

Workshop MSRI 5/5-2003

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Perturbations of self-adjoint operators with periodic flow - branching levels.

joint work in progress with M. Hitrik

## 0. Introduction.

Melin-Sj (2001) found large stable classes of analytic non-selfadjoint pseudodiff. op. in dimension with spectrum given by a Bohr-S. condition. (absence of small denominators.)

Current project with Hitrik 2002-20XY

Small non-s.a. perturbations of s.a. operators.

So far, mainly the case when the classical flow of the unperturbed part is periodic. ②

Let  $M = \begin{cases} \mathbb{R}^2 \\ \text{or} \\ \text{Compact 2-dim analytic} \\ \text{m fld.} \end{cases}$

When  $M = \mathbb{R}^2$ , let

$$P_\varepsilon = P^w(x, hD_x, \varepsilon; h)$$

$[0, \varepsilon_0[ \ni \varepsilon \longmapsto P(x, \xi, \varepsilon; h) \in C^\infty$

space of holomorphic function  $f(x, \xi)$  in tubular neighborhood of  $\mathbb{R}^4$  with  $|f(x, \xi)| \leq C m(\text{Re}(x, \xi))$

$$m \geq 1 \quad m(X) \leq C_0 \langle X - Y \rangle^{N_0} m(Y)$$

$$P_\varepsilon(x, \xi, \varepsilon; h) \sim \sum_{j=0}^{\infty} P_{j, \varepsilon}(x, \xi) h^j, \quad h \rightarrow 0 \quad (3)$$

in this space

$$|P_{0, \varepsilon}(x, \xi)| \geq \frac{1}{C} m(x, \xi), \quad |(x, \xi)| \geq C'$$

Write  $P = P_{0, 0}$ ,  $P_\varepsilon = P_{0, \varepsilon}$

When  $M$  compact, assume  $P_\varepsilon$  is a  $h$ -differential operator ....

(Replace  $m(x, \xi)$  by  $\langle \xi \rangle^m$ ,  $m > 0$ .)

Assume  $P_{\varepsilon=0}$  is self-adjoint.

Then near 0,  $\sigma(P_\varepsilon)$  is discrete and contained in  $\{ |Im z| \leq O(\varepsilon) \}$ .

Assume  $\bar{p}^{-1}(0) \cap T^*M$  is connected and  $dp \neq 0$  there.

Let  $H_p = P'_\xi \cdot \frac{\partial}{\partial x} - P'_x \cdot \frac{\partial}{\partial \xi}$ . Assume (4)

The  $H_p$ -flow is periodic on  $\bar{p}^{-1}(E) \cap T^*M$  with period  $T(E) > 0$  analytic in  $E$ ,  $E \in \text{neigh}(0, \mathbb{R})$ .

Write  $P_\varepsilon = P + i\varepsilon q + O(\varepsilon^2 m)$

Following the tradition in the self-adjoint case (A. Weinstein, Colin de Verdière, Guillemin, Boutet de Monvel, Helffer - Robert, Dozias, Ivrii ...) introduce <sup>Grigis</sup>

$$\langle q \rangle = \frac{1}{T(E)} \int_{-T(E)/2}^{T(E)/2} q_0 \exp(t H_p) dt,$$

on  $\bar{p}^{-1}(E)$ .

Using phase space exponential weights  $\approx$  Conjugation by a Fourier int. op. with complex phase



Can assume

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$$P_\varepsilon = P + i\varepsilon \langle q \rangle + \mathcal{O}(\varepsilon^2)$$

on (new) real phase space.

Assume  $\langle q \rangle$  real or more generally

$\text{Im} \langle q \rangle$  is an analytic function of  $P$  and  $\text{Re} \langle q \rangle$ .

Let  $F_0 \in \left[ \min_{\bar{P}'(0)} \langle \text{Re} q \rangle, \max_{\bar{P}'(0)} \langle \text{Re} q \rangle \right]$ .

