)} + \hbar^2 \Gamma_{sc}^{(2)} + \dots}}{|\det S''_{sc}|^{1/2}}$$

- 1-d playpen ($x \rightarrow x^4 + \dots$)

- method 1) Feynman

$$+ \text{diagram 1} + \text{diagram 2} + \dots$$

no transl. invariance

method 2)

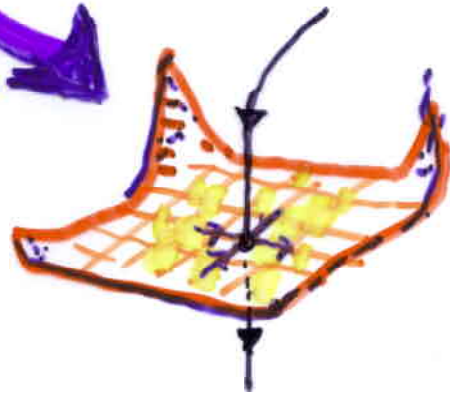
Poincaré meets Feynman

$$\int [d\phi] e^{S[\phi]}$$



$$\phi = h(\tilde{\phi})$$

$$\int [d\tilde{\phi}] \left\| \frac{\partial \phi}{\partial \tilde{\phi}} \right\| e^{-\frac{i}{\hbar} \tilde{\phi} \Delta \tilde{\phi}}$$



penalty

"trivialized"

not

Feynman graphs

method 3)

Transfer operators

$$\mathcal{L} \rightarrow L_{ab} : \hbar^{56} (!)$$

"somewhat classical matters"

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