

LIUVILLE QUANTUM GRAVITY, KPZ
& SCHRAMM-LOEWNER EVOLUTION

Bertrand Duplantier

Institut de Physique Théorique, Saclay, France

Scott Sheffield

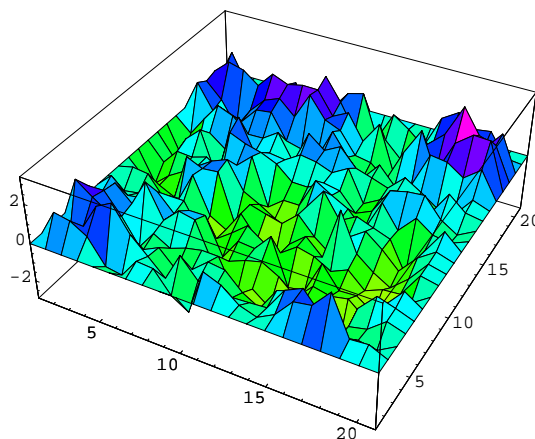
MIT Math, USA

STATISTICAL MECHANICS & CONFORMAL INVARIANCE

MSRI

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Gaussian Free Field (GFF)



Distribution h with *Gaussian weight* $\exp\left[-\frac{1}{2}(h, h)_{\nabla}\right]$, and **Dirichlet inner product** in domain D

$$\begin{aligned}(f_1, f_2)_{\nabla} &:= (2\pi)^{-1} \int_D \nabla f_1(z) \cdot \nabla f_2(z) d^2z \\ &= \text{Cov}\left((h, f_1)_{\nabla}, (h, f_2)_{\nabla}\right)\end{aligned}$$

◇ STARRING THE GFF! (Courtesy of N.-G. Kang) ◇

LIOUVILLE QG

RANDOM MEASURE

$$d\mu = "e^{\gamma h} d^2z"$$

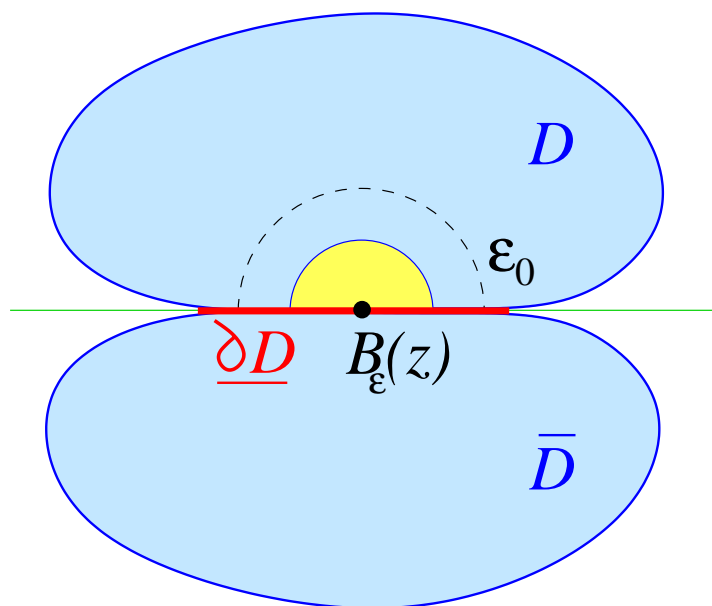


THE EMERGENCE OF QUANTUM GRAVITY

(Courtesy of N.-G. Kang)



Bulk & Boundary Liouville Quantum Gravity



- Circle averages $h_\epsilon(z)$, $z \in D$ (Dirichlet)
- GFF with free boundary conditions on $\underline{\partial D}$
- Half-circle averages $\hat{h}_\epsilon(z)$, $z \in \underline{\partial D}$.

- Regularization

$h_\varepsilon(z)$ mean value of h on circle $\partial B_\varepsilon(z)$

- Variance

$$\text{Var } h_\varepsilon(z) = \log[\mathbf{C}(z, D) / \varepsilon]$$

$\mathbf{C}(z, D)$ conformal radius of D viewed from z

$h_\varepsilon(z)$ Gaussian random variable

$$\mathbb{E} e^{\gamma h_\varepsilon(z)} = e^{\gamma^2 \text{Var } h_\varepsilon(z) / 2} = \left(\frac{\mathbf{C}(z, D)}{\varepsilon} \right)^{\gamma^2 / 2} \quad \square$$

QUANTUM AREA MEASURE

$$d\mu_\varepsilon := \exp[\gamma h_\varepsilon(z)] \varepsilon^{\gamma^2/2} d^2z$$

