

Quantum ergodicity

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11 mai 2015

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- Quantum ergodicity on manifolds (comparing < 0 curvature, > 0 curvature and 0 curvature).
- QE on large regular graphs.

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M a compact riemannian manifold, of dimension d.

$$\Delta \psi_k = -\lambda_k \psi_k$$
$$\|\psi_k\|_{L^2(M)} = 1$$

in the limit $\lambda_k \longrightarrow +\infty$.

We study the weak limits of the probability measures on M,

$$|\psi_k(x)|^2 d\operatorname{Vol}(x)$$

 $\lambda_k \longrightarrow +\infty.$

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This question is linked with the ergodic theory for the geodesic flow / billiard flow.

Hence the name quantum ergodicity.

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Let $(\psi_k)_{k\in\mathbb{N}}$ be an orthonormal basis of $L^2(M)$, with

$$-\Delta\psi_k = \lambda_k\psi_k, \qquad \lambda_k \le \lambda_{k+1}.$$

QE theorem (simplified) :

Theorem (Shnirelman, Zelditch, Colin de Verdière)

Assume that the action of the geodesic flow is **ergodic** for the Liouville measure. Let $a \in C^0(M)$. Then

$$\frac{1}{N(\lambda)}\sum_{\lambda_k\leq\lambda}\left|\int_M a(x)|\psi_k(x)|^2d\mathrm{Vol}(x)-\int_M a(x)d\mathrm{Vol}(x)\right|^2\longrightarrow 0.$$

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Equivalently, there exists a subset $\mathcal{S} \subset \mathbb{N}$ of density 1, such that

$$\int_{M} a(x) |\psi_k(x)|^2 d\operatorname{Vol}(x) \xrightarrow{k \in S} \int_{M} a(x) d\operatorname{Vol}(x).$$

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Equivalently,

$$|\psi_k(x)|^2 d\operatorname{Vol}(x) \xrightarrow[k \to +\infty]{k \to +\infty} d\operatorname{Vol}(x)$$

in the weak topology.

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The full statement uses analysis on phase space, i.e. $T^*M = \{(x, \xi), x \in M, \xi \in T^*_x M\}.$

For $a = a(x, \xi)$ a "reasonable" function on T^*M , we can define an operator on $L^2(M)$,

$$a(x, D_x)$$

Say $a \in S^0(T^*M)$ if a is smooth and 0-homogeneous in ξ (i.e. a is a smooth fn on the sphere bundle).

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$$-\Delta\psi_k = \lambda_k\psi_k, \qquad \lambda_k \leq \lambda_{k+1}.$$

For $a \in S^0(T^*M)$, we consider

 $\langle \psi_k, a(x, D_x)\psi_k \rangle_{L^2(M)}.$

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This amounts to $\int_M a(x) |\psi_k(x)|^2 d \operatorname{Vol}(x)$ if a = a(x).

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Assume that the action of the geodesic flow is **ergodic** for the Liouville measure. Let $a(x,\xi) \in S^0(T^*M)$. Then

$$\frac{1}{N(\lambda)}\sum_{\lambda_k\leq\lambda}\left|\langle\psi_k,a(x,D_x)\psi_k\rangle_{L^2(M)}-\int_{|\xi|=1}a(x,\xi)dxd\xi\right|^2\longrightarrow 0.$$

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For any bounded operator K on $L^2(M)$, define the quantum variance

$$Var_{\lambda}(K) = \frac{1}{N(\lambda)} \sum_{\lambda_k \leq \lambda} \left| \langle \psi_k, K \psi_k \rangle_{L^2(M)} \right|^2$$

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The proof start from the trivial observation that

 $Var_{\lambda}([\sqrt{-\Delta}, K]) = 0$

for any K.

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$$Var_{\lambda}([\sqrt{-\Delta},K])=0$$

for any K.

In addition, if $K = a(x, D_x)$ is a pseudodifferential operator with $a \in S^0(T^*M)$, then

$$[\sqrt{-\Delta}, a(x, D_x)] = (Xa)(x, D_x) + b(x, D_x)$$

where b is -1-homogeneous in ξ and X is the derivation along the geodesic flow.



