

Upper bounds on the number of the scattering poles

PLAMEN STEFANOV

Purdue University

Classical Resonances

Let P be a compactly supported perturbation of the Laplacian, $n \geq 2$, **odd**, such that $P = -\Delta$ outside $B(0, R)$. Assume the “black-box” conditions.

Typical example:

$$P = -\Delta_g + V(x)$$

with Dirichlet or Neumann B.C. on the boundary of the obstacle $\mathcal{O} \subset B(0, R)$ (can be empty).

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Classical Resonances: Poles of the meromorphic continuation of the resolvent

$$\chi(P - \lambda^2)^{-1}\chi \quad \text{from } \Im\lambda > 0 \text{ to } \mathbf{C},$$

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Semi-classical Resonances

$P(h) = -h^2\Delta$ outside $B(0, R)$. **S.C. Resonances:** poles of

$$\chi(P(h) - z)^{-1}\chi$$

from $\Im z > 0$ to a neighborhood of $E_0 > 0$ in \mathbf{C} . Scattering matrix $S(z)$.

Reference Operator $P^\#$ (or $P^\#(h)$) on the “perturbed torus” containing $B(0, R)$.

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$$N(r) = \{\lambda - \text{resonance}, |\lambda| \leq r\}.$$

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Melrose, Sjöstrand, Zworski, Vodev, etc.:

$$\begin{aligned} N(r) &\leq C_n(r+1)^n, & \text{if } n^\# = n, \\ N(r) &\leq C\Phi(Cr), & \text{otherwise,} \end{aligned}$$

where $N^\#(r) \leq \Phi(r)$.

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In S.C. case: $N(\Omega) = O(h^{-n^\#})$, where Ω is a small neighborhood of $E_0 > 0$.

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- Sharp estimates in various neighborhoods of \mathbf{R} ?

$O(|\lambda|^{-N})$ (respectively $O(h^N)$) close to the real axis

If $P = -\Delta_g + V$ in \mathbf{R}^n , then for some $N \gg 1$, [S]

$$\#\text{Res}(P) \cap \{1 \leq |\lambda|; -\Im\lambda \leq C|\lambda|^{-N}\} \leq C_n r^n (1 + o(1))$$

with

$$C_n = 2(2\pi)^{-n} \text{vol}(\mathcal{T}),$$

\mathcal{T} being the *trapped set* in $\{p_0(x, \xi) \leq 1\}$. In a presence of an obstacle, $N = \infty$.

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In the **S.C. case**,

$$\#\text{Res}(P(h)) \cap \{0 < a \leq \Re z \leq b; 0 \leq -\Im z \leq Ch^{-N}\} \leq C_n h^{-n} (1 + o(1)),$$

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Those estimates are sharp in some cases (potential barrier).

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“Close” to \mathbf{R} , we have (semi-classical max principle)

$$\frac{1}{\text{dist}(\lambda^2, \text{spec}(P_\theta))} \leq \|(P_\theta - \lambda^2)^{-1}\| \leq \frac{|\lambda|^M}{\text{dist}(\lambda^2, \text{spec}(P_\theta))}$$

with $M = M(n) \dots$

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with $M = M(n)$... but there are still difficulties connected to the fact that the resonant states are not asymptotically orthogonal, may cluster, etc.

$O(1)$ (respectively $O(h)$) close to the real axis

SJÖSTRAND: hyperbolic flow, additional assumptions:

$$\text{Res}(P(h)) \cap \left\{ [E_1, E_2] + i[-\delta, 0] \right\} \leq Ch^{-n} \delta^{\mu/2-\epsilon},$$

where $Ch \leq \delta \leq 1/C$ and μ is the Minkowski codimension of \mathcal{T} . If $\delta \sim h$, we have

$$O(h^{-d/2-\epsilon})$$

resonances with $d = n - \mu$ (the Minkowski dimension of \mathcal{T}).