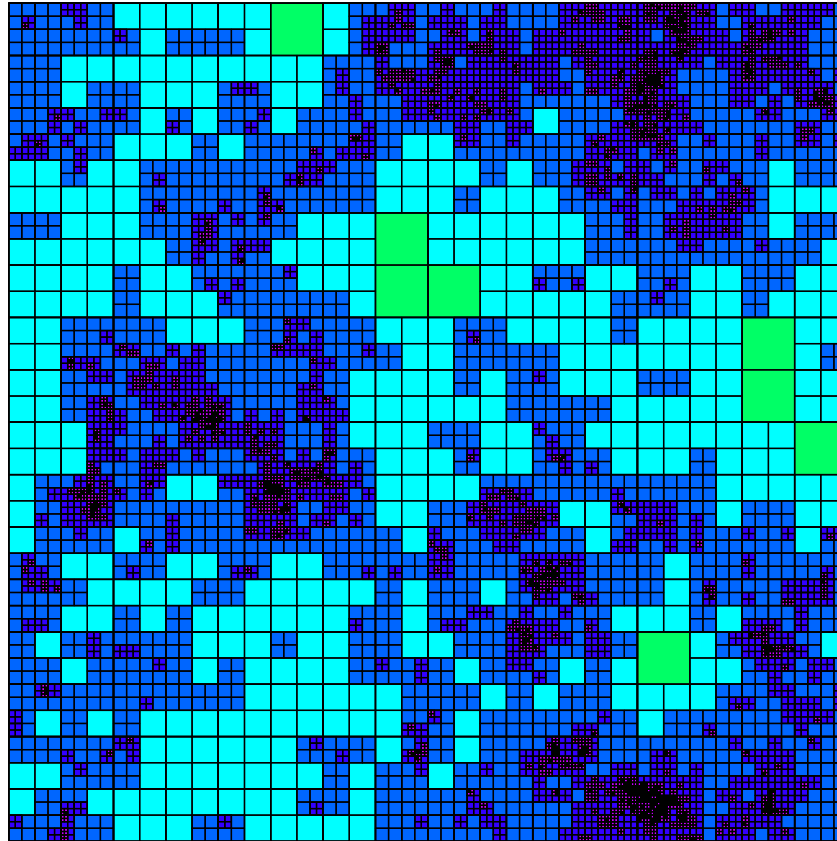
converges, as $\varepsilon \rightarrow 0$ and for $\gamma < 2$, to a random measure, denoted by $e^{\gamma h(z)} d^2z$.

QUANTUM BOUNDARY MEASURE

$$d\hat{\mu}_\varepsilon := \exp\left[\frac{\gamma}{2} \hat{h}_\varepsilon(z)\right] \varepsilon^{\gamma^2/4} dz$$

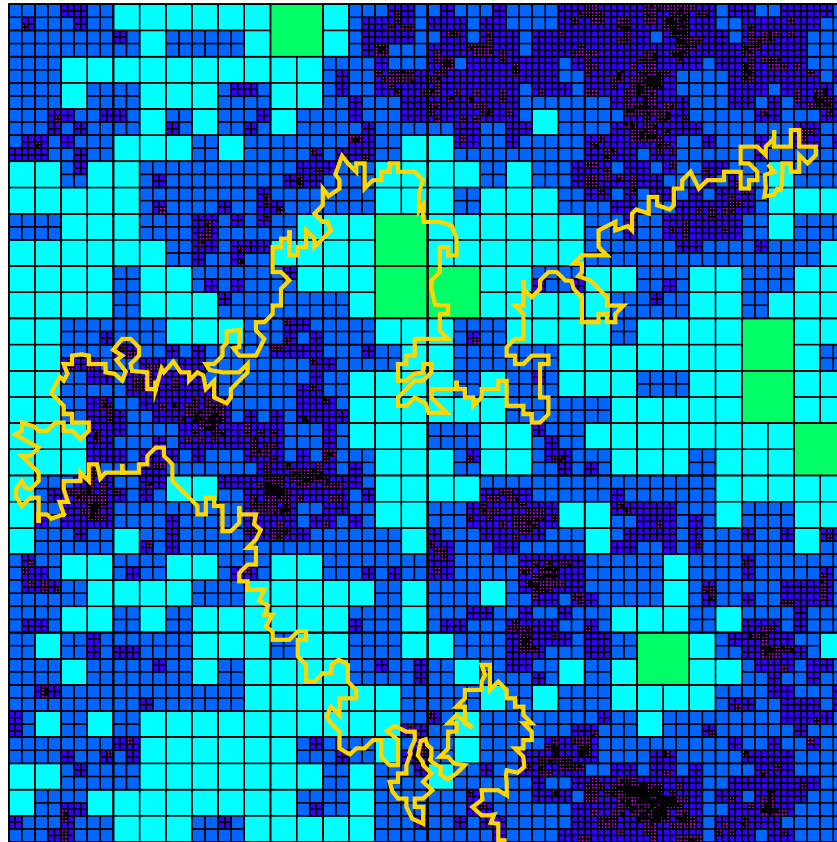
converges, as $\varepsilon \rightarrow 0$ and for $\gamma < 2$, to a *boundary* random measure, denoted by $e^{(\gamma/2)h(z)} dz$.