This implies that

 $Var_{\lambda}((Xa)(x, D_x)) \underset{\lambda \longrightarrow +\infty}{\longrightarrow} 0.$

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In addition,

$$Var_{\lambda}(a(x, D_x)) \leq C \int_{|\xi|=1} |a(x, \xi)|^2 dx d\xi.$$



This implies that

$$Var_{\lambda}((Xa)(x, D_x)) \xrightarrow{\lambda \longrightarrow +\infty} 0.$$

In addition,

$$Var_{\lambda}(a(x,D_x)) \leq C \int_{|\xi|=1} |a(x,\xi)|^2 dx d\xi.$$

If the geodesic flow is ergodic, this implies

$$Var_{\lambda}(a(x, D_x)) \underset{\lambda \longrightarrow +\infty}{\longrightarrow} 0$$

if a has zero mean.

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QUE conjecture :

Conjecture (Rudnick, Sarnak 94)

On a negatively curved manifold, we have convergence of the whole sequence : $\langle \psi_k, a(x, D_x)\psi_k \rangle_{L^2(M)} \longrightarrow \int_{|\xi|=1} a(x, \xi) dx d\xi$.

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Proven by E. Lindenstrauss in the special case of arithmetic congruence surfaces, for joint eigenfunctions of the Laplacian and the Hecke operators.

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A-Nonnenmacher (06) proved a weaker statement valid in greater generality.

Let M have negative curvature. Assume

$$\langle \psi_k, a(x, D_x) \psi_k \rangle_{L^2(M)} \longrightarrow \int_{|\xi|=1} a(x, \xi) d\mu(x, \xi)$$

Then μ must have positive Kolmogorov-Sinai entropy.

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(Jakobson-Bourgain 97, Jaffard 90, A-Macià 2012) It's not possible for a sequence of eigenfunctions to concentrate on a closed geodesic.

On each cylinder of periodic orbits, the limit measure must be absolutely continuous.

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Since the 90s there has been the idea of using graphs as a testing ground/toy model for quantum chaos.

Smilansky, Kottos, Alon,... Keating, Berkolaiko, Winn, Piotet, Marklof...

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Here we focus on the case of <u>large regular (discrete) graphs</u>. Let $G_N = (V_N, E_N)$ be a (q + 1)-regular graph of size N $(V_N = \{1, ..., N\})$.

We look at the limit $N \longrightarrow +\infty$.

We assume that G_N has "few" short loops (= converges to a tree in the sense of Benjamini-Schramm).

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Theorem

(A-Le Masson, 2013) Assume that G_N has "few" short loops and that it forms an <u>expander</u> family. Let $(\phi_i^{(N)})_{i=1}^N$ be an ONB of eigenfunctions of the laplacian on G_N . Let $a = a_N : V_N \longrightarrow \mathbb{C}$ be such that $|a(x)| \leq 1$ for all $x \in V_N$.

Then

$$\lim_{N\longrightarrow +\infty} \frac{1}{N} \sum_{i=1}^{N} \left| \sum_{x \in V_N} a(x) |\phi_i^{(N)}(x)|^2 - \bar{a} \right|^2 = 0.$$

• Also works on shrinking spectral intervals

• Applies to <u>random regular graphs</u>. In that case there also exists a probabilistic proof (Geisinger 2013) in the case where a(x) is chosen independently of G_N .

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More general version

Theorem

(A-Le Masson, 2013) Assume that G_N has "few" short loops and that it forms an <u>expander</u> family. Let $(\phi_i^{(N)})_{i=1}^N$ be an ONB of eigenfunctions of the laplacian on G_N . Let $K_N : V_N \times V_N \longrightarrow \mathbb{C}$ be a matrix such that $d(x, y) > D \Longrightarrow K_N(x, y) = 0$. Assume $|K_N(x, y)| \le 1$. Then $\lim_{N \longrightarrow +\infty} \frac{1}{N} \sum_{i=1}^N \left| \langle \phi_i^{(N)}, K_N \phi_i^{(N)} \rangle - \overline{K_N}(\lambda_i) \right|^2 = 0.$

$$\overline{K_N}(\lambda_i) = \sum_{x,y} K(x,y) \Phi_{sph,\lambda_i}(d(x,y)).$$

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Theorem

(Brooks-Lindenstrauss 2011) Assume that G_N has "few" loops of length $\leq c \log N$. For $\epsilon > 0$, there exists $\delta > 0$ s.t. for every eigenfunction ϕ ,

$$B \subset V_N, \sum_{x \in B} |\phi(x)|^2 \ge \epsilon \Longrightarrow |B| \ge N^{\delta}.$$

Proof also yields that $\|\phi\|_{\infty} \leq |\log N|^{-1/4}$.

