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Semi-classical Methods in Superconductivity.

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(after Helffer-Moscane...)

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The question in Superconductivity

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Analyze the minima or local minima of the Ginzburg-Landau functional.

$$G(\Psi, A) = \int_{\Omega} \{ |(\nabla - i\kappa A)\Psi|^2 dx$$

$$+ \int_{\Omega} \frac{\kappa^2}{2} (|\Psi|^2 - 1)^2 dx$$

$$+ \kappa^2 \int_{\Omega} |\operatorname{curl} A - \mathcal{H}|^2 dx$$

"Parameters"

Ω

Bounded open set in \mathbb{R}^n ($n=2, 3$)

κ

Ginzburg-Landau parameter

\mathcal{H}

External magnetic field
(permits to vary the intensity)

Normal state:

$$\psi = 0 \quad ; \quad A = \sigma A_e$$

$$\text{curl } A_e = \partial \mathcal{L}$$

Fact:

For σ large, the normal state is a global minimizer

(Gi-Th)

Goal When decreasing σ , find the critical value σ_c for which the normal solution is no more a global (local) minimizer.

Analyze the behavior of the minimizer, for σ near σ_c

Asymptotic regime: κ large

Fact: $\sigma_c \sim \kappa$

Natural step.

Analyze the Hessian of the
G-L functional at $(0, \sigma A_e)$

For σ large, the Hessian is definite
positive (module a gauge degeneracy)

$$\mu^{(1)}(\kappa \sigma_e A_e) \approx \kappa^2$$

where $\mu^{(1)}(A)$ is the lowest
eigenvalue of the Neumann
realization of

$$(\nabla - iA)^2 \text{ in } \Omega$$

Semi-classical parameter:

$$\hbar = 1/\kappa\sigma$$

Semi-classical question

Analyze as $\hbar \rightarrow 0$,
the ground-state energy of

$$\sum_{j=1}^n (\hbar D_{x_j} - A_j)^2$$

(Neumann realization in \mathbb{R})

and the localization of
the corresponding
ground state.

This is at least a first step
for understanding where the
minimizer is no more vanishing,
when bifurcating of the normal
solution (for decreasing σ) ($\sigma \lesssim \sigma_c$)

References (in Mathematics) ⁶

$$(\hbar D - A)^2 + V$$

Helffer - Sjöstrand 1984 - 87

B. Simon *in connexion with Solid State Physics*

Magnetic bottles

1988 - 89

Montgomery, 1995

Helffer - Mohamed 1996

Can't hear

The zero locus of a magnetic field

Superconductivity 1972 Baumann-Phillips - Tang 1998

Bernoff - Sternberg 1998

Gingari - Phillips 1999

Delpino - Fellner - Steinberg 2000

Lu - Pan 1999 - 2001

Helffer - Monneau 2001, 2002

Helffer - Pan 2002

Bonnaillie 2002 Pan

2001

Main Results $n=2$

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Variable magnetic fields: $B(x) = \text{curl } A$

$$b = \inf_{x \in \Omega} |B(x)|$$

$$b' = \inf_{x \in \partial\Omega} |B(x)|$$

Th 1 (Lu-Pan) $n=2$

$$\lim_{h \rightarrow 0} \frac{\lambda^1(h)}{h} = b \quad (\text{Dirichlet})$$

$$\lim_{h \rightarrow 0} \frac{\mu^1(h)}{h} = \min[b, \delta_0 b'] \quad (\text{Neumann})$$

$$0 < \delta_0 < 1$$

For Neumann:

- Localization at the points where the infimum is attained

- $B(x) = \text{cte}$
↳ localization at the boundary.

When

k is large,

Superconductivity appears

first at the

boundary.

The decay is like

$$\exp - \frac{d[x, 2\Omega]}{\sqrt{2}}$$

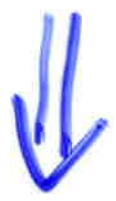
$$d > 0$$

Exponential decay \therefore proof uses

Agmon estimates, like in He-Sj.