Goal: determine all eigenvalues in

a rectangle  $\left] -\frac{1}{\mathcal{O}(1)}, \frac{1}{\mathcal{O}(1)} \right] \left[ +i\varepsilon \right] F_0 - \frac{1}{\mathcal{O}(1)}, F_0 + \frac{1}{\mathcal{O}(1)} \left[$

for  $h \ll \varepsilon \leq \mathcal{O}(h^\delta)$  in general, and

for  $h^2 \ll \varepsilon \leq \mathcal{O}(h^\delta)$  when the

subprincipal symbol of  $P_{\varepsilon=0}$  vanishes.

Let

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$$\Lambda_{0, F_0} = \{p \in T^*M; p(p) = 0, \text{Re}\langle q \rangle = F_0\}.$$

Assume

$T(0)$  is the minimal period for the  $H_p$ -flow on  $\Lambda_{0, F_0}$ .

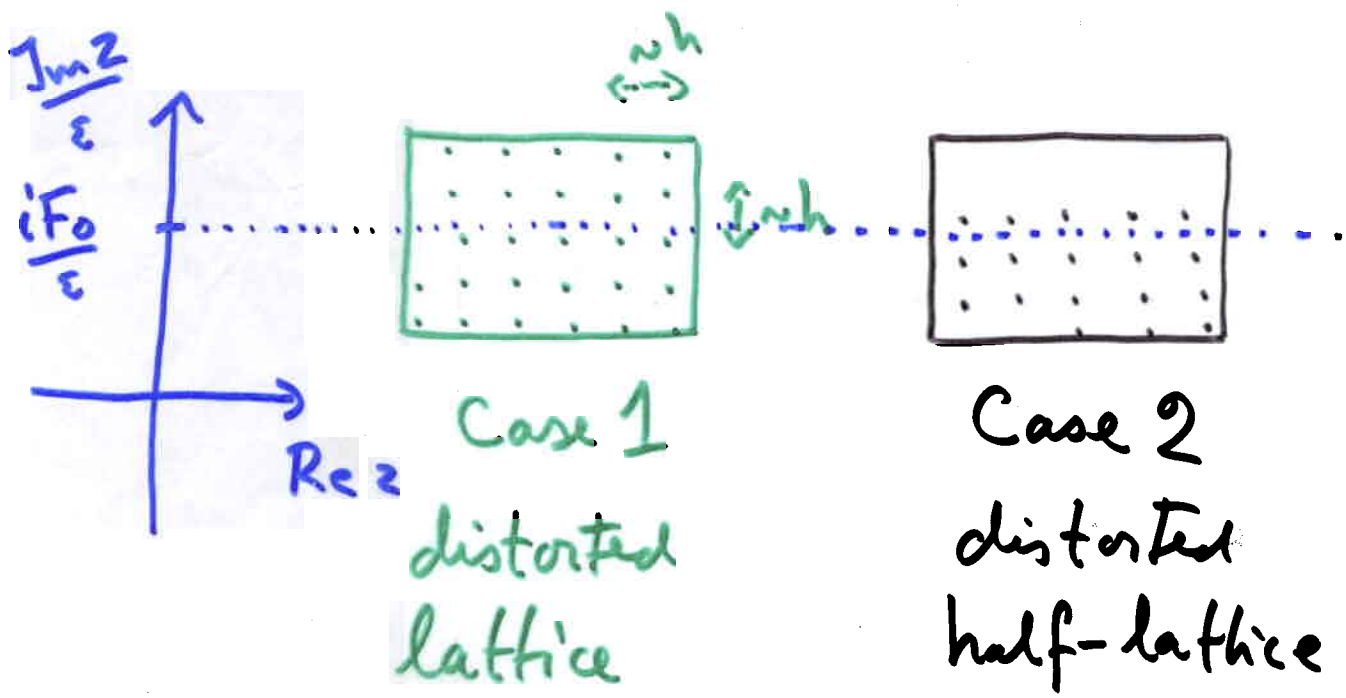
Then  $\bar{p}^{-1}(0)/\exp \mathbb{R} H_p$  is a 2-dim. symplectic mfld near the image of  $\Lambda_{0, F_0}$ .  $\text{Re}\langle q \rangle$  is an analytic function on this mfld.