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G = (V, E) graph, |V| = N, $\mathcal{A} : \ell^2(V) \longrightarrow \ell^2(V)$ (self-adjoint) defined by

$$\mathcal{A}f(x)=\sum_{y\sim x}f(y).$$

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Introduction. The Shnirelman theorem. Quantum unique ergodicity? Other geometries QE on discrete graphs

Sketch of proof : 1) work with non-backtracking RW instead of simple RW

G = (V, E) graph, |V| = N, $\mathcal{A} : \ell^2(V) \longrightarrow \ell^2(V)$ (self-adjoint) defined by

$$\mathcal{A}f(x)=\sum_{y\sim x}f(y).$$

Define B = set or oriented edges of G, and $\mathcal{B} : \ell^2(B) \longrightarrow \ell^2(B)$ by

$$\mathcal{B}f(e) = \sum_{o(e')=t(e), e' \neq \hat{e}} f(e').$$

Note that \mathcal{B} is not self-adjoint but $\mathcal{B} = I\mathcal{B}^*I$ where I is the edge-reversal involution

$$If(e) = f(\hat{e}).$$

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For regular graphs, the spectrum and eigenfunctions of \mathcal{B} are explicit in terms of those of \mathcal{A} .

• each eigenvalue $\lambda = 2\sqrt{q} \cos(s \ln q)$ ($s \in \mathbb{R} \cup i\mathbb{R}$) of \mathcal{A} gives rise to two eigenvalues $q^{1/2 \pm is}$ of \mathcal{B} .

• the N(q-1) other eigenvalues are ± 1 , each with multiplicity $\frac{N(q-1)}{2}$ (rank of fundamental group of G-1).

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 each eigenvalue λ = 2√q cos(s ln q) (s ∈ ℝ ∪ iℝ) of A gives rise to two eigenvalues q^{1/2±is} of B. The eigenfunction φ_s for A gives rise to the two eigenfunctions of B,

$$f_s^\pm(e)=\phi(t(e))-rac{1}{q^{1/2\pm is}}\phi(o(e)).$$

• the N(q-1) other eigenvalues are ± 1 , each with multiplicity $\frac{N(q-1)}{2}$ (rank of fundamental group of G - 1).

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Introduction. The Shnirelman theorem. Quantum unique ergodicity? Other geometries QE on discrete graphs 2) Definition of the quantum variance for the

2) Definition of the quantum variance for the non-backtracking operator

Consider $K(e, e') : B \times B \longrightarrow \mathbb{C}$ with the property that

$$K(e,e') \neq 0 \Longrightarrow \exists k \leq D, \mathcal{B}^k(e,e') \neq 0.$$

(K may be seen as a function on the set of geodesic segments of length $\leq D$).

Define

$$Var(K) = rac{1}{N} \sum_{j=1}^{N} |\langle If_{s_j}^-, Kf_{s_j}^+ \rangle|^2.$$

This is built so that

$$Var([\mathcal{B}, K]) = 0$$

for all K.

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Introduction.The Shnirelman theorem.Quantum unique ergodicity?Other geometriesQE on discrete graphs3) Dynamical interpretation

Notice that

$$[\mathcal{B},K]=dK$$

("derivative along geodesic flow") and that

$$Var(K) \leq C rac{1}{N} \sum_{e,e'} |K(e,e')|^2$$

if $D \leq girth$.

Uniform mixing of \mathcal{B} on a family of graphs (=expanding property) then implies that

$$Var(K_N) \xrightarrow[N \to \infty]{} 0$$

if $\operatorname{Tr}(K_N \mathcal{B}^k) = 0$ for all k.

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This method seems adaptable to

- quotients (with large girth) of $(F_d, \langle a_1, \ldots, a_d \rangle)$ with weights $p(x, xa_i) = p(a_i)$ symmetric.
- $\Delta + v$, with v deterministic, having some kind of periodicity; probably also v(x) random iid (Anderson model).
- some non-regular graphs??

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For a lot of graphs (not only regular ones) there is an explicit way to transform solutions of

$$(\Delta + \mathbf{v})\phi = \lambda\phi$$

in $\ell^2(V)$ to solutions of $\mathcal{B}f = \alpha_{\lambda}f$ in $\ell^2(B)$, where α_{λ} is a function on B.

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Perspectives

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in $\ell^2(V)$ to solutions of $\mathcal{B}f = \alpha_{\lambda}f$ in $\ell^2(B)$, where α_{λ} is a function on B.

When you have a family of graphs and you need to control the behaviour of the functions α_{λ} ... <□ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ >