Klanerman

 F_{19} $\sqrt{2}$ $\frac{\chi}{\chi}$] a null znd ford forms.
 $\hat{\chi}, \hat{\chi}$ are traceless parts. $Fig 2$ doubly null fulnation \overline{u} $\overline{\mathbf{u}}$ S **parger** $F_{9}3$ $\frac{1}{2}$ normally $t=0$ but if deformed the outgoing here will have to
shrink locally. $F_{19}4$ befored ZK. Killing V.F. along hon-zon. maximal hypersurface

ON THE FORMATION OF TRAPPED SURFACES

SERGIU KLAINERMAN

Initial data slide: We can specify the traceless part of the null 2nd fundamental forms, as in figure 1.

Main ideas slide: See figure 2.

Intuition of result (not slide): Figure 3. If I start with something deformed, the area is going down in a local region, and so $tr\chi$ is negative locally. But this won't work globally since can't do it at all points. This result says I can deform the initial surface everywhere except around a point enough so that I get $tr\chi$ negative everywhere except the point, and then use the short pulse to get the rest of the trapped surface. $tr\chi$ is controlled from the first part, the existence result.

Local rigidity slide: Need to extend Z . If I can extend to region where T is timelike, I can appeal to analyticity, or the result on the slide.

For 1st result, we want the space to be strongly gravitating, but we want double null foliation to exist. Is it difficult to avoid caustics, but still strongly gravitating enough to make trapped surfaces? No. The existence part controls this, and it is hard, but past that, it is not difficult.

ON THE FORMATION OF TRAPPED **SURFACES**

Sergiu Klainerman

Princeton University

November 19, 2013

- **a.** RIGIDITY. Does the Kerr family $\mathcal{K}(a, m)$, $0 \le a \le m$, exhaust all possible, stationary, asymptotically flat, vacuum black holes ?
- b. STABILITY. Is the Kerr family stable under arbitrary small perturbations ?
- c. COLLAPSE. Can black holes form starting from reasonable initial data configurations ? Formation of trapped surfaces.

RIGIDITY.

STABILITY.

COLLAPSE. Can black holes form starting from reasonable initial data configurations ? Formation of trapped surfaces.

- Penrose Singularity Theorem (1969)
- Christodoulou's Trapping Theorem(2008)
- Kl-Rodnianski(2010)
- Kl-Luk-Rodnianski(2013)

PENROSE SINGULARITY THEOREM

THEOREM. Space-time (M, g) cannot be future null geodesicaly complete, if

- Ric(g)(L, L) \geq 0, \forall L null
- M contains a non-compact Cauchy hypersurface
- M contains a closed trapped surface S

Null second fundamental forms χ, χ

QUESTIONS

- Quantitative version of the incompleteness theorem ?
- Can trapped surfaces form in evolution ? In vacuum ?
	-
- Significance of the uniformity condition?
	-
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- Can trapped surfaces form in evolution ? In vacuum ? Does the existence of a trapped surface implies the presence of a Black Hole ?

True if weak cosmic censorship holds true.

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True if weak cosmic censorship holds true.

• Significance of the uniformity condition?

Can singularities form starting with non-isotropic, initial configurations?

PROBLEM. Specify an open set of regular initial conditions, free of trapped surfaces, on a space-like or null configuration whose future development contains a trapped surface.

- Heuristics ? In the absence of spherical symmetry it is not at
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- **Heuristics ?** In the absence of spherical symmetry it is not at all clear how such a large, uniform distortion can be produced.
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DIFFICULTIES

- **Heuristics ?** In the absence of spherical symmetry it is not at all clear how such a large, **uniform** distortion can be produced.
- **Semi-Global.** Need to control the MFGHD of an initial data for a far longer time than that provided by the classical existence results [Y. C. Bruhat(1952)), Rendall(1990)]

THEOREM[[Christ(2008)]. (∃) open set of regular, vacuum, data whose MGFHD contains a trapped surface.