Case 1.  $F_0$  is not a critical value of  $\text{Re}\langle q \rangle$ .  $\Rightarrow \Lambda_{0, F_0}$  is a Lagrangian torus.

Case 2.  $F_0$  is a non-deg. maximum or minimum, attained at a single point.

# Spectrum (Hitrik-Sj 2003):

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(In case 1 result still valid for  $\epsilon$  up to  $\epsilon_0 > 0$  small but independent of  $h$ . Sj 2002)

Remaining case generically:

**Case 3:**  $F_0$  is a non-degenerate saddle point level, attained at only 1 critical point.

This talk

## 1. Reductions:

$Re\langle q \rangle = F_0 :$



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Conjugation by microlocally defined elliptic Fourier integral operator

$\rightsquigarrow \tilde{P}_\varepsilon$  on  $L^2_{(S_0, S_1, S_2)}(S^1_t \times \mathbb{R}_x)$

a space of microlocally defined Floquet periodic functions with Floquet parameters  $e^{iS_0/h}, e^{iS_1/h}, e^{iS_2/h}$ ,  $S_j =$  action differences.

Leading symbol of  $\tilde{P}_\varepsilon$  :

$\tilde{P}_\varepsilon = g(\tau) + i\varepsilon \langle g \rangle(\tau, x, \tau) + \mathcal{O}(\varepsilon^2)$   
 $\uparrow_{\text{real}, g' \neq 0}$

Further averagings in  $t$  eliminate the  $t$ -dependence from the full symbol and we get :



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a) General case:

$$\tilde{P}_\varepsilon(\tau, x, \xi; h) =$$

$$g(\tau) + \varepsilon [i \langle q \rangle(\tau, x, \xi) + \theta(\varepsilon) + \frac{h}{\varepsilon} P_1 + h \frac{h}{\varepsilon} P_2 + \dots]$$

$\varepsilon, \frac{h}{\varepsilon}$  small parameters if  $h \ll \varepsilon \leq h^\delta$ .

b) The subprincipal symbol of  $P_{\varepsilon=0}$  is 0:

$$\frac{h}{\varepsilon} P_1 = h \tilde{P}_1 : \tilde{P}_\varepsilon(\tau, x, \xi; h) =$$

$$g(\tau) + \varepsilon [i \langle q \rangle(\tau, x, \xi) + \theta(\varepsilon) + \frac{h^2}{\varepsilon} P_2 + h \tilde{P}_1 + h^2 \tilde{P}_2 + \dots]$$

$\varepsilon, \frac{h^2}{\varepsilon}$  small parameters if  $h^2 \ll \varepsilon \leq h^\delta$ .

These reductions are mod  $\mathcal{O}(h^\infty)$  in the space of holomorphic symbols.

To compute the spectrum "it suffices" to expand in Fourier series in  $t$  and study

$$\tilde{P}_\varepsilon(h(k - \frac{k_0}{4}) - \frac{s_0}{2\pi}, x, hD_x; h) \quad k \in \mathbb{Z}$$

$k_0 \in \mathbb{Z}$  is fixed.

Likely that this gives the spectrum in the respective cases  $h \ll \varepsilon \leq h^\delta$ ,  $h^2 \ll \varepsilon \leq h^\delta$ , but we recently discovered a technical difficulty when the lost operator becomes too non-self-adjoint. We then need to work in modified spaces to avoid pseudo-spectral behaviour and these spaces should vary with  $h$ .

(covers the case of the damped wave eq.)

However, in case b) for  $\varepsilon \sim h$  and probably in both cases a, b for  $\varepsilon \leq O(h)$ , we can work on ordinary  $L^2$  and we seem to get complete spectral results. Remainder of the talk: 1D-case.