Conjecture

$\alpha \rightarrow \theta(\alpha)$ is monotonic



The ground state is localized near the corners of minimal angle.



localization

(2005)

Only proved by Bonnaillie in the form

The ground state is localized near the corners for which $\theta(\alpha)$ is minimal.

Th2 ([DFS] + [BS] + [LP] + [HeMo])

Ass Constant magnetic field - $b \neq 0$

$$\mu^1(h) = b \Theta_0 h + c_{\uparrow} \max_{x \in \mathbb{R}} K(x) h^{3/2} + O(h^{3/2+2})$$

NEUMANN!

$h > 0$

Any associated ground state is localized at the points of maximal curvature.

"Corners"

(²⁰⁰² Bonneauille, ²⁰⁰¹ Jafarulluh, ²⁰⁰¹ Pen)

$$\lim_{h \rightarrow 0} \frac{\mu^1(h)}{h} = \min_{d_j} b \Theta(d_j)$$

$$0 < \Theta(d_j) \leq \Theta_0$$

$$\lim_{L \rightarrow 0} \Theta(L) = 0$$

angle of the corner

conjectured to be $< \Theta_0$

Main Results $n=3$ | (Hello, Pan?) III

Theorem 3 [Pan]

If the magnetic field is constant, then

i) $\lim_{h \rightarrow 0} \frac{\psi^{(h)}(x)}{h} = \mu_0 B$

ii) Any ground state is localized at the boundary [Pan]

iii) Any ground state is localized at the points of $\partial\Omega$ where $\partial\Omega$ is denoted by Γ_B

$B(x) = \text{curl } A(x)$ is tangent to $\partial\Omega$.

Ex:

$$\alpha x_1^2 + \beta x_2^2 + \gamma x_3^2 \leq 1$$

$$\alpha x_1 B_1 + \beta x_2 B_2 + \gamma x_3 B_3 = c$$

Curvature effects in dimension 3. He-Mo Pa?

Invariants_

$$\Gamma_B = \{x \in \Omega ; \vec{B}_0 \cdot \vec{n}_x = 0\}$$

Assume Γ_B is regular,

and \vec{B} not tangent to Γ_B

↑ normal at x

$$K_{m,B}(x) = K_x(T(x) \wedge \vec{n}_x, \frac{\vec{B}}{|\vec{B}|})$$

↑ second fundamental form

Effective potential on Γ_B :

$$\tilde{\gamma}_0(x) = \left(\frac{1}{2}\right)^{2/3} \hat{\gamma}_0 \delta_0^{1/3} \left(|K_{m,B}(x)|^{2/3} (\delta_0 + (1-\delta_0) |\langle T(x), \frac{\vec{B}}{|\vec{B}|} \rangle|^{2/3}) \right)$$

$$\hat{\gamma}_0 = \inf_{x \in \Gamma_B} \tilde{\gamma}_0(x), \quad \delta_0 \in]0, 1[$$

Theorem 4

$$\begin{aligned} \mu^{\pm}(h) = & \theta_0 h \\ & + \hat{\gamma}_0 \theta^{2/3} h^{4/3} \\ & + O(h^{7+4/3}) \end{aligned}$$

$h > 0$

Expected:

Localization near the points of V_B

where $\hat{\gamma}_0 = \tilde{\gamma}_0(\alpha)$

Unexplained constants

$\theta_0, \theta(\kappa), \hat{V}_0, \delta_0$

Sketch of the proof through analysis of models.