Random Surface



Euclidean squares of similar quantum area

Random Surface & Fractal Sets



Euclidean & Quantum Fractal Measures

The d -dimensional *Euclidean* or analogously *quantum measure* of planar *fractal* sets is characterized by scaling properties:

- Rescale a d -dimensional fractal $X \subset \mathcal{D} \subset \mathbb{C}$ via the map $z \rightarrow \psi(z) = bz$, $b \in \mathbb{C}$ (so that the Euclidean area of the domain \mathcal{D} is multiplied by $|b|^2$); then the d -dimensional *Euclidean fractal measure* of X is multiplied by $|b|^d = |b|^{2-2x}$, where x (the *Euclidean scaling weight*) is defined by $d := 2 - 2x (\leq 2)$.
- If X is a *fractal subset* of a *random surface* $\mathcal{S} := (\mathcal{D}, h)$, and we rescale \mathcal{S} so that its quantum area increases by a factor of $|b|^2$, then the *quantum fractal measure* $Q(X, h)$ of X is multiplied by $|b|^{2-2\Delta}$, where Δ is the analogous *quantum scaling weight*.

- $Q(\psi(X, h)) = Q(X, h)$ whenever ψ is conformal and

$$\psi(\mathcal{D}, h) := (\psi(\mathcal{D}), h \circ \psi^{-1} - Q \log |\psi'|)$$

$$Q := \frac{\gamma}{2} + \frac{2}{\gamma}.$$

The pair $\mathcal{S} = (\mathcal{D}, h)$ describes the same Liouville quantum surface (up to coordinate change) as the conformally transformed pair $\psi(\mathcal{D}, h)$.

- The *Knizhnik, Polyakov, Zamolodchikov (KPZ)* relation

$$x = (\gamma^2/4)\Delta^2 + (1 - \gamma^2/4)\Delta$$

is then equivalent to

$$d = \alpha Q - \alpha^2/2,$$

where $d = 2 - 2x$ and $\alpha := \gamma(1 - \Delta)$.

SLE - GFF (QG) COUPLING

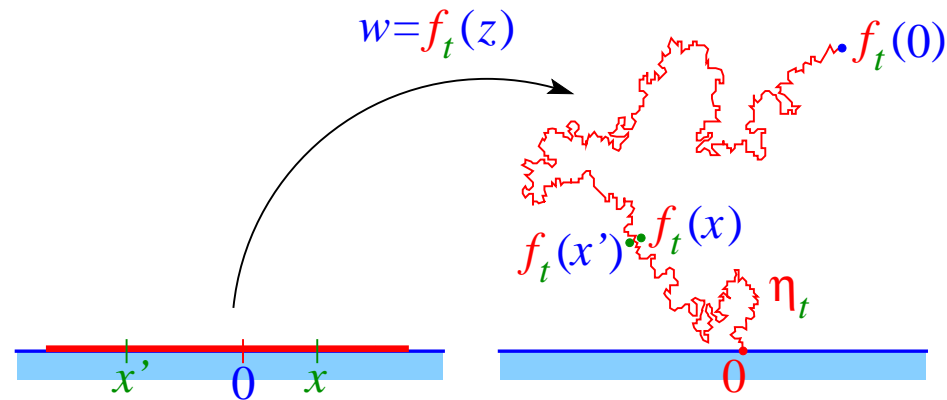
Dubédat, 2009

Sheffield, arXiv:1012.4797

*D. & Sheffield, PRL **107**, 131305 (2011), arXiv:1012.4800*

Miller & Sheffield, arXiv:1201.14896-98

“Zipping-up” SLE Map



Let f_t be the (reverse) SLE_κ conformal map

$$z \in \mathbb{H} \rightarrow w = f_t(z) \in \mathbb{H} \setminus \eta_t,$$

with trace η_t and tip $f_t(0)$ [$t = 0$, $f_0(z) = z$]. It satisfies the stochastic differential equation (B_t standard Brownian motion)

$$df_t(z) = -2dt/f_t(z) - \sqrt{\kappa}dB_t.$$

(Reverse) SLE Martingale

Real stochastic process in the upper-half plane:

