# <span id="page-0-0"></span>Random versus Deterministic Approach in the Study of Wave and Dispersive Equations

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## <span id="page-2-0"></span>Introduction

- Lots of progress in the last 20 years in the study of nonlinear dispersive and wave equations.
- The thrust of this body of work has focused on deterministic aspects of wave phenomena.
- Yet there remain some important open questions especially in the supercritical case.

Consider for example the Cauchy IVP for the p-NLS equation:

 $\int$  *iu*<sub>t</sub> + ∆*u* = ±|*u*|<sup>*p*−1</sup>*u*,  $u(x, 0) = u_0(x) \in H^s$  *x*  $\in \mathbb{R}^n$  or  $\mathbb{T}^n$ 

**Recall:** the scale invariant norm is  $s_c := \frac{n}{2} - \frac{2}{(p-1)}$ .

*H*<sup>s</sup> data with  $s > s_c$  is subcritical;  $s = s_c$  is critical;  $s < s_c$  is supercritical.

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#### **Critical Data Results** :

- $\blacktriangleright$  Global well-posedness and scattering for energy-critical NLS in  $\mathbb{R}^n$ 
	- <sup>F</sup> *Defocusing*: Bourgain; Grillakis; Colliander-Keel-S.-Takaoka-Tao; Killip-Visan, X. Zhang, Dodson.
	- \* Focusing: Kenig-Merle (concentrated compactness /rigidity method) and Killip-Visan.
- $\blacktriangleright$  Global well-posedness and scattering for mass-critical NLS in  $\mathbb{R}^n$ 
	- **★ radial: Killip, Tao, Visan, X. Zhang.**
	- <sup>F</sup> *nonradial*: Dodson
- $\triangleright$  Global well-posedness and scattering for 'energy-supercritical' ( $s_c > 1$ ) defocusing NLW and NLS **under the assumption of a uniform in time bound on the scale invariant norm** by Kenig and Merle; Killip and Visan; Bulut.
	- $\star$  In spirit of Escauriaza, Seregin and Sverak for the Navier-Stokes equation.

#### **Supercritical Data Results:** (?)

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# <span id="page-4-0"></span>p-NLS: Deterministic GWP Results on T *n*

#### **Critical Data Results**:

- $\triangleright$  Global well-posedness for energy-critical NLS
	- $\star$  *Defocusing and n* = 3: lonescu-Pausader (large data, based on a work by Ionescu-Pausader-S.); and previously Herr-Tzvetkov-Tataru (small data).

- $\triangleright$  Global well-posedness for mass-critical NLS
	- \* (?) In fact there are no even local results at the L<sup>2</sup> level!

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### <span id="page-5-0"></span>Deterministic  $\rightarrow$  Nondeterministic Approach

Bourgain considered the L<sup>2</sup>-critical<sup>1</sup>:

Theorem (Rational Torus; **Bourgain**(96'))

 $\int$  *iu*<sub>t</sub> +  $\Delta u = |u|^2 u - (\int |u|^2 dx)$  *u* 

Given  $\mu$  a probability measure on the space of initial data X (eg.  $X = H^s$ )

<sup>1</sup>In 93' Bourgain had proved LWP [for](#page-6-0) $s > 0$  and GWP in  $H^1(\mathbb{T}^2)$  $H^1(\mathbb{T}^2)$  $H^1(\mathbb{T}^2)$  for [c](#page-4-0)[ub](#page-7-0)[ic](#page-8-0) [N](#page-2-0)[L](#page-8-0)[S](#page-9-0)  $\Omega$ 

### <span id="page-6-0"></span>Deterministic  $\rightarrow$  Nondeterministic Approach

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Theorem (Rational Torus; **Bourgain**(96'))

 $\int$  *iu*<sub>t</sub> + ∆*u* = |*u*|<sup>2</sup>*u* − ( $\int$ |*u*|<sup>2</sup>*dx*) *u*  $u(x, 0) = u_0(x), \quad x \in \mathbb{T}^2$  (Rational),

*is almost sure globally well-posed* **below** *L* 2 *; i.e. for* **supercritical data**  $u_0 \in H^{-\varepsilon}$ .

Given  $\mu$  a probability measure on the space of initial data X (eg.  $X = H^s$ )

<sup>1</sup>In 93' Bourgain had proved LWP [for](#page-7-0) $s > 0$  and GWP in  $H^1(\mathbb{T}^2)$  $H^1(\mathbb{T}^2)$  $H^1(\mathbb{T}^2)$  for [c](#page-4-0)[u](#page-5-0)[b](#page-7-0)[ic](#page-8-0) [N](#page-2-0)[L](#page-8-0)[S](#page-9-0)  $\Omega$ 

## <span id="page-7-0"></span>Deterministic  $\rightarrow$  Nondeterministic Approach

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*is almost sure globally well-posed* **below** *L* 2 *; i.e. for* **supercritical data**  $u_0 \in H^{-\varepsilon}$ .