Resonances in sectors $0 \leq -\arg \lambda \ll 1$

SJÖSTRAND, ZWORSKI: Define $P_\epsilon^\#$ in an ϵ -neighborhood K_ϵ of the convex hull of the “black box” with Dirichlet B.C. Then

$$\#\text{Res}(P) \cap \{-\theta < \arg \lambda \leq 0\} \leq 2(1 + C\epsilon)N_\epsilon^\#(r) + C\epsilon r^n,$$

where $\epsilon = \theta^{2/7}$ and $N_\epsilon^\#(r)$ satisfies reasonable conditions.

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ZERZERI: K_ϵ might be an ϵ -neighborhood of the “black box” (no convexity) but then we have to add $2(2\pi)^{-n} \times$ the measure of the “trapped rays” $\times r^n$ (not the same as \mathcal{T} above).

Resonances in larger sectors $0 \leq -\arg \lambda < \pi/2 - \delta$ ($-\Im z = O(1)$)

SJÖSTRAND, PETKOV-ZWORSKI: “bottle theorem”

$$N^\#((1 - \epsilon)r) - E_-(r) \leq N_\delta(r) \leq N^\#((1 + \epsilon)r) + E_+(r),$$

where

$$0 \leq E_\pm(r) \leq C_1 \epsilon r^{n^\#} + C_2(R_0, \epsilon) r^n + \dots$$

$C_{1,2}$ independent of P ! In particular, if P depends on a parameter so that $N^\#(r)/r^n \gg 1$ (and $n^\# = n$), then we have almost an asymptotic formula! If $n^\# > n$, i.e., if $N^\#(r)$ dominates, (hypoelliptic operators), then we have true asymptotics (VODEV, SJÖSTRAND, PETKOV-ZWORSKI).

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Seimi-classical version: PETKOV-ZWORSKI

Resonances in a disk $|\lambda| \leq r$

$$P = -\Delta + V(r), \quad \text{supp } V \subset [0, R]$$

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$$C_n = ?$$

Black box in $B(0, R)$, $n^\# = n$ odd. Define the continuous functions

$$M(r) = \int_0^r \frac{N(t)}{t} dt, \quad M^\#(r) = \int_0^r \frac{N^\#(t)}{t} dt.$$

Assume also that $N^\#(r+1) - N^\#(r) = O(r^{n-\epsilon})$, with some $\epsilon \in (0, 1]$.

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Theorem 1 *There exists $C_n > 0$ depending on the dimension only, such that*

$$M(r) \leq 2M^\#(r) + C_n R^n r^n + o(r^n). \quad (1)$$

As a consequence, for any $0 < \epsilon \leq 1$,

$$N(r) \leq 2(1 + \epsilon)N^\#((1 + \epsilon)r) + C_n R^n r^n / \epsilon + o_\epsilon(r^n). \quad (2)$$

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If we know more about the regularity of $N(r)$, then (1) is true for $N(r)$ and $N^\#(r)$.

The proof follows R. FROESE (where $P = -\Delta + V$), and PETKOV-ZWORSKI (semiclassical version). Denote $s(\lambda) =$ scattering matrix, $\sigma(\lambda) =$ scattering phase.

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Lemma 1 *For any $r > 0$ we have*

$$M(r) = 2 \int_0^r \frac{\sigma(t)}{t} dt + \frac{1}{2\pi} \int_0^\pi \log |s(re^{i\theta})| d\theta.$$

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Proof: Integrate along the semi-disk $[-r, r] \cup r \exp(i[0, \pi])$ using $s'/s = 2\pi i\sigma'$.

$$\begin{aligned} N(t) &= \frac{1}{2\pi i} \oint \frac{s'(z)}{s(z)} dz = \Im \frac{1}{2\pi} \oint \frac{s'(z)}{s(z)} dz \\ &= \int_{-t}^t \sigma'(z) dz + \frac{1}{2\pi} \int_0^\pi t \frac{d}{dt} \log |s(te^{i\theta})| d\theta \\ &= 2\sigma(t) + \frac{1}{2\pi} \int_0^\pi t \frac{d}{dt} \log |s(te^{i\theta})| d\theta. \end{aligned}$$

Divide by t and integrate to get the lemma.