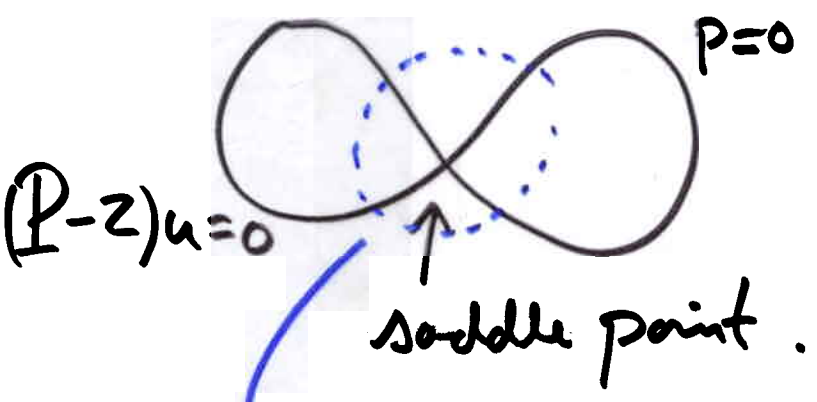
## 2. Spectrum of branching 1-dim. $\textcircled{11}$ pseudodiff. operators.

Let  $P(x, hD_x; h)$  analytic pseudo defined near  $(0,0) \in \mathbb{R}_x \times \mathbb{R}_\xi$ . If  $p$  is real with a saddle point at  $(0,0)$  and  $p(0,0) = 0$  then  $\exists F(\tau; h) = f(\tau) + hf_1(\tau) + \dots$  s.t.

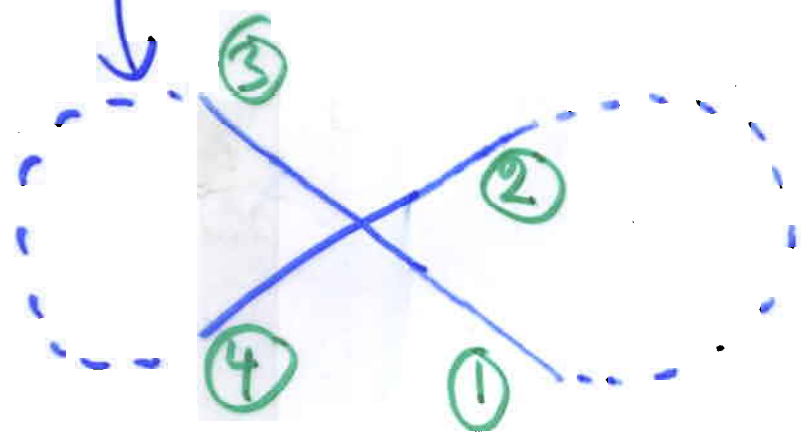
$$U^{-1} F(P; h) U = -P_0$$

$P_0 = \frac{1}{2}((hD_x)^2 - x^2)$ . Still true after making  $p$  slightly complex ( $U$  is then a Fourier int. op. with complex phase) (Helffer-Sj, Horn, C. März, Parisse-Colin de Verdière, Fujise-Ramond, Buslaev-Grigis....)

Let  $p$  be real defined near



and add perturbations making  $P$  non-self-adjoint.



$$(P_0 + \mu)v = 0$$

$$\text{cf } (-\hbar^2 \partial_x^2 + V + \mu) = 0$$



Near  $(j)$ :  $v = v_j \times$  normalized WKB solution

Transition matrix:

$$\begin{pmatrix} v_2 \\ v_1 \end{pmatrix} = \begin{pmatrix} b_{2,3} & b_{2,4} \\ b_{1,3} & b_{1,4} \end{pmatrix} \begin{pmatrix} v_3 \\ v_4 \end{pmatrix}$$



$\det B = 1$   $b_{j,k}$  can be  
computed explicitly

Let  $S_{3,4}^{(M)} =$  extension action from (4) to (3)

$S_{1,2} =$  - " - (2) to (1)