#### **Very informal definition of almost sure well-posedness**

Given  $\mu$  a probability measure on the space of initial data X (eg.  $X = H^s$ )

There exists *Y*  $\subset$  *X*, with  $\mu$ (*Y*) = 1 and such that for any  $u_0 \in Y$  there exist  $T > 0$  and a unique solution *u* to the IVP in  $C([0, T], X)$  that is also stable in the appropriate topology.

<sup>1</sup>In 93' Bourgain had proved LWP [for](#page-8-0) $s > 0$  and GWP in  $H^1(\mathbb{T}^2)$  $H^1(\mathbb{T}^2)$  $H^1(\mathbb{T}^2)$  for [c](#page-4-0)[u](#page-5-0)[bic](#page-8-0) [N](#page-2-0)[L](#page-8-0)[S](#page-9-0)  $\Omega$  <span id="page-8-0"></span>Bourgain's interest was to construct an **invariant Gibbs measure** derived from the PDE above viewed as an infinite dimension Hamiltonian system $^2\colon$ 

- <sup>1</sup> Established local well posedness for 'typical elements' in the support of the measure; i.e. for random data  $\phi^\omega$  in  $H^{-\varepsilon}(\mathbb{T}^2)$ , ( an 'almost sure' -in the sense of probability- LWP in  $H^{-\varepsilon}(\mathbb{T}^2)$ ).
- <sup>2</sup> Proved that the associated Gibbs measure is **invariant** and used it to extend the local result to a global one in the almost sure sense.

Furthermore, Bourgain showed that almost surely in  $\omega$  the nonlinear part of the solution

 $w := u - S(t) \phi^{\omega}$ 

is smoother than the linear part.

**Note:** *This result still keeps open the question of (a.s.) global well-posedness when data are in H<sup>s</sup>* , 0 ≤ *s* < 1 *since there are no invariant measures and no conservation laws. More later*.

<sup>2</sup>After Lebowitz, Rose and Speer's and Zhidkov's works.  $A \cup B \cup A \cup B \cup A \cup B \cup A \cup B \cup B$  $QQ$ Gigliola Staffilani (MIT) [Random and Deterministic Approach](#page-0-0) October 19th – 30th, 2015 7/34

## <span id="page-9-0"></span>On Randomized Data

In Bourgain's case, for the cubic NLS on  $\mathbb{T}^2$ , the typical element in the support of the Gibbs measure (the invariant measure) consists of **randomized data**:

$$
\phi^{\omega}(x) = \sum \frac{g_n(\omega)}{|n|} e^{i\langle x, n \rangle} \in H^{-\varepsilon}(\mathbb{T}^2),
$$

where {*gn*(ω)}*<sup>n</sup>* are i.i.d. standard (complex) (Gaussian) random variables on a probability space  $(\Omega, \mathcal{F}, \mathbb{P})$ .

### Remark

*Note that if*

$$
\phi(x)=\sum_{n\in\mathbb{Z}^2}\frac{1}{|n|}e^{i\langle x,n\rangle},
$$

*then*  $\phi \in H^{-\epsilon}$  and also  $\phi^{\omega}(x)$  defines almost surely in  $\omega$  a function in H<sup>- $\epsilon$ </sup> but  $\underline{\text{not}}$  *in*  $H^s$ ,  $s \geq 0$ !

In other words *randomization* does **not** improve regularity in terms of derivatives!

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## <span id="page-10-0"></span> $Randomization = Better estimates$

But there is an *almost sure* improved regularity -akin to the role of Kintchine inequalities in Littlewood-Paley theory- that stems from classical results of **Rademacher, Kolmogorov**, **Paley** and **Zygmund** proving that random series on the torus enjoy better *L <sup>p</sup>* bounds than deterministic ones. For example, consider the *Rademacher Series* :

$$
f(\omega) := \sum_{n=0}^{\infty} a_n r_n(\omega) \qquad \omega \in [0,1), \quad a_n \in \mathbb{C}
$$

where

$$
r_n(\omega):=\operatorname{sign} \sin(2^{n+1}\pi\,\omega)
$$

We have:

If  $a_n \in \ell^2$  the sum  $f(\omega)$  converges a.e. and moreover

### Classical Theorem

If  $a_n \in \ell^2$  then the sum  $f(\omega)$  belongs to  $L^p([0,1))$  for all  $p \geq 2$ . More precisely,

$$
\left(\int_0^1|f|^p\,d\omega\right)^{1/p}\approx_p\|a_n\|_{\ell^2}
$$

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# <span id="page-11-0"></span>Large Deviation-type Estimates