To estimate the scattering phase, we apply results by T. CHRISTIANSEN to get

$$M(r) = 2 \left[M^\#(r) - \tau_n \frac{r^n}{n} \right] + \frac{1}{2\pi} \int_0^\pi \log |s(re^{i\theta})| d\theta + O(r^{n-\epsilon}),$$

with $\tau_n = (2\pi)^{-n} \text{Vol}(\mathbf{T}_R \times B(0, 1))$.

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To estimate the scattering determinant $s(\lambda)$, we proceed in the usual way. The scattering operator $S(\lambda)$ is given by

$$S(\lambda) = I + c_n \lambda^{n-2} \mathbf{E}_-(\lambda) [\Delta, \chi_2] R(\lambda) [\Delta, \chi_1]^t \mathbf{E}_+(\lambda), \quad (3)$$

where

$$[\mathbf{E}_\pm(\lambda)f](\omega) = \int e^{\pm i\lambda\omega \cdot x} f(x) dx = \hat{f}(\mp\lambda\omega), \quad \omega \in S^{n-1},$$

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We estimate the characteristic values of the scattering amplitude which reduces to an estimate of

$$\Delta_\theta^m e^{i\lambda x \cdot \theta}, \quad |x| \leq R.$$

We need to work here in a sector $0 < \delta \leq \arg \lambda \leq \pi - \delta$. Using a standard argument, to cover the missing sectors $0 \leq \arg \lambda \leq \delta$ and $\pi - \delta \leq \arg \lambda \leq \pi$, we use the fact that $|s| = 1$ on \mathbf{R} and the Phragmen-Lindelöff principle.

Modified estimate:

$$M(r) \leq 2 \left[M^\#(r) - \tau_n \frac{r^n}{n} \right] + \frac{1}{n} C_n R^n r^n + o(r^n).$$

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Here

$$M^\#(r) - \tau_n \frac{r^n}{n} = \begin{cases} 0, & \text{if } P = -\Delta + V, \\ -\frac{1}{n} (2\pi)^{-n} \int_{x \in \mathcal{O}, |\xi| \leq 1} dx d\xi & \text{in obstacle scattering,} \\ \frac{1}{n} (2\pi)^{-n} \int_{|\xi| \leq 1} (|g|^{1/2} - 1) dx d\xi & \text{for } P = -\Delta_g. \end{cases}$$

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$$nM(r) \leq 2nM^\#(r) + C'_n R^n r^n + o(r^n),$$

where $C'_n = C_n - (2\pi)^{-n} \text{vol}(B(0, 1))^2$, and

$$nM^\#(r) = \begin{cases} (2\pi)^{-n} \int_{|x| \leq R, |\xi| \leq 1} dx d\xi, & \text{if } P = -\Delta + V, \\ (2\pi)^{-n} \int_{x \in B(0, R) \setminus \mathcal{O}, |\xi| \leq 1} dx d\xi & \text{in obstacle scattering,} \\ (2\pi)^{-n} \int_{|x| \leq R, |\xi| \leq 1} |g|^{1/2} dx d\xi & \text{for } P = -\Delta_g. \end{cases}$$