Proposition  $\mu$  is an eigenvalue  
of  $P \Leftrightarrow G(\mu; h) = 0$

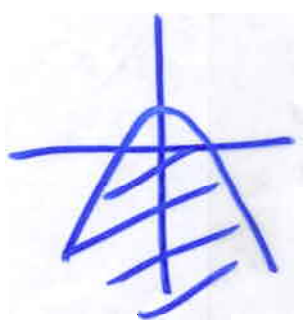
where  $G(\mu; h) = a_1 + a_2 + a_3 + a_4$   
 $= e^{-\frac{\pi M}{2h}} \left( e^{\frac{i}{h}(S_{1,2} + S_{3,4})} b_{2,3} + e^{\frac{i}{h}S_{1,2} + \frac{\pi M}{h}} \right.$   
 $\left. + e^{\frac{i}{h}S_{3,4} + \frac{\pi M}{h}} + b_{1,4} \right)$

Cf: Fujie-Ramond

Davies, Redpath, Shkrikov

If  $\frac{|M|}{h} \gg 1$ ,  $|\arg \mu - \frac{\pi}{2}| < \pi - \frac{1}{\alpha}$ :

Case 1.



$$a_4 = a_{4+} + a_{4-}$$

$$a_{4\pm} = e^{-\mathcal{O}_{2,3}(h) + \mathcal{O}_-(\frac{h}{\mu}) + \frac{i}{h}(-\mu \ln \frac{\mu}{i} - \frac{1-\mu^2}{2})} e^{\pm \frac{\pi M}{h}}$$

$$a_2 = e^{\frac{i}{h} S_{1,2} + \frac{\pi M}{2h}} = e^{\frac{i}{h} \varphi_{4\pm}} = e^{\frac{i}{h} \varphi_2}$$

$$a_3 = e^{\frac{i}{h} S_{3,4} + \frac{\pi M}{2h}} = e^{\frac{i}{h} \varphi_3}$$

$$a_1 = e^{\frac{i}{h}(S_{1,2} + S_{3,4} + \mu \ln \frac{\mu}{i} + \frac{1-\mu^2}{2}) - \mathcal{O}_-(\frac{h}{\mu}) + \mathcal{O}_{2,3}(h)}$$

No zero if one  $a_j$ ,  $j = 1, 2, 3, 4 \pm$  dominates over the others.

Look for the curves  $\Gamma_{j,h}$  where

$|a_j| = |a_k|$  and especially where

$|a_j|, |a_k|$  dominate over the other  $a_j$ .

Can neglect  $a_4^-$  when  $\text{Re } \mu \geq 0$

$a_4^+$  - " -  $\text{Re } \mu \leq 0$

Can neglect  $\Gamma_{2,3}$  since

$$a_2 a_3 = a_1 a_4^+$$

$\text{Re } \mu \geq 0$ :

$$(\text{Im } \mu) \ln \frac{1}{|\mu|} = \begin{cases} \text{Im } S_{3,4}(\mu) + \bar{\Delta}(\mu) : \Gamma_{3,4}^+ = \Gamma_{1,2} \\ \text{Im } S_{1,2}(\mu) + \bar{\Delta}(\mu) : \Gamma_{2,4}^+ = \Gamma_{1,3} \\ \frac{1}{2}(\text{Im } S_{1,2} + \text{Im } S_{3,4}) + \bar{\Delta}(\mu) : \Gamma_{1,4}^+ \end{cases}$$