Analysis of the models
n=2

Exercise 1

$$D_t^2 + t^2$$

harmonic oscillator

Application

$$D_{x_1}^2 + (D_{x_2} - x_1)^2$$

bottom 1

Exercise 2

$$D_t^2 + (t-\xi)^2$$

harmonic oscillator
 on \mathbb{R}^+ , Neumann

$$\xi \rightarrow \mu(\xi)$$

$$\inf_{\xi} \mu(\xi) = \mu_0 = \mu(\xi_0)$$

unique

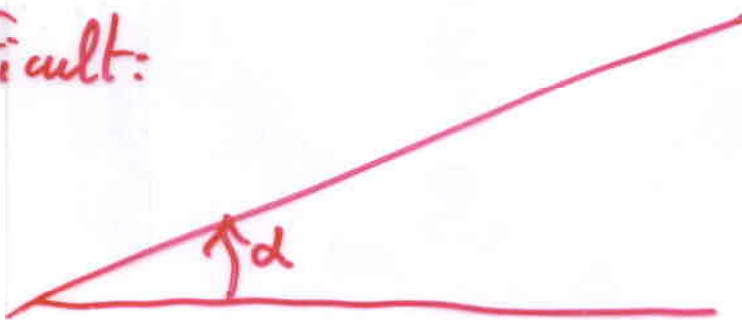
$$0 < \delta_0 = \frac{1}{2} \mu''(\xi_0) < 1$$

Application

$$D_{x_1}^2 + (D_{x_2} - x_1)^2$$

on $\mathbb{R}^+ \times \mathbb{R}$, Neumann
 bottom μ_0

More difficult:
 Exercise 3



Bottom
 of essential
 spectrum
 is μ_0

$\mu(\alpha)$
 bottom
 of
 the spec-
 trum



positive curvature

$$\textcircled{+} \delta - \frac{2\pi_3}{R} b^{\frac{1}{2}} + O(a)$$

& large



negative curvature

$$\textcircled{+} \delta + \frac{2\pi_c}{r} b^{\frac{1}{2}} + O(a)$$

& large

Exercise 4

$$D_{x_1}^2 + D_{x_2}^2 + \left(\tau - c_3 \partial_{x_1} - \sin \partial_{x_2} \right)^2; \quad \tau > 0$$

+ Neumann at $x_1 = 0$

$$\sigma(\theta) = \inf_{\tau} \sigma(\theta, \tau)$$

$$\sigma(0) = 0; \quad \sigma\left(\frac{\pi}{2}\right) = 1$$

$\sigma \nearrow$ on $\left[0, \frac{\pi}{2}\right]$

Application

$$D_{x_1}^2 + D_{x_2}^2 + \left(D_{x_3} - c_3 \partial_{x_1} - \sin \partial_{x_2} \right)^2$$

The bottom of the spectrum is minimal
for $\theta = 0$

Curvature model.

$$\begin{aligned}
 & \left(\hbar D_{\tilde{r}} - \sin \theta t \right)^2 \\
 & + \left(\hbar D_S + \cos \theta t + \kappa \frac{\tilde{r}^2}{2} \right)^2 \\
 & + \hbar^2 D_t^2
 \end{aligned}$$

angle between T and $\frac{\vec{B}}{|\vec{B}|}$
 parametrization of \vec{r}_B
 curvature $\kappa_{m,B}(z)$
 normal variable to Ω

$$\begin{aligned}
 & \rho \left(\hbar^{1/6} \sin \theta D_{\tilde{r}} + \cos \theta \left(\sigma + \frac{\kappa}{2} \tilde{r}^2 \right) \right) \\
 & + \left(\hbar^{1/6} \cos \theta D_{\tilde{r}} - \sin \theta \left(\sigma + \frac{\kappa}{2} \tilde{r}^2 \right) \right)
 \end{aligned}$$

$$\begin{aligned}
 & \textcircled{H}_0 + \hbar^{1/3} \left[\delta_0 \left[\sin \theta D_{\tilde{r}} + \cos \theta \left(\sigma + \frac{\kappa}{2} \tilde{r}^2 \right) \right]^2 \right. \\
 & \quad \left. + \left(\cos \theta D_{\tilde{r}} - \sin \theta \left(\sigma + \frac{\kappa}{2} \tilde{r}^2 \right) \right)^2 \right]
 \end{aligned}$$