$$h_0(z) := \frac{2}{\sqrt{\kappa}} \log |z|,$$

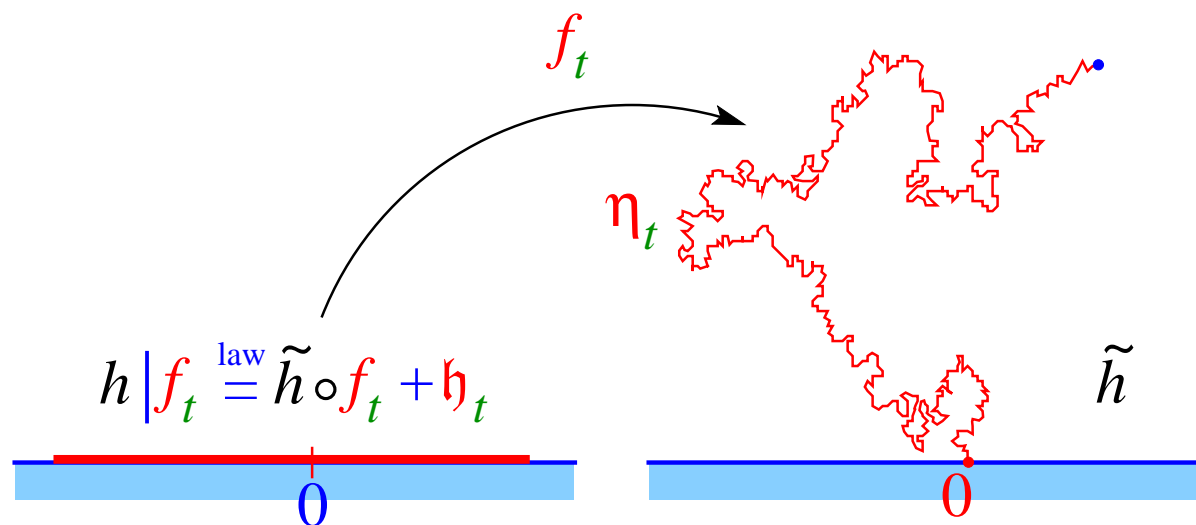
$$h_t(z) := h_0 \circ f_t(z) + Q \log |f_t'(z)|.$$

This process $h_t(z)$ is a *martingale* (so that $\mathbb{E}h_t(z) = h_0(z)$) for the particular choice:

$$Q = \sqrt{\kappa}/2 + 2/\sqrt{\kappa},$$

for which $dh_t(z) = -\Re[2/f_t(z)]dB_t$.

SLE–GFF Coupling

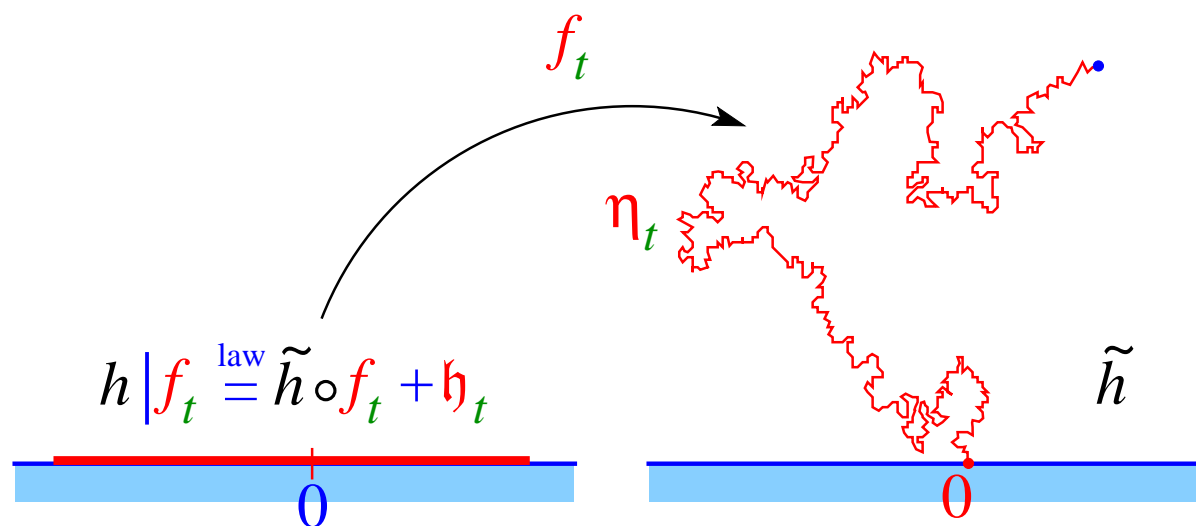


Define $h := \tilde{h} + \mathfrak{h}_0$, sum of the GFF \tilde{h} on \mathbb{H} with *free boundary conditions* on \mathbb{R} , and of the deterministic function \mathfrak{h}_0 . Given f_t , the conditional law of h (denoted by $h|f_t$) is

$$h(z)|f_t \stackrel{(\text{law})}{=} \tilde{h} \circ f_t(z) + \mathfrak{h}_t(z),$$

where $\tilde{h} \circ f_t$ is the pullback of the free boundary GFF \tilde{h} .

SLE–GFF Coupling



$$h(z)|f_t \stackrel{\text{(law)}}{=} \tilde{h} \circ f_t(z) + \mathfrak{h}_t(z)$$

To sample h , one can first sample the B_t process (which determines f_t), then sample independently the f.b.c. GFF \tilde{h} and take the above sum [Sheffield, 2010].