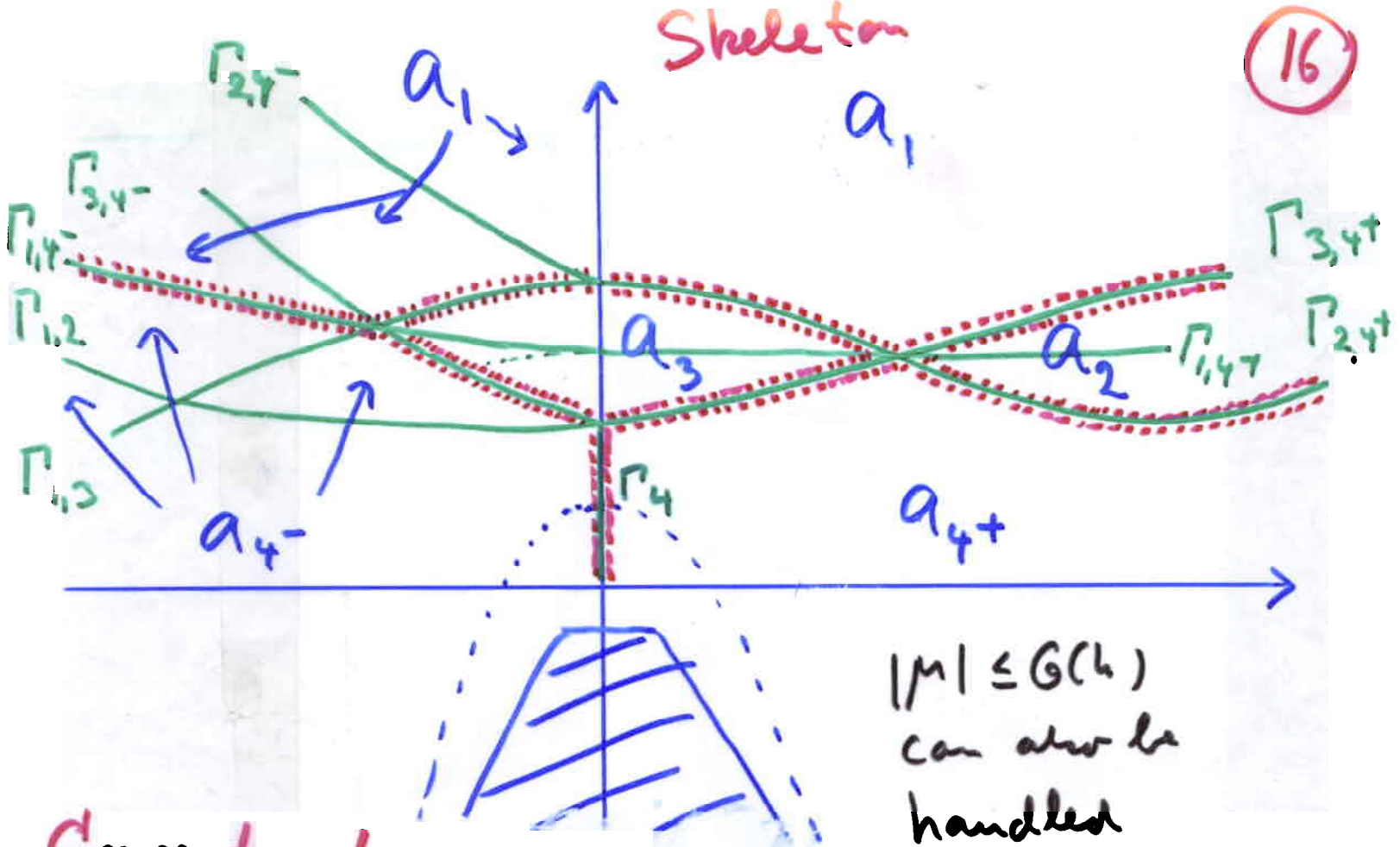
$\text{Re } \mu \leq 0$ :

$$(\text{Im } \mu) \ln \frac{1}{|\mu|} = \begin{cases} \text{Im } S_{3,4} + \bar{\Delta}(\mu) - 2\pi \text{Re } \mu : \Gamma_{3,4}^- \\ \text{Im } S_{1,2} + \bar{\Delta}(\mu) - 2\pi \text{Re } \mu : \Gamma_{2,4}^- \\ \text{Im } S_{1,2} + \bar{\Delta}(\mu) : \Gamma_{1,3} \\ \text{Im } S_{3,4} + \bar{\Delta}(\mu) : \Gamma_{1,2} \\ \frac{1}{2}(\text{Im } S_{1,2} + \text{Im } S_{3,4}) + \bar{\Delta} - \pi \text{Re } \mu : \Gamma_{1,4}^- \end{cases}$$