### **Proposition**

*Let* {*gn*(ω)} *be a sequence of complex i.i.d. zero mean Gaussian random variables on a probability space*  $(\Omega, \mathcal{F}, \mathbb{P})$  *and*  $\{c_n\} \in \ell^2$ . Define

$$
\mathsf{F}(\omega):=\sum_n\, \mathsf{c}_n g_n(\omega)
$$

*Then, there exists C* > 0 *such that for every q* ≥ 2 *we have*

$$
\left\|\sum_{n}c_n g_n(\omega)\right\|_{L^q(\Omega)}\leq C\sqrt{q}\left(\sum_{n}|c_n|^2\right)^{\frac{1}{2}}.
$$

As a consequence from Chebyshev's inequality there exists *C* > 0 such that for every  $\lambda > 0$ ,

$$
\mathbb{P}(\{\omega\,:\,|F(\omega)|>\lambda\,\})\,\leq\exp\left(\frac{-C\,\lambda^2}{\|F(\omega)\|_{\!\!\textnormal{L}^2(\Omega)}^2}\right)\!.
$$

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<span id="page-12-0"></span>More generally one uses the following, where *k* would represent the number of random terms in a multilinear estimate at hand,

Proposition (Large Deviation-type) *Let*  $d \geq 1$  *and*  $c(n_1, \ldots, n_k) \in \mathbb{C}$ . Let  $\{(g_n)\}_{1 \leq n \leq d}$  as above. For  $k \geq 1$  denote  $by A(k, d) := \{(n_1, \ldots, n_k) \in \{1, \ldots, d\}^k, n_1 \leq \cdots \leq n_k\}$  and  $\mathcal{F}_k(\omega) = \sum c(n_1, \ldots, n_k) g_{n_1}(\omega) \ldots g_{n_k}(\omega).$ *A*(*k*,*d*) *Then for p*  $>$  2  $||F_k||_{L^p(\Omega)} \lesssim \sqrt{2}$  $\overline{k+1}(p-1)^{\frac{k}{2}}\|F_k\|_{L^2(\Omega)}.$ 

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<span id="page-13-0"></span>**As a consequence from Chebyshev's inequality for every** λ > 0**,**

$$
\mathbb{P}(\{\omega:|F_k(\omega)|>\lambda\})\leq \exp\left(\frac{-C\,\lambda^{\frac{2}{k}}}{\|F_k(\omega)\|_{L^2(\Omega)}^{\frac{2}{k}}}\right).
$$

This result follows from the hyper-contractivity property of the Ornstein-Uhlenbeck semigroup by writing  $G_n = H_n + iL_n$  where  $\{H_1, \ldots, H_d, L_1, \ldots, L_d\}$  are real centered independent Gaussian random variables with the same variance.

#### (**c.f. Tzvetkov; Thomann-Tzvetkov**)

Given  $\delta > 0$ , the large deviation result above with

 $\lambda = \delta^{-\frac{k}{2}}\|F_k(\omega)\|_{L^2(\Omega)}$ 

says that in a set  $\Omega_\delta$  with  $\mathbb{P}(\Omega_\delta^c)< e^{-\frac{1}{\delta}}$  we can replace  $|F_k(\omega)|^2$  by  $||F_k(\omega)||_{L^2(\Omega)}^2$ .

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## <span id="page-14-0"></span>An Example from Bourgain's Work:

Take again

$$
\phi^{\omega}(x) = \sum \frac{g_n(\omega)}{|n|} e^{i\langle x, n \rangle}
$$

and look at the cubic Wick ordered nonlinearity, involving its free evolution  $S(t)\phi^{\omega}(x)$ , and that Bourgain had to estimate in  $L^2$ :

 $\|F_3(\omega)\|_{\mathfrak{h}_m^2\mathfrak{h}_m^2},$ 

where

$$
F_3(\omega)=\sum_{S_{n,m}}\frac{1}{|n_1|}\frac{1}{|n_2|}\frac{1}{|n_3|}g_{n_1}(\omega)\overline{g_{n_2}}(\omega)g_{n_3}(\omega)
$$

where

 $S_{n,m} = \{ (n_1, n_2, n_3) / n_1 - n_2 + n_3 = n; n_1, n_3 \neq n_2; m = |n_1|^2 - |n_2|^2 + |n_3|^2 \}$ 

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Naively we could just use C-S to estimate  $\|F_3(\omega)\|_{\beta,\beta}^2$  and obtain

$$
\sum_{n,m}\left|\sum_{S_{n,m}}\frac{1}{|n_1|}\frac{1}{|n_2|}\frac{1}{|n_3|}g_{n_1}(\omega)\overline{g_{n_2}}(\omega)g_{n_3}(\omega)\right|^2\lesssim \sum_{n,m}|S_{n,m}|\sum_{S_{n,m}}\frac{1}{|n_1|^2}\frac{1}{|n_2|^2}\frac{1}{|n_3|^2}
$$

where  $|S_{n,m}|$  is the cardinality of  $S_{n,m}$  and it translates into a **loss** of derivatives.