The conditional expectation w.r.t. \tilde{h} is the *martingale*

$$\mathbb{E}[h(z)|f_t] = \mathfrak{h}_t(z).$$

Neumann Green function

Consider the *Neumann Green function* in \mathbb{H} ,

$G_0(y, z) := -\log(|y - z||y - \bar{z}|)$, and define the

time-dependent $G_t(y, z) := G_0(f_t(y), f_t(z))$, i.e., G_0 taken at

image points under f_t . A calculation of the Green function's

variation shows that $-dG_t(y, z) = d\langle \mathfrak{h}_t(y), \mathfrak{h}_t(z) \rangle$

(*Hadamard's formula*). Integrating w.r.t. t yields the

covariation of the \mathfrak{h}_t martingales

$$\langle \mathfrak{h}_t(y), \mathfrak{h}_t(z) \rangle = G_0(y, z) - G_t(y, z).$$

In the $y \rightarrow z$ limit:

$$\langle \mathfrak{h}_t(z), \mathfrak{h}_t(z) \rangle = C_0(z) - C_t(z),$$

where $C_t(z) := -\log [\Im f_t(z) |f'_t(z)|]$.

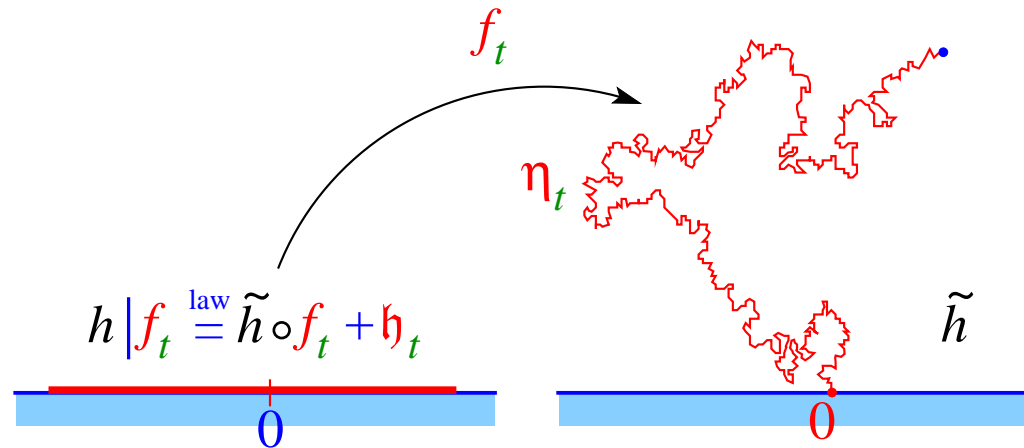
SLE–GFF Coupling

Define the *covariance*: $\text{Cov}[A, B] := \mathbb{E}[AB] - \mathbb{E}[A]\mathbb{E}[B]$.

Recall that the Green's function $G_0(y, z) = \text{Cov}[\tilde{h}(y), \tilde{h}(z)]$, thus $G_t(y, z) = \text{Cov}[\tilde{h} \circ f_t(y), \tilde{h} \circ f_t(z)]$. The random distribution $\tilde{h} \circ f_t$ and the set of (time changed) Brownian motions \mathfrak{h}_t are Gaussian processes, whose respective covariance G_t and covariation $\langle \mathfrak{h}_t, \mathfrak{h}_t \rangle$ thus add to constant G_0 :

$$\begin{aligned} G_t(y, z) + \langle \mathfrak{h}_t(y), \mathfrak{h}_t(z) \rangle &= G_0(y, z) \\ \text{Cov}[\tilde{h} \circ f_t(y), \tilde{h} \circ f_t(z)] + \langle \mathfrak{h}_t(y), \mathfrak{h}_t(z) \rangle &= \text{Cov}[\tilde{h}(y), \tilde{h}(z)] \\ &= \text{Cov}[h(y), h(z)] \quad \square \end{aligned}$$

Liouville Invariance



Recall that $h := \tilde{h} + \mathfrak{h}_0$, and $\mathfrak{h}_t := \mathfrak{h}_0 \circ f_t + Q \log |f_t'|$. Hence $\tilde{h} \circ f_t + \mathfrak{h}_t = h \circ f_t + Q \log |f_t'|$. For $Q = \gamma/2 + 2/\gamma$, this is the transformation law of the GFF h under the conformal map f_t^{-1} . The pair $(\mathbb{H}, \tilde{h} \circ f_t + \mathfrak{h}_t) = f_t^{-1}(\mathbb{H} \setminus \eta_t, h)$ describes the same random surface as the pair $(\mathbb{H} \setminus \eta_t, h)$: *Given f_t , the image under f_t of the measure $e^{\gamma h(z)} d^2z$ in \mathbb{H} is a random measure whose law is the *a priori* (unconditioned) law of $e^{\gamma h(w)} d^2w$ in $\mathbb{H} \setminus \eta_t$.*

Liouville Quantum Measure

$$(e^{\gamma h(z)} | f_t) d^2 z \stackrel{(\text{law})}{=} e^{\gamma h(w)} d^2 w \quad (\text{conformal invariance})$$

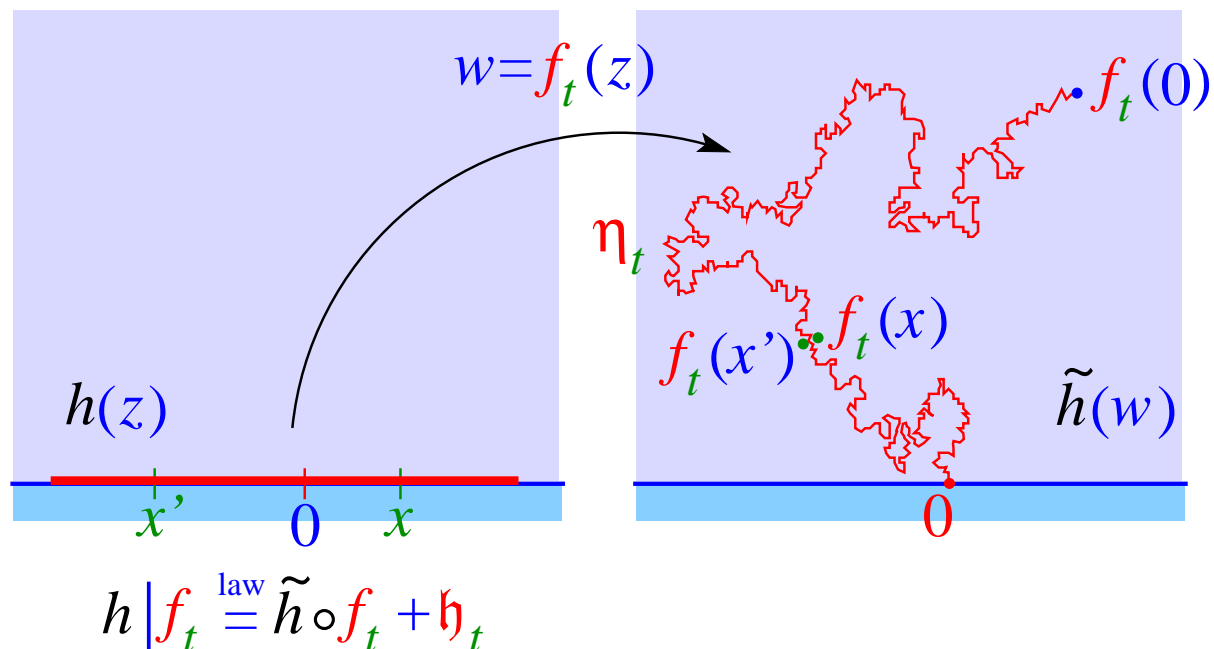
for $d = 2 = \gamma Q - \gamma^2/2$, i.e., $Q = \gamma/2 + 2/\gamma = \sqrt{\kappa}/2 + 2/\sqrt{\kappa}$

$$\implies \gamma = \sqrt{\kappa} \wedge (4/\sqrt{\kappa}), \quad \gamma' = 4/\gamma$$

- $\gamma \leq 2$: *KPZ prediction* $\gamma = (\sqrt{25 - c} - \sqrt{1 - c})/\sqrt{6}$ for the *central charge* $c = \frac{1}{4}(6 - \kappa)(6 - 16/\kappa) \leq 1$ of the SLE's CFT coupled to gravity.
- $\gamma' = 4/\gamma > 2$: *Duality* property of Liouville quantum gravity; the quantum measure develops atoms with localized area.

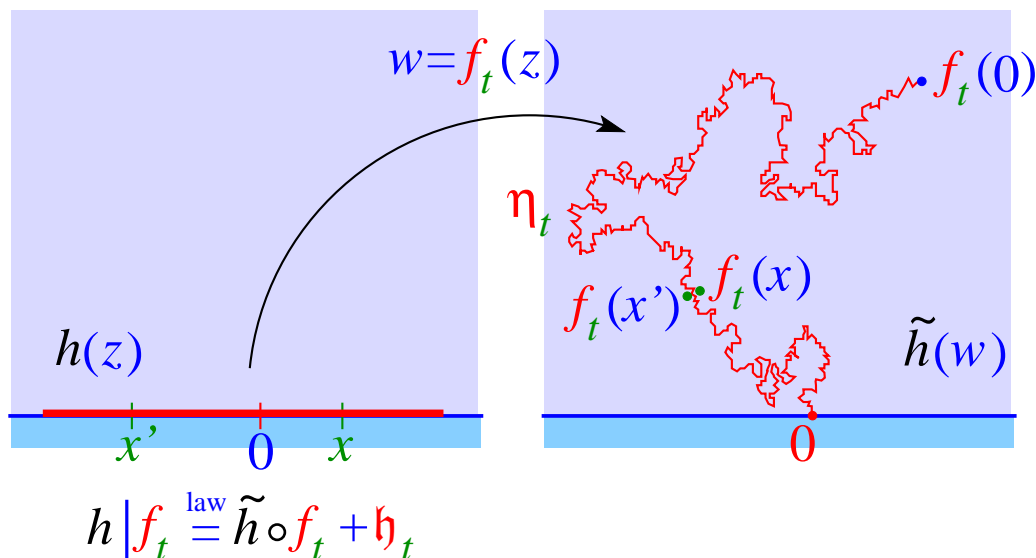
Conformally welding two γ -Liouville quantum surfaces produces SLE_{κ} .