-
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- ³ Similar result for data given at past null infinity.

THEOREM[Kl-Luk-Rodnianski(2013)] Result holds true for non-isotropic data concentrated near one null geodesic generator.

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THEOREM[Kl-Luk-Rodnianski(2013)] Result holds true for non-isotropic data concentrated near one null geodesic generator.

- **4** Combines all ingredients in Christodoulou's theorem with a deformation argument along incoming null hypersurfaces.
- 2 Reduces to a simple differential inequality on $\ S_{0,0}=H_0\cap \underline{H}_0.$

Characteristic Data. Ric(g) = 0. Hypersurfaces $H_0 \cup \underline{H}_0 \subset \mathbb{R}^{1+3}$, $S_{0,0} = H_0 \cap \underline{H}_0$.

• Conformal metrics $([\gamma]_u, S_u)$, $([\gamma]_u, S_u)$.

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THEOREM[Rendall]. For any given, regular, characteristic initial data set there exists a unique future development near $S_{0,0}$.

THEOREM[Rendall]. For any given, regular, characteristic initial data set there exists a unique, **future** development near $S_{0,0}$, foliated by a double null foliation (u, u) .

D Max. future development $\mathcal{D}(u_*, \underline{u}_*)$, $0 \le u \le u_*, 0 \le \underline{u} \le \underline{u}_*.$ **2** Instead of $[\gamma]_u$, $[\gamma]_{\underline{u}}$ one can prescribe the shears $\hat{\chi}_0$, $\hat{\chi}_0$.

THEOREM[Chr. 2008]. Specify δ -regular initial data on $H_0 \cup H_0$ whose future development $\mathcal{D}(u_*,\delta)$ contains a trapped surface.

-
-
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REMARKS.

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- **3** Companion result at past null infinity

INITIAL DATA

• Flat on
$$
H_0(\mu = 0)
$$
.

 \circ δ - pulse on $H_0(u = 0)$.

Christodoulou's δ -pulse.

 $\hat{\chi}_0(\underline{u},\omega) = \delta^{-1/2}h(\underline{u}/\delta,\omega).$

Definition[Christodoulou's $C(\delta, B)$ -data]

$$
\sum_{i\leq 5}\sum_{k\leq 3}\delta^{\frac{1}{2}+k}||\nabla_{\underline{u}}^k\nabla^i\hat{\chi}_0||_{L^\infty(H_0)}\leq B.
$$

$$
\sum_{0 \le k \le 2} \delta^{1/2} \| (\delta \nabla_4)^k \hat{\chi}_0 \|_{L^2(H_0)} \le B
$$

$$
\sum_{0 \le k \le 1} \sum_{1 \le m \le 4} \delta^{1/2} \| (\delta^{1/2} \nabla)^{m-1} (\delta \nabla_4)^k \nabla \hat{\chi}_0 \|_{L^2(H_0)} \le B
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Definition [KI-Rodn's $KR(\delta, B)$ -data]

$$
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$$

$$
\sum_{0 \le k \le 1} \sum_{1 \le m \le 4} \delta^{1/2} \| (\delta^{1/2} \nabla)^{m-1} (\delta \nabla_4)^k \nabla \hat{\chi}_0 \|_{L^2(H_0)} \le B
$$

$$
\delta^{1/2} \|\hat{\chi}_0\|_{L^\infty(H_0)} \leq B
$$

THEOREM[Christ(2008), Kl-Rodn(2010)]

1 Given $B > 0$, $u_* < 1$ there exists $\delta > 0$ sufficiently small such that if the data verify $C(\delta, B)$ or $KR(\delta, B) \Rightarrow$ the maximal future development contains $\mathcal{D}(u_*,\delta)$.

1 Given $B > 0$, $u_* < 1$, $\exists \delta > 0$ small s.t. if data verify $C(\delta, B)$ or $KR(\delta, B) \Rightarrow$ the maximal future development $\supset \mathcal{D}(u_*, \delta)$.

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surface $S_{u_{*},\delta} \subset \mathcal{D}(u_{*},\delta)$ is trapped.

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³ A companion result can be proved for formation of scars.

- **1** Based on a continuity argument, one derives detailed information of all connection coefficients and curvature components with respect to an adapted null frame.
-

$$
\frac{d}{d\underline{u}} \text{tr}\,\chi = -|\hat{\chi}|^2 + \text{err}(\delta)
$$

$$
\frac{d}{du}|\hat{\chi}|^2 - \frac{2}{r}|\hat{\chi}|^2 = \text{err}(\delta)
$$

$$
\operatorname{tr}\chi(u,\underline{u},\omega) = \frac{2}{r(u,0)} - \frac{1}{r^2(u,0)}\int_0^{\underline{u}} |\hat{\chi}_0|^2(\underline{u}',\omega)d\underline{u}' + O(\delta)
$$

- **1** Based on a continuity argument, one derives detailed information of all connection coefficients and curvature components with respect to an adapted null frame.
- **2** Transport equations

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$$
\n
$$
\text{Need } \qquad \qquad \left| \frac{\overline{2(r_0 - u)}}{r_0^2} < \inf_{\omega \in \mathbb{S}^2} M_0(\omega) \right| < \frac{2}{r_0}
$$

Theorem[Kl-Luk-Rodn(2013)] Assume $C(\delta, B)$. If in addition,

$$
\sup_{\omega \in S_{0,0}} M_0(\omega) \geq M_* > 0.
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then, for any, $0 < u_* < 1$, $1 - u_*$ small, $\exists \delta = \delta(B, M_*) > 0$ such that $\mathcal{D}(u_*,\delta)$ contains a trapped surface.

- Uses all ingredients in Christodoulou's theorem combined with a new deformation argument along $H_\delta.$
-

$$
1 - u_* < R < 1
$$
\n
$$
-\Delta R + R^{-1} |\nabla R|^2 + R < 2^{-1} M_0
$$

- Such functions exist if $M_0 > 0$.
-
- Uses all ingredients in Christodoulou's theorem combined with a new deformation argument along $H_\delta.$
- Surface $\underline{\boldsymbol{\mu}} = \delta, \boldsymbol{r} = R(\omega) \quad [\ R \in \mathcal{C}^2(\mathcal{S}_0)]$ is trapped if

$$
\begin{array}{rcl} & 1-u_* & <\ \ R < 1 \\ -\Delta R + R^{-1} |\nabla R|^2 + R & <\ \ 2^{-1} M_0 \end{array}
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DEFORMED TRAPPED SURFACE

RIGIDITY.

Does the Kerr family $\mathcal{K}(a, m)$, $0 \le a \le m$, exhaust all possible, stationary, asymptotically flat, vacuum black holes ?

STABILITY.

COLLAPSE.

STATIONARY BLACK HOLES

Stationary, asymptotically flat, solutions of the EVE,

$\text{Ric}(g) = 0.$

Asymptoticaly flat, globally hyperbolic, Lorentzian manifold

 \bullet Metric g has an asymptotically timelike, Killing vectorfield T,

$$
\mathcal{L}_T g = 0.
$$

STATIONARY BLACK HOLES

Stationary, asymptotically flat, solutions of the EVE,

 $\text{Ric}(g) = 0.$

DEFINITION [External Black Hole]

Asymptoticaly flat, globally hyperbolic, Lorentzian manifold with boundary (M, g) , diffeomorphic to the complement of a cylinder $\subset \mathbb{R}^{1+3}$.

 \bullet Metric g has an asymptotically timelike, Killing vectorfield τ ,

$$
\mathcal{L}_\mathcal{T}\mathcal{g}=0.
$$

KERR FAMILY $\mathcal{K}(a, m)$

Boyer-Lindquist (t, r, θ, φ) coordinates.

$$
-\frac{\rho^2\Delta}{\Sigma^2}(dt)^2+\frac{\Sigma^2(\sin\theta)^2}{\rho^2}\Big(d\varphi-\frac{2amr}{\Sigma^2}dt\Big)^2+\frac{\rho^2}{\Delta}(dr)^2+\rho^2(d\theta)^2,
$$

$$
\begin{cases}\n\Delta = r^2 + a^2 - 2mr; \\
\rho^2 = r^2 + a^2(\cos\theta)^2; \\
\Sigma^2 = (r^2 + a^2)^2 - a^2(\sin\theta)^2\Delta.\n\end{cases}
$$
\n**Stationary.**\n
$$
\mathbf{T} = \partial_t
$$
\n**Axisymmetric.**\n
$$
\mathbf{Z} = \partial_\varphi
$$

$$
-\frac{\Delta}{r^2}(dt)^2+\frac{r^2}{\Delta}(dr)^2+r^2d\sigma_{\mathbb{S}^2}
$$

KERR FAMILY $\overline{\mathcal{K}}(a, m)$

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$$

Stationary. $\mathbf{T} = \partial_t$ Axisymmetric. $\mathbf{Z} = \partial_{\omega}$

Schwarzschild. $a = 0, m > 0$, static, spherically symmetric.

$$
-\frac{\Delta}{r^2}(dt)^2+\frac{r^2}{\Delta}(dr)^2+r^2d\sigma_{\mathbb{S}^2}
$$

RIGIDITY CONJECTURE. Kerr family $K(a, m)$, $0 \le a \le m$, exhaust all stationary, asymptotically flat, regular vacuum black holes.

- True in the static case. [Israel, Bunting-Masood ul Ulam]
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- T-pseudoconvexity

NEW RIGIDITY RESULT

THEOREM. Assume (M, g) nondegenerate, regular, black hole with a stationary Killing v-field $\;\mathsf T$ and $\;\|\mathsf g(\mathsf T,\mathsf T)\|_{L^\infty(\mathsf S_0)}$ small. \Rightarrow (M, g) is stationary, axially symmetric, hence a $\mathcal{K}(a, m)$ with small a.

MAIN STEPS

- \bigcirc (M, g) admits a a second rotational Killing vector-field Z which commutes with T , in a small neighborhood Ω of the horizon.
- $\textbf{2}$ If $\| \textbf{g}(\textsf{T},\textsf{T}) \|_{L^{\infty}(\textsf{S}_{0})}$ sufficiently small $\Rightarrow \textsf{T}$ is strictly time-like in the complement of $Ω$.

T BECOMES TIMELIKE !

Assume. Maximal hypersurface Σ_0 , passing through bifurcate sphere S_0 . Decompose **T**,

$$
T = nT_0 + X
$$
, $g(T,T) = -n^2 + |X|^2$

$$
\nabla_i X_j + \nabla_j X_i = 2nk_{ij}.
$$

\n
$$
\Delta n = |k|^2 n,
$$

Remark. $n|_{S_0} = 0$, $n = 1$ at infinity.

$$
\nabla_i X_j + \nabla_j X_i = 2nk_{ij}.
$$

$$
\Delta n = |k|^2 n,
$$

Step 1. By Hopf Lemma $\nu(n) > 0$ in a quantitative fashion. **Step 2.** $\int_{\Sigma_0} n|k|^2$ is small. Integrate the identity $n|k|^2 = k^{ij}nk_{ij} = \nabla^i X^j k_{ij} = \nabla^i (X^j k_{ij}).$ **Step 3.** Use Steps 1,2 to show that k is uniformly small. **Step 4.** Show, by a propagation argument, that X remains small. Step 5. Deduce that T becomes strictly time-like away from a small neighborhood of S_0 .

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Pseudo-convexity

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 ${\bf D}^2 f(X,X)(p) < 0.$