

East model: Aging through Hierarchical Coalescence

Fabio Martinelli

Dept. of Mathematics, University Roma 3, Italy

with: A. Faggionato, C.Roberto, C. Toninelli

Overview of the talk

1 I Part : the East model

- East model : definition and main features.
- Previous results on the relaxation to the equilibrium.
- High density, non-equilibrium dynamics.
- Staircase behavior of the density ;
- Aging for density-density time auto-correlation ;
- Scaling limit for the inter-vacancy interval length.

2 II Part : Hierarchical Coalescence Process.

- Main features of a hierarchical coalescence process.
- Universality for the HCP.
- Implications for the East model :

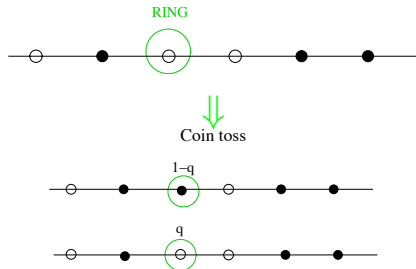
3 Open problems.

East Model [Jackle-Eisinger '91]

- Configuration space $\Omega = \{0, 1\}^{\mathbb{Z}}$
- $\sigma \in \Omega, \quad \sigma(x) \in \{0, 1\} \begin{cases} \nearrow 1 : \text{there is a particle at site } x \\ \searrow 0 : \text{there is no particle at site } x \end{cases}$
- Glauber dynamics with **kinetic constraint** : the right neighbor should be empty for a flip to take place.



Nothing happens



Main Features of the East Model

- $\sigma \in \Omega := \{0, 1\}^{\mathbb{Z}}, \quad \forall x \in \mathbb{Z} : \sigma(x) \in \{0, 1\}$
- $\{\sigma_t\}_{t \geq 0}$ denotes the process with initial config. σ

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$$\mathcal{L}f(\sigma) := \sum_{x \in \mathbb{Z}} (1 - \sigma(x+1))(\mu_x(f) - f(\sigma))$$

- Reversible w.r.t. μ (Bernoulli product measure of density $p = 1 - q$).
- Introduced by physicists to model liquid/glass transition.
- Constraint on the allowed moves (more effective as $q \downarrow 0$) simulates *geometric constraints* on molecules in highly dense liquids.
- As $q \downarrow 0$ constraints slow down the dynamics. The liquid freezes into an amorphous solid state $\sim$ glass.

Relaxation to equilibrium : basic results

Definition

$$T_{\text{rel}} := (\text{spectral gap of } \mathcal{L})^{-1}$$

Theorem

- $T_{\text{rel}} < +\infty \quad \forall q \in (0, 1)$ (Aldous, Diaconis '02)
- $T_{\text{rel}} \simeq \left(\frac{1}{q}\right)^{|\log_2(q)|/2}$ as $q \downarrow 0$ (Cancrini, M., Roberto, Toninelli '08)

Theorem (Cancrini, M., Schonmann, Toninelli. '09)

Let $\nu \neq \mu$ be a different product measure. Then

$$|E_\nu(f(\sigma_t)) - \mu(f)| \leq C_f \exp(-t/2 T_{\text{rel}})$$

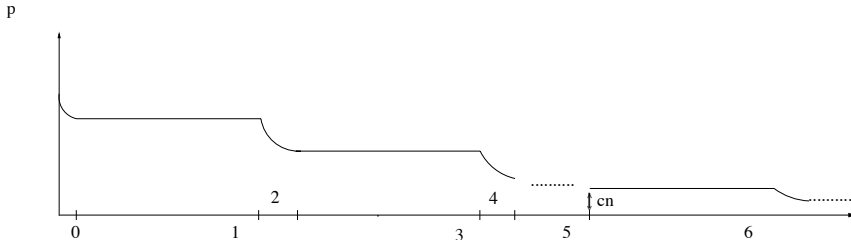
High density non-equilibrium dynamics

Basic Setting

- 1 Start at time zero from a quite general *renewal law* ν independent from q ;
- 2 Run the Glauber dynamics with $q \ll 1$;
- 3 Focus on the *pre-asymptotic* behavior not too far from the origin (or in large q -independent interval) and up to time scales $T = T(q)$ with $1 \ll T \ll T_{\text{rel}}$.
- 4 Physically we are considering a *quench* from *low* to *high* density, i.e. from the liquid to the glass phase ;

Plateau behavior of the density

- Numerical simulations and non-rigorous theoretical analysis suggest that, as $q \rightarrow 0$:
 - (a) the model does not reach equilibrium ;
 - (b) time auto-correlation function shows aging ;
 - (c) the density profile exhibits plateau behavior.