Conformal Welding



Conformal welding: the *quantum boundary lengths* of any pair of real segments $[0, x]$ and $[x', 0]$ such that $f_t(x) = f_t(x')$ on the SLE trace are *a.s. equal* for $h = \tilde{h} + \mathfrak{h}_0$ [Sheffield, 2010].

Liouville Quantum Gravity & SLE



- Exponential martingales yield SLE quantum measures:

$$\mathbb{E}[h|f_t] = \mathfrak{h}_t, \quad \mathbb{E}(e^{\alpha h}|f_t) = \exp[\alpha \mathfrak{h}_t - (\alpha^2/2)\langle \mathfrak{h}_t, \mathfrak{h}_t \rangle]$$

[D. & Sheffield, PRL 107, 131305 (2011)]

SLE Exponential Martingales & KPZ

$$\mathcal{M}_t^\alpha(z) := \mathbb{E}(e^{\alpha h(z)} | f_t), \quad \alpha \in \mathbb{R}$$

$$(e^{\alpha h(z)} | f_t) d^2 z \stackrel{(\text{law})}{=} |f'_t(z)|^{d-2} e^{\alpha h(w)} d^2 w$$

$$d := \alpha Q - \alpha^2 / 2 \quad (\text{KPZ})$$

where $w = f_t(z)$, $d^2 w = |f'_t(z)|^2 d^2 z$.

SLE Exponential Martingales

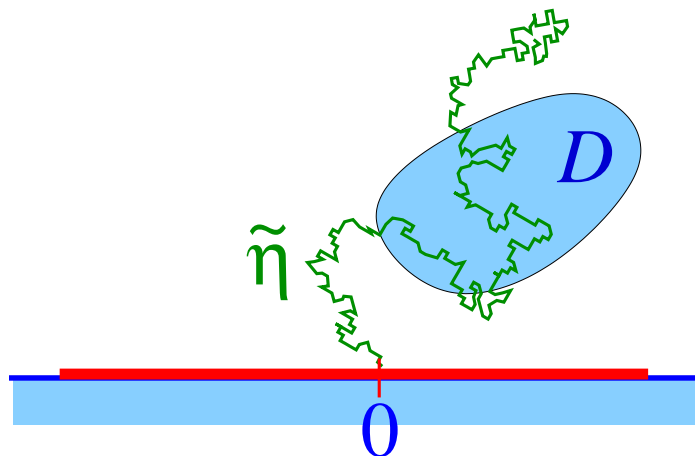
- Conditional expectation w.r.t. GFF h : $\mathbb{E}[h(z)|f_t] = \mathfrak{h}_t(z)$.
- Conditional expectations of exponentials:

$$\begin{aligned}\mathcal{M}_t^\alpha(z) &:= \mathbb{E}(e^{\alpha h(z)}|f_t), \quad \alpha \in \mathbb{R} \\ &= \exp[\alpha \mathfrak{h}_t(z) - (\alpha^2/2)C_t(z)] \\ &= |f'_t(z)|^d |w|^{2\alpha/\sqrt{\kappa}} (\mathfrak{I}w)^{-\alpha^2/2}; \quad d := \alpha Q - \alpha^2/2 \\ C_t(z) &:= \langle \mathfrak{h}_t(z), \mathfrak{h}_t(z) \rangle = \log[\mathfrak{I}f_t(z)|f'_t(z)|]\end{aligned}$$

where $w = f_t(z)$; $\mathcal{M}_t^\alpha(z)$ is an *exponential martingale* with respect to the *Brownian motion driving the SLE process*:

$$\mathbb{E}\mathcal{M}_t^\alpha(z) = \mathcal{M}_0^\alpha(z) = |z|^{2\alpha/\sqrt{\kappa}} (\mathfrak{I}z)^{-\alpha^2/2}.$$

SLE Natural Length



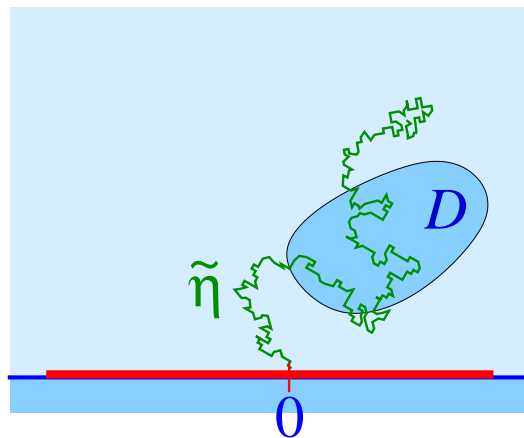
Expected (w.r.t. the $\text{SLE}_{\kappa \in [0,8]}$ law) *length* of an infinite SLE $\tilde{\eta}$ in D (Lawler & Sheffield, 2009)

$$v(D) = \int_D G(z) d^2z,$$

SLE Green's function in \mathbb{H} :

$$G(z) := |z|^a |\Im z|^b, \quad a = 1 - 8/\kappa, \quad b = 8/\kappa + \kappa/8 - 2.$$

SLE Quantum Length



$$h = \tilde{h} + \mathfrak{h}_0$$

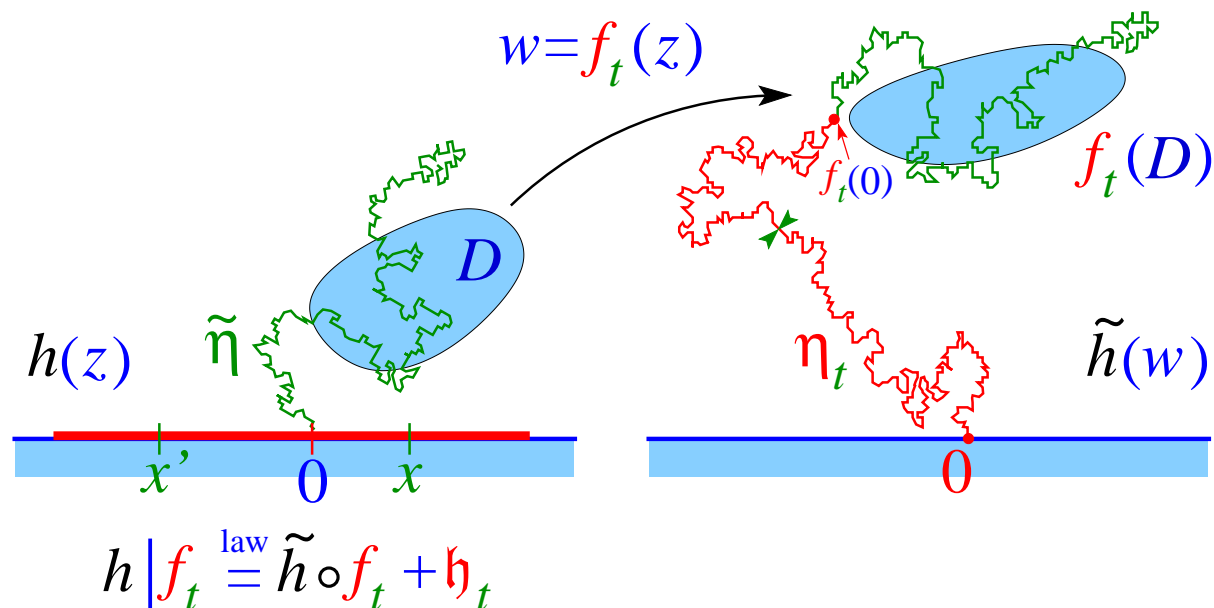
Expected (w.r.t. $\tilde{\eta}$, given h) Liouville *quantum length* v_Q in D

$$v_Q(D, h) := \int_D e^{\alpha h(z)} G(z) d^2z,$$

$\alpha = \sqrt{\kappa}/2$ ($= \gamma/2$ for $\kappa \leq 4$, and $\gamma'/2$ for $\kappa > 4$) satisfies KPZ for the SLE Hausdorff dimension $d = 1 + \kappa/8$.

[Doob-Meyer, second moment method.]

Expected SLE Quantum Length



$$\mathbb{E}[\mathbf{v}_Q(D, h) | f_t] = \int_D \mathcal{M}_t^\alpha(z) G(z) d^2z$$

$$\mathbb{E} \mathbf{v}_Q(D, h) = \int_D \mathcal{M}_0^\alpha(z) G(z) d^2z = \int_D (\sin \vartheta)^{8/\kappa - 2} d^2z,$$

with $\vartheta := \arg z$. It is finite for $\kappa \in [0, 8)$ and coincides with the *Euclidean area* of D for $\kappa = 4$.

PERSPECTIVES

- *Scaling limits of discrete models on random planar graphs*
- *Quantum wedges and cones*
- *Quantum bubbles and foam ($\gamma\gamma' = 4$ duality)*
- *Geodesics & random metrics*