Exponential concentration of the zeros to

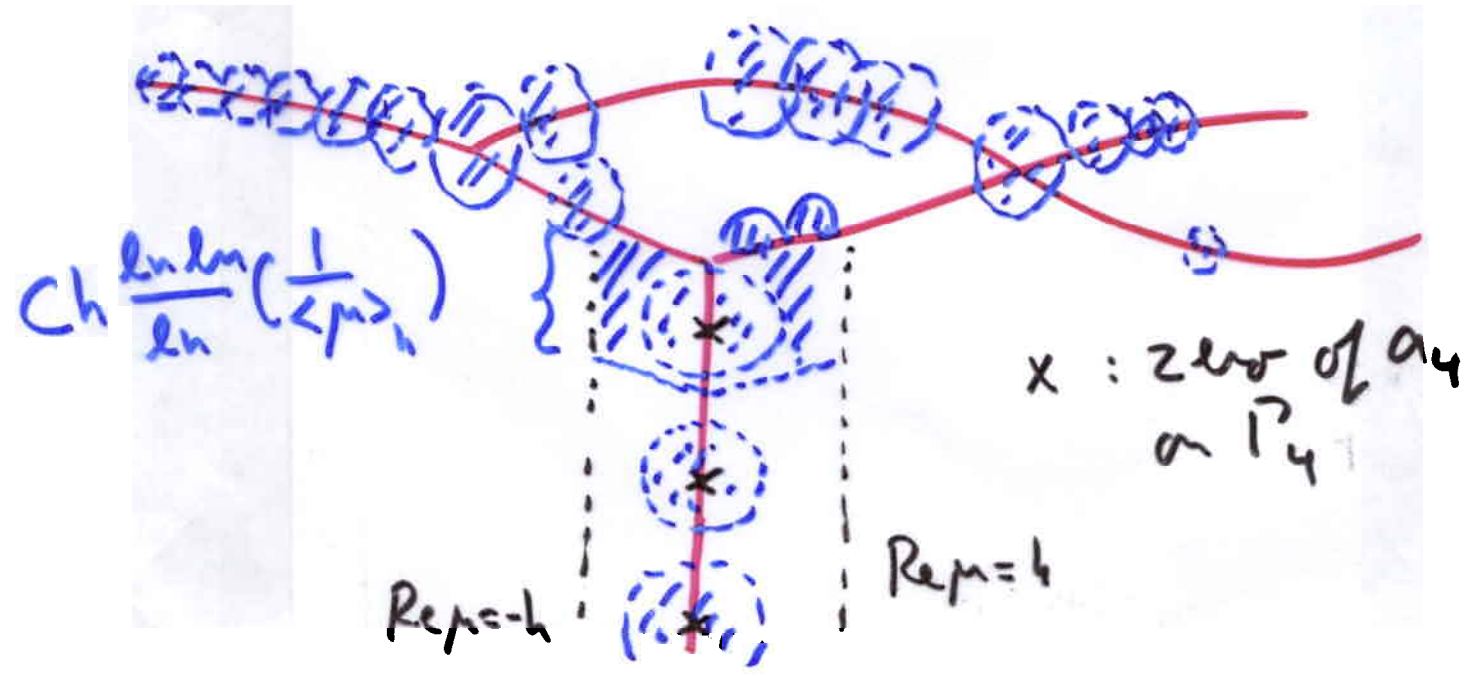
$\Gamma_{2,4}^+ \cup \Gamma_{3,4}^+$  for  $\text{Re } \mu \gg h$ :

$$G = a_4^+ \left[ \left(1 + \frac{a_2}{a_4^+}\right) \left(1 + \frac{a_3}{a_4^+}\right) + e^{-2\pi \mu/h} \right]$$



Concentration of zeros to the skeleton:

$$\langle \mu \rangle_n = (h^2 + |\mu|^2)^{\frac{1}{2}} \quad 0 = D(\mu; \frac{Ch}{\ln \langle \mu \rangle_n})$$

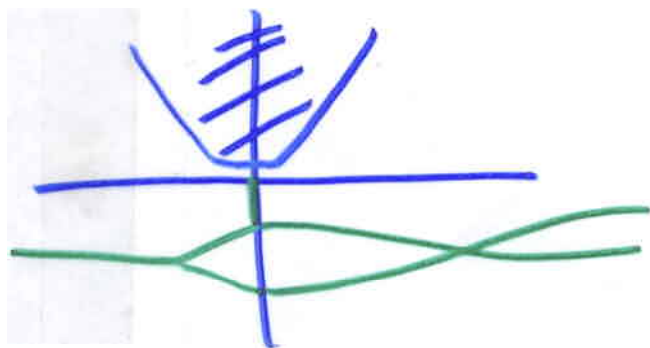




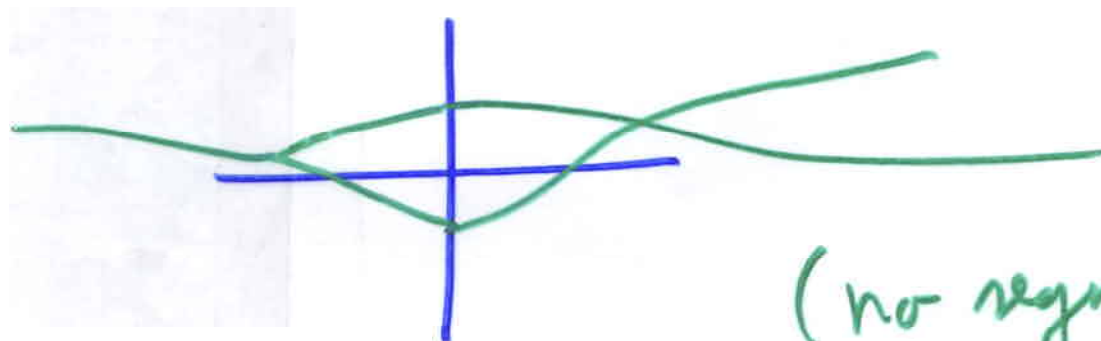
Other possibilities:

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Case 2:  $|\arg \mu + \frac{\pi}{2}| < \pi - \frac{1}{C}$ :



When the skeleton reaches both of  $i\mathbb{R}_{\pm}$ :



(no segment  
in  $i\mathbb{R}$ )

Discussion with Fujie-Ramond:

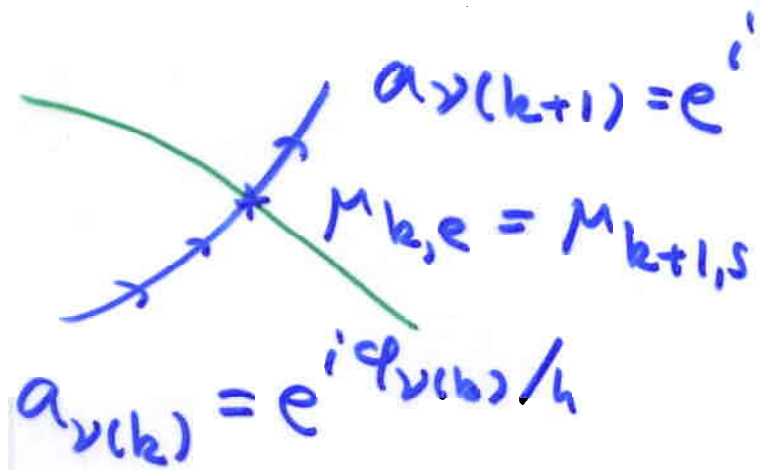
Can determine the dominant term  $a_j$   
away from the skeleton by

Stokes line navigation

$h^0$  of zeros inside a closed curve  $\gamma$  crossing the skeleton transversally at finitely many points

$$\mu_{k,e} = \mu_{k+1,s}, \quad k=0,1,\dots,N-1$$

$$(N-1+1=0)$$



$$\Rightarrow \frac{1}{2\pi h} \sum_{k=0}^{N-1} (\varphi_{\nu(k)}(\mu_{k,e}) - \varphi_{\nu(k+1)}(\mu_{k+1,s})) + O(1) + O(\max \ln \ln \frac{1}{\langle \mu_{k,e} \rangle_h})$$

maximum over all crossing points in the exceptional region.

Inspiration: E.B.Davies : 1D-systems.