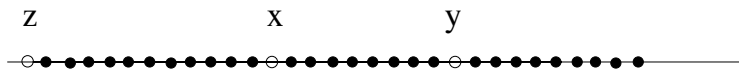
A First Key Observation

As $q \downarrow 0$:

- 1 Dynamics is dominated by **killing the excess vacancies**.
- 2 Dynamics $\iff$ **coarsening of intervals delimited by consecutive vacancies**.
- 3 To kill a vacancy at x the dynamics *must* bring a vacancy at $x + 1$ from the nearest (East-ward w.r.t. x) vacancy $\Rightarrow$ *cooperative* relaxation.
- 4 Cooperative relaxation requires a certain number of auxiliary *extra* vacancies to be first created and then destroyed $\Rightarrow$ *Energy Barriers*.
- 5 *Metastable* effects very relevant.
- 6 Key question : what is the *structure* of the energy barriers ?

Energy Barriers Structure and Activation Times

$$y - x \in [2^{n-1} + 1, 2^n], \quad 1 \leq n \ll \log_2(1/q)$$



Combinatorial argument (Evans, Sollich + Chung-Diaconis-Graham 01) : during the killing of vacancy at x at least n extra vacancies between x and y .

Energy barrier $\Delta E_n = n \Rightarrow$ Activation Time $t_n := (1/q)^n$.

Metastability : Actual killing is *random and instantaneous* (w.r.t. to the expected time t_n) and occurs on scale $t_{n-1} = \frac{1}{q^{n-1}}$.

Killing vacancy at $x \Leftrightarrow$ **coalescing domain** $[x, y]$ with $[z, x]$

Active and stalling periods

Definition (Hierarchy of activation times)

$$t_n^+ := t_n^{1+\epsilon} = \left(\frac{1}{q^n}\right)^{1+\epsilon}, \quad t_n^- := t_n^{1-\epsilon} = \left(\frac{1}{q^n}\right)^{1-\epsilon}$$

Definition

Domain $[x, x + \ell]$ is of class n if $\ell \in [2^{n-1} + 1, 2^n]$, $n \geq 1$.
Vacancy at x is of class n if it is the left border of a domain of class n .

Definition

- $[t_n^-, t_n^+]$ is called the n -th active period ;
- $[t_n^+, t_{n+1}^-]$ is called the n -th stalling period (nothing happens).

Evolution during the n -th active period.

- Recursively assume that at time t_n^- all vacancies in e.g. $[-2^N, 2^N]$, $N \gg 1$, are of class at least n w.h.p.
- Vacancies of class *larger* than n do not disappear w.h.p.
- When a vacancy of the n -th class disappears $\Leftrightarrow$ the class of the vacancy to its *left* becomes at least $n + 1$ ($2^n + 2^n = 2^{n+1}$).

- Thus vacancies of class n either disappear *directly* or *increase* their class w.h.p.
- At the end of the period (t_n^+) w.h.p. *all* vacancies are of class at least $n + 1$ and were already present at t_n^- ;
- Between the *stalling* period $[t_n^+, t_{n+1}^-]$ w.h.p. no vacancy is killed $\Rightarrow$ the recursion step is proved.

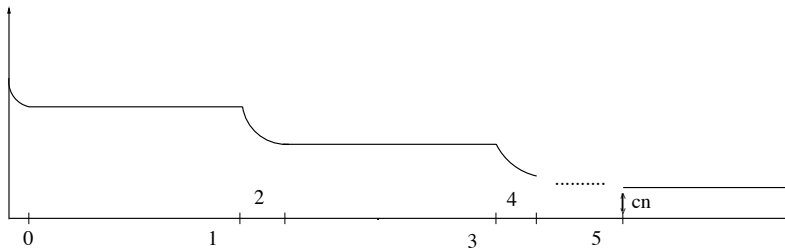
Staircase Behavior

Theorem

$\forall \sigma \in \{0,1\}^{\mathbb{Z}}, \forall n \in \mathbb{N}$ there exists $c_n(\sigma)$ s.t.

$$\lim_{q \downarrow 0} \sup_{t \in [t_n^+, t_{n+1}^-]} |\mathbb{P}_\sigma(\sigma_t(0) = 0) - c_n(\sigma)| = 0$$

p



Height of plateaux

Definition

Let ν be a (non-trivial) probability law over the integers with finite mean. We define $\text{Ren}(\nu)$ as the stationary renewal distribution on $\mathbf{Z}$ with interval law ν

Theorem

If the initial distribution is $Q = \text{Ren}(\nu)$ then

$$\lim_{n \rightarrow \infty} \lim_{q \downarrow 0} \mathbb{E}_Q(c_n)(2^n + 1) = 1$$

Coarsening

Let $\tilde{X}$ be the non-negative random variable such that

$$\mathbb{E}(\exp(-s\tilde{X})) = 1 - \exp\left\{-\int_1^\infty \frac{e^{-sx}}{x} dx\right\} \quad s > 0$$

Theorem

If the initial distribution is $Q = \text{Ren}(\nu)$ then for any bounded function f and any $k \in \mathbb{Z}$

$$\lim_{n \uparrow \infty} \lim_{q \downarrow 0} \sup_{t \in [t_n^+, t_{n+1}^-]} \left| \mathbb{E}_Q(f(X_k^{(n+1)}(t))) - \mathbb{E}(f(\tilde{X})) \right| = 0$$

where $X_k^{(n)}(t) := (x_{k+1}(t) - x_k(t))/(2^{n-1} + 1)$

$\implies$ Mean domain length grows as $t^{\log 2 / |\log q|}$

Aging

Theorem

Fix $\sigma \in \{0, 1\}^{\mathbb{Z}}$, $\forall n, m \in \mathbb{N}$ there exists $c_{n,m,x}(\sigma)$ s.t.

$$\lim_{q \downarrow 0} \sup_{t \in [t_n^+, t_{n+1}^-]} \sup_{s \in [t_m^+, t_{m+1}^-]} |\text{Cov}_{\sigma}(\sigma_t(x); \sigma_s(x)) - c_{n,m,x}(\sigma)| = 0$$

If $Q = \text{Ren}(\nu)$ set $\rho_x := Q(\sigma(x) = 0)$, then

$$\lim_{n \rightarrow \infty} \lim_{m \rightarrow \infty} \lim_{q \downarrow 0} \left| \mathbb{E}_Q(c_{n,m,x}) - \frac{\rho_x}{(2^n + 1)} \left(1 - \frac{\rho_x}{(2^m + 1)} \right) \right| = 0$$

Hierarchical Coalescence Process (HCP)

- Main motivation : physical modeling of non-equilibrium coarsening dynamics of 1D systems (Derrida, Bray, Godreche,...) ;
- Features : infinite sequence of 1D coalescence processes $\{\xi_n\}_{n=1}^{\infty}$ with “end of ξ_n = beginning of ξ_{n+1} ”.
- The n -th process live in the n -th *epoch* and *only* intervals with length $\in [d_n, d_{n+1})$ are *active* i.e. merge with very general rates with the left or right neighbor.
- We assume $d_n \uparrow$ and $2d_n \geq d_{n+1}$.
- Examples : $d_n = n$ (“Paste-all-model”), $d_n = a^n$, $a \in (1, 2]$, $d_n = 2^{n-1} + 1 \Leftrightarrow$ the East model.

Some Facts

- 1 Assume that at the beginning of the first epoch the intervals form a renewal process
- 2 Physicists observed a large scale (i.e. $n \rightarrow \infty$) universality of rescaled variables (e.g. (domain length)/ d_n).
- 3 Different models show the same behavior independently of the merging rates and of the active domains range. For example the “Paste-all-model” behaves as the East model a fact refereed to as “very surprising” in the physics literature.
- 4 Physicists were not able to prove the scaling hypothesis nor were able to classify the universality classes according to the initial renewal process.

HCP definition I

$\Omega^{(n)}$ is the subset of $\Omega := (0, 1)^{\mathbb{Z}}$ with each domain (=interval among consecutive zeros) of class at least n .
Class 0 if $d = 1$, class $n \geq 1$ if $d \in [d_n, d_{n+1})$.

n -th epoch Coalescence Process

Initial configuration : $\sigma_0^{(n)} \in \Omega^{(n)}$

Dynamics :

- independent exponential clock on each domain $[x, y]$ with rate $\lambda_n(y - x)$;
- $\lambda_n(d) > 0$ iff $d \in [d_n, d_{n+1})$;
- when clock rings on $[x, y]$ and if the domain is still present $\Rightarrow$ incorporate left domain, i.e. erase point x .

At infinite time the configuration $\sigma_\infty^{(n)}$ belongs to $\Omega^{(n+1)}$

HCP definition II (infinite epochs)

- 1 Start from $\sigma \in \Omega = \Omega^{(0)}$ and run the 0-th epoch coalescence process for an infinite time to get $\sigma_{\infty}^{(0)}$;
- 2 Start from $\sigma_{\infty}^{(0)} \in \Omega_1$ and run the 1-th epoch coalescence process for an infinite time to get $\sigma_{\infty}^{(1)}$;
- 3 ...
- 4 Start from $\sigma_{\infty}^{(n-1)} \in \Omega^{(n)}$ and run the n -th epoch coalescence process for an infinite time to get $\sigma_{\infty}^{(n)}$
- 5 ...

HCP Main Results

Theorem

Assume $\sigma \sim Q = \text{Ren}(\nu)$ with ν any finite mean probability measure on $[1, \infty)$. Then (i) $\sigma_t^{(n)}$ is distributed with $\text{Ren}(\nu_t^{(n)})$; (ii) Let $X^{(n)}$ be distributed with $\nu_0^{(n)}$ and $Z^{(n)} := X^{(n)}/d_n$. Then $Z^{(n)}$ weakly converges to $Z^{(\infty)}$ with

$$\mathbb{E}(e^{-sZ^{(\infty)}}) = 1 - \exp \left\{ - \int_1^\infty \frac{e^{-sx}}{x} dx \right\}$$

(iii) If instead $Q = \text{Ren}(\nu|0)$ with ν in the domain of attraction of an α -stable law then

$$\mathbb{E}(e^{-sZ^{(\infty)}}) = 1 - \exp \left\{ - \alpha \int_1^\infty \frac{e^{-sx}}{x} dx \right\}$$

Hints for the proof

- 1 The Laplace transforms $\{g^{(n)}(s)\}_{n \geq 1}$ of $\{Z^{(n)}\}_{n \geq 1}$ satisfy a highly non-linear system of recursive identities

$$1 - g^{(n)}(s a_{n-1}) = (1 - g_{n-1}(s))e^{h^{(n-1)}(s)} \quad \forall n \geq 2.$$

where

$$a_n = d^{(n+1)} / d^{(n)}, \quad h^{(n)}(s) = \mathbb{E} \left(e^{-sZ^{(n)}} \chi_{1 \leq Z^{(n)} < a_n} \right)$$

- 1 One observes that a family of fixed points is given by

$$g_c^{(\infty)}(s) = 1 - \exp \left(-c \int_1^\infty dx \frac{e^{-sx}}{x} \right), \quad c \in (0, 1]$$

- 2 One then *proves* that there exists a non-negative measure $dt^{(n)}(x)$ such that

$$g^{(n)}(s) = 1 - \exp \left(- \int_1^\infty dt^{(n)}(x) \frac{e^{-sx}}{x} \right)$$

- 3 The recursive identities for the $\{g^{(n)}(s)\}_{n \geq 1}$ become identities for the measures $\{t^{(n)}\}_{n \geq 1}$ and one can *prove* that $t^{(n)} \rightarrow c \times \text{Leb}$ if

$$\lim_{s \downarrow 0} -s \frac{d}{ds} g^{(1)}(s) / (1 - g^{(1)}(s)) = c (\in (0, 1])$$

HCP Universality and generalizations

- Previous result $\Rightarrow$ *Universality Classes*.
- One can generalize to $Q = \text{Ren}(\mu, \nu)$: scaling for the position of the first vacancy + same scaling for the rescaled domain length.
- Same method for triple coalescence (merging with left *and* right neighboring intervals) but *different* asymptotics.

East model vs HCP

Discussion about activation times and energy barriers suggest that, as $q \downarrow 0$, the East dynamics should be approximated by a suitable HCP such that ;

- 1 $d_n = 2^{n-1} + 1$;
- 2 time t_n^+ should corresponds to the end of the n -th coalescence ;
- 3 It remains to determine the appropriate rates $\lambda_n(d) \Leftrightarrow$ survival propability of a vacancy with domain d (of class n).
- 4 The rates are found solving a large deviation problem for the East dynamics. All what is needed is that the rates only weakly depend on q .

Main result

- The final result is that, on finite, large volume *independent* from q , the marginal for the East process and for the HCP above are close in $\|\cdot\|_{\text{TV}}$ for $q \downarrow 0$.
- Staircase behavior and aging follow at once from the scaling limit of the HCP.

References

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