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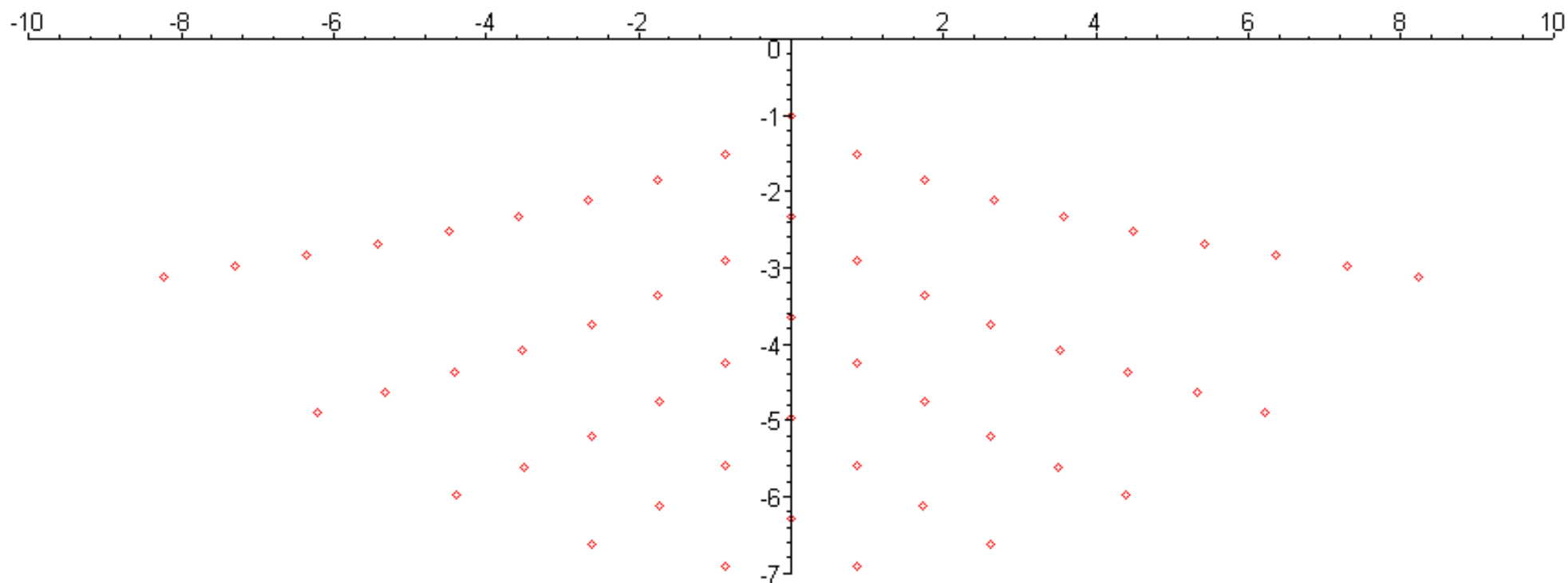
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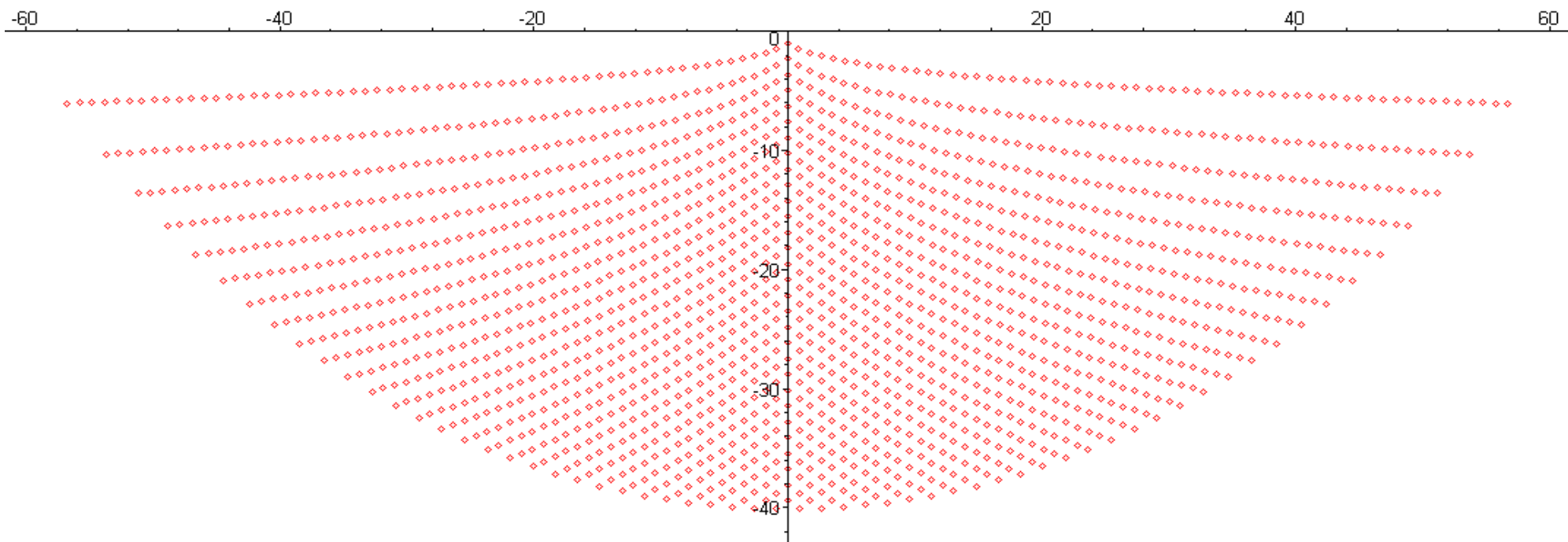
$$C'_n = ?$$

Resonances of the unit sphere, Dirichlet B.C.

$n = 3$, resonances are the zeros of $H_{k+1/2}^{(1)}(\lambda)$, $k = 0, 1, \dots$ Multiplicity = $2k + 1$.



More of them:



OLVER: Asymptotically, the sphere resonances are given by

$$\lambda_{k,s} \sim \nu \zeta^{-1} \left(-\nu^{-\frac{2}{3}} e^{\frac{i\pi}{3}} a_s \right), \quad s = 1, 2, \dots, k; \quad k = 1, 2, \dots,$$

where $\nu = k + \frac{1}{2}$, and $\dots < a_2 < a_2 < a_1 < 0$ are the zeros of $\text{Ai}(s)$, and

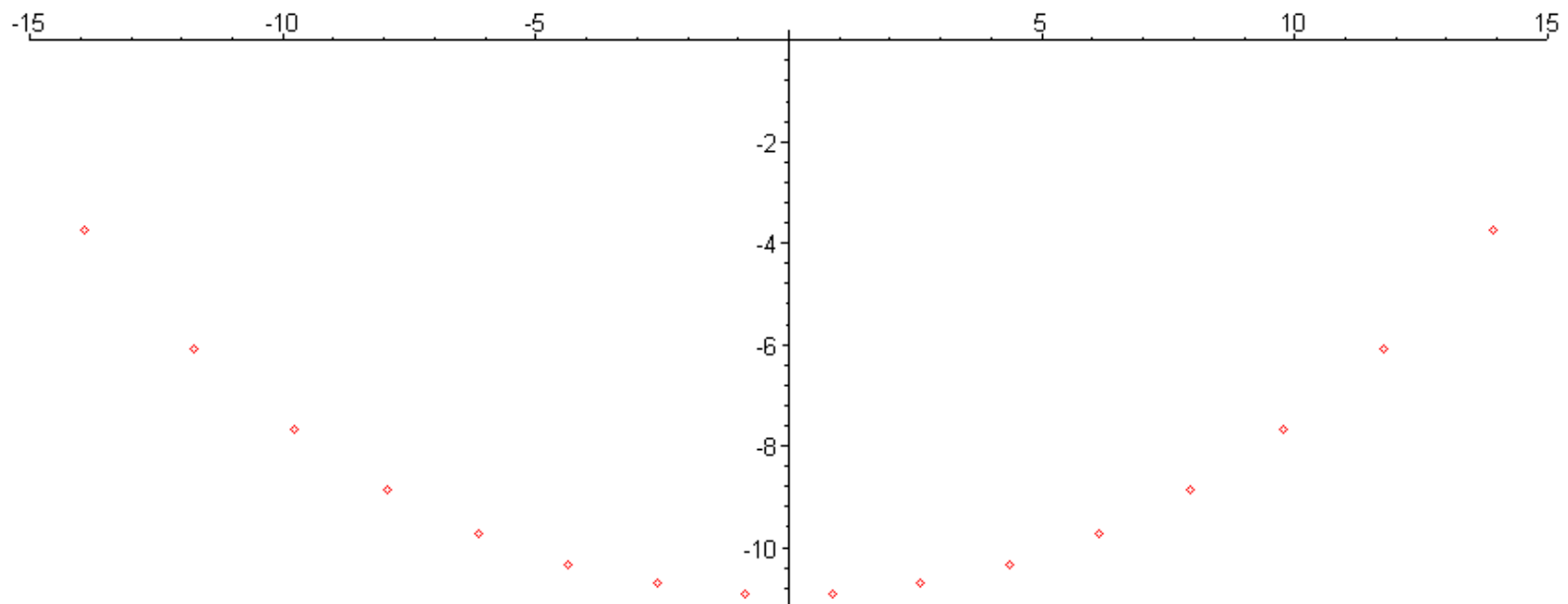
$$\frac{2}{3} \zeta^{\frac{3}{2}}(z) = \log \left(\frac{1 + \sqrt{1 - z^2}}{z} \right) - \sqrt{1 - z^2}$$

OLVER: Asymptotically, the sphere resonances are given by

$$\lambda_{k,s} \sim \nu \zeta^{-1} \left(-\nu^{-\frac{2}{3}} e^{\frac{i\pi}{3}} a_s \right), \quad s = 1, 2, \dots, k; \quad k = 1, 2, \dots,$$

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$\lambda_{k,s}$ for $k = 16$

Define $N_{S^2}(r) = \#\text{Res}(P) \cap \{|\lambda| < r\}$. We have

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C_{S^2} can be expressed as a certain integral involving ζ^{-1} .

Resonances of $P = -\Delta + \mathbf{1}_{B(0,R)}$, $n = 3$

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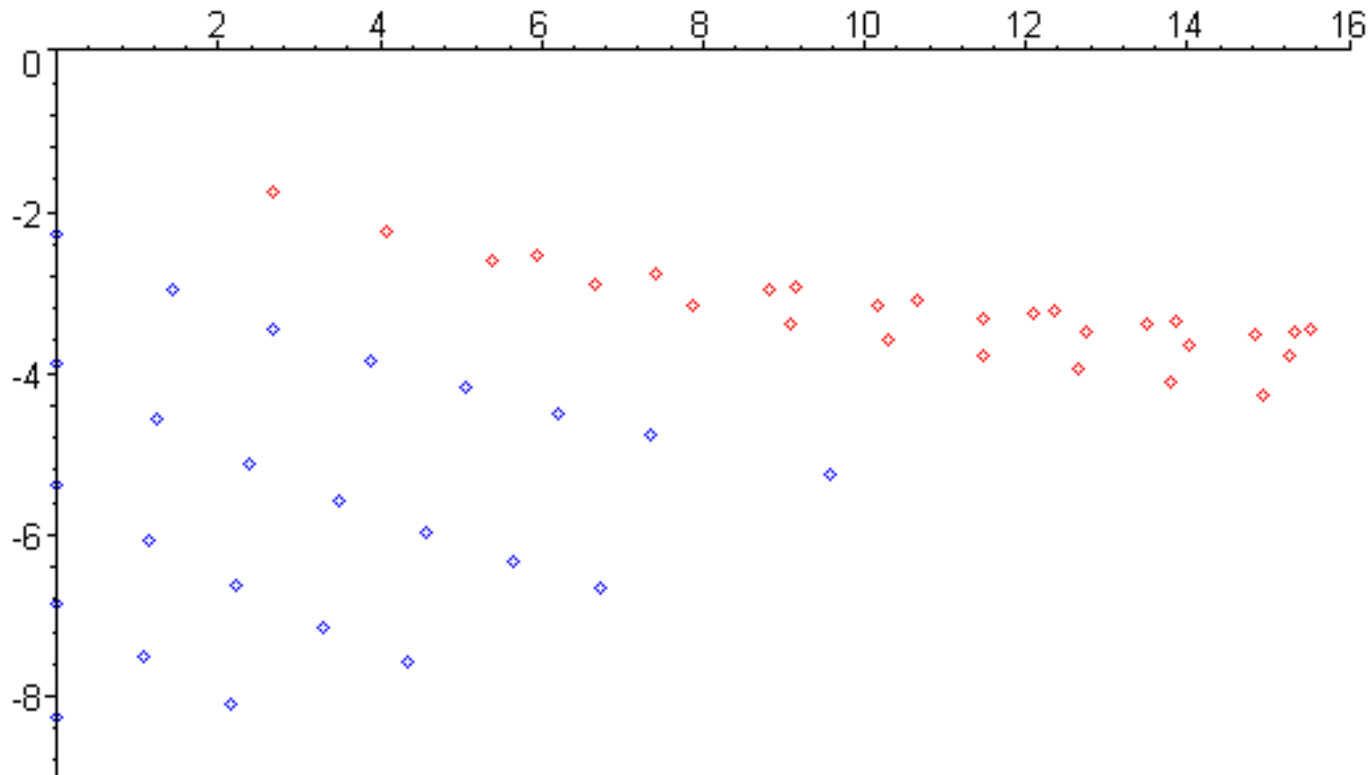
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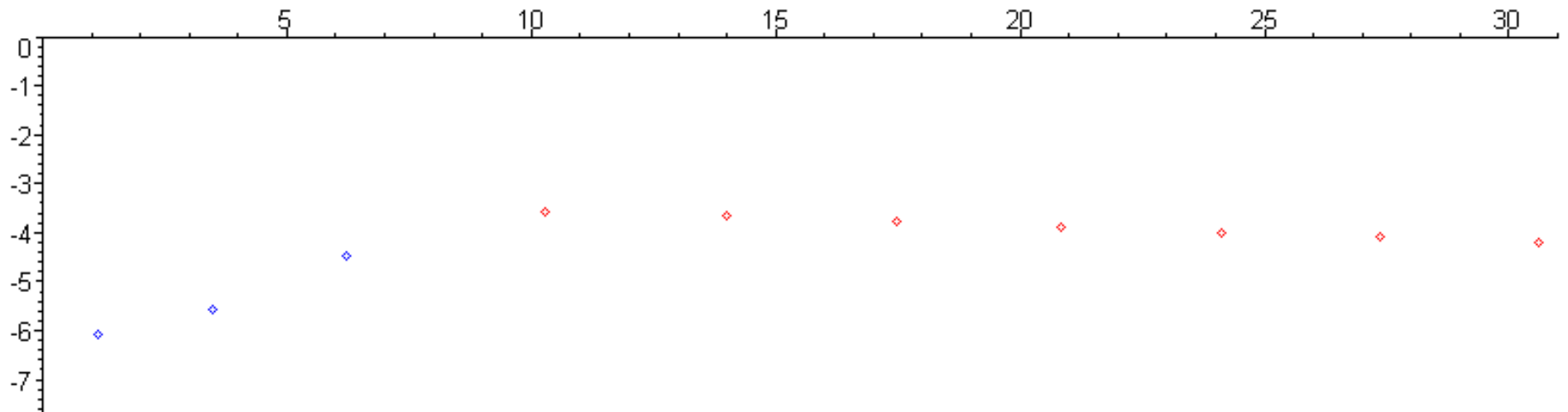
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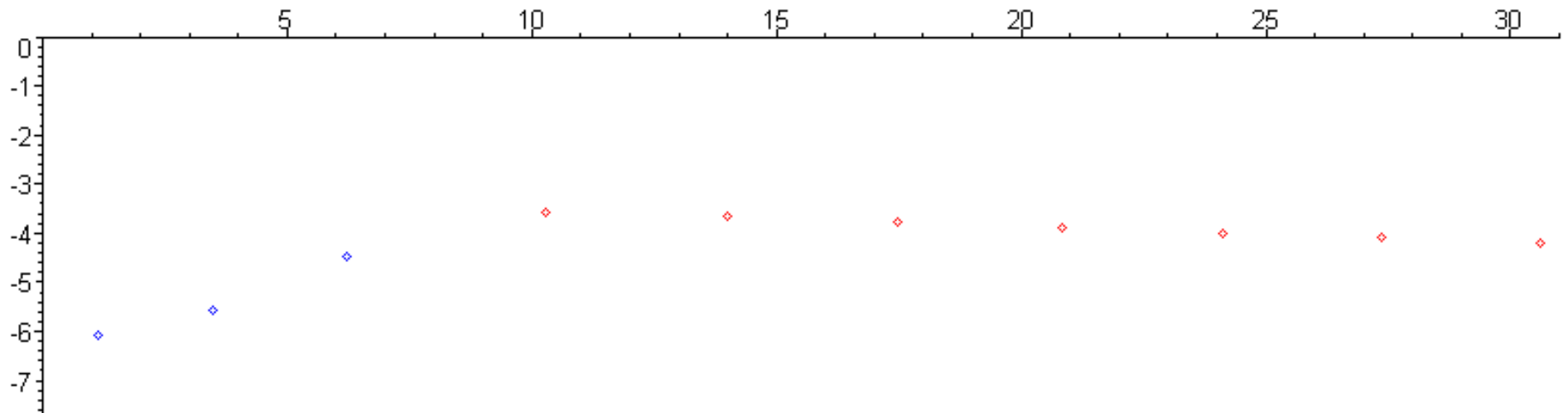
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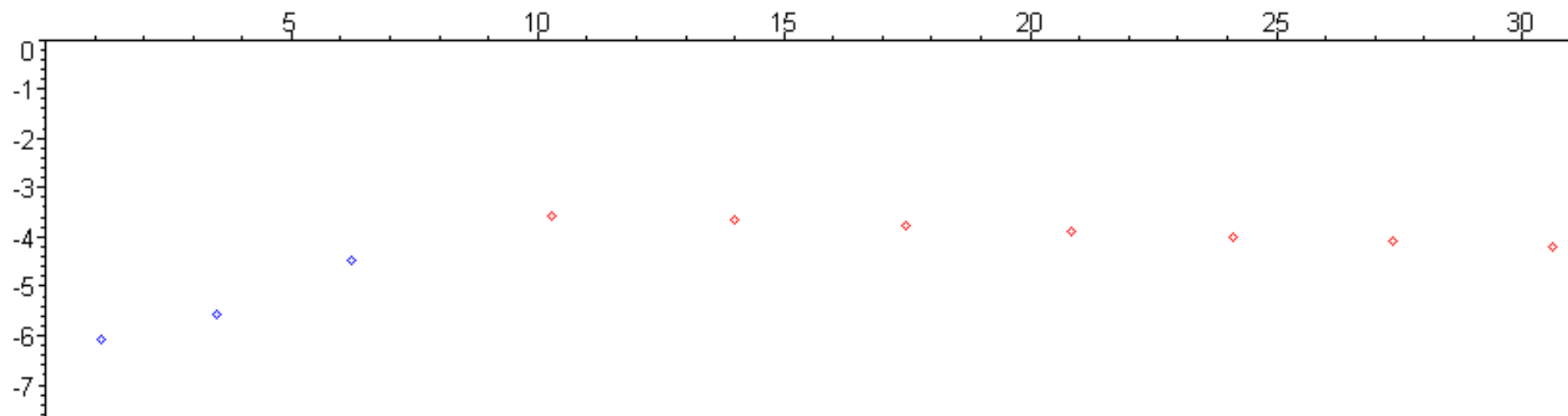
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It appears that

$$N(r) = C_B r^3 + C_{S^2} r^3 + o(r^3),$$

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So, at least in this case,

$$C'_n = C_{S^{n-1}}$$

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Let as above, $P^\#$ be a reference operator equal to P in $B(0, R)$ with Dirichlet B.C. on $\partial B(0, R)$.

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