Now in the Large Deviation Estimate, take

$$
\lambda = \delta^{-1} ||F_3(\omega)||_{L^2(\Omega)}.
$$

Then in a set  $\Omega_\delta$  of measure 1  $e^{-\frac{1}{\delta^\alpha}}$  one has

$$
||F_3(\omega)||^2_{\mathcal{L}^2_{n,m}} = \sum_{n,m} |F_3(\omega)|^2 \lesssim \delta^{-2} \sum_{n,m} ||F_3(\omega)||^2_{L^2(\Omega)}
$$
  
=  $\delta^{-2} \sum_{n,m} \sum_{S_{n,m}} \sum_{S'_{n,m}} \int_{\Omega} \frac{g_{n_1}}{|n_1|} \frac{\overline{g_{n_2}}}{|n_2|} \frac{g_{n_3}}{|n_3|} \frac{\overline{g_{n'_1}}}{|n'_1|} \frac{\overline{g_{n'_2}}}{|n'_2|} \frac{\overline{g_{n'_3}}}{|n'_3|} d\omega$ 

and by **independence** of the random variables the RHS contracts to

$$
\|F_3(\omega)\|_{l^2_n l^2_m}^2 \lesssim \delta^{-2} \sum_{n,m} \sum_{S_{n,m}} \frac{1}{|n_1|^2} \frac{1}{|n_2|^2} \frac{1}{|n_3|^2}
$$

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# <span id="page-16-0"></span>Randomization without invariant measure (a.s LWP)

In this vein, consider

$$
\begin{array}{lll} \text{(IVP)} & \left\{ \begin{array}{ll} u_t + P(D)u = F(u) & x \in M, \ t > 0 \\ u(x, 0) = \phi(x), \end{array} \right. \end{array}
$$

with  $\phi \in X^s$ , a set of initial data of regularity  $s$  small.

If *M* is a compact manifold of dimension *d*, no boundaries and with a countable basis of eigenvectors  $\{h_n(x)\}$  for the Laplacian, then we randomize  $\phi$  as

$$
\phi^{\omega}(x) := \sum_{n \in \mathbb{Z}^d} a_n g_n(\omega) h_n(x).
$$

If  $M = \mathbb{R}^d$  then we randomize  $\phi$  as

$$
\phi^{\omega}(x):=\sum_{n\in\mathbb{Z}^d}g_n(\omega)P_n\phi(x),
$$

where *P<sup>n</sup>* is a projection operator on cubes of size *one* on the frequency space.

### General procedure to prove a.s LWP

- Assume  $v^{\omega}$  is the **linear evolution** with initial datum  $\phi^{\omega}$ .
- Use the fact that *v* <sup>ω</sup> has better *L <sup>p</sup>* or multilinear estimates than φ *almost surely* to show that  $w = u - v^{\omega}$  solves a **difference equation** that lives in a smoother space than *X s* . Obtain for *w* a *deterministic* local well-posedness.

### Remark (Important)

*The difference equation that w solves is not back to merely being at a 'smoother' level but rather it is a hybrid equation with nonlinearity* = = *supercritical (but random)* + *deterministic (smoother).*

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### <span id="page-18-0"></span>a.s Local to a.s Global: known mechanisms

• Invariant Gibbs or weighted Wiener measures- when available.

The use of the invariance of the **measure** has limitations since in **higher dimensions** its support (data) lives on extremely **rough spaces** where the multilinear analysis needed to control the nonlinear terms of the equation is so far not possible. **In higher dimensions usually a radial assumption is put in place**.

- Sometimes may use energy methods (**Burq-Tzvetkov** and **Pocovnicu** for NLW; **Nahmod-Pavlovic-S.** for NS)
- Sometimes may use adaptation to this setting of Bourgain's *high-low method* (**Colliander- T. Oh**, NLS; **Bulut**, **Luehrmann-Mendelson**, NLW)

#### **These methods have limitations!**

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<span id="page-19-0"></span>• Randomization techniques have now been used with or without the help of the invariant measure in several contexts:

After Bourgain's work in 94-96'; in 07-08 work by Burq-Tzvetkov (NLW, supercritical), T. Oh's (coupled KdV system, subcritical) and Tzvetkov (NLS, subcritical). Lots of work followed:

- <sup>I</sup> Schrodinger Equations ¨ : **Bourgain**, **Tzvetkov**, **Thomann**, **Thomann-Tzvetkov**, **Nahmod-Oh-Rey-Bellet-S.**, **Nahmod-Rey-Bellet-Sheffield-S.**, **Burq-Thomann-Tzevtkov**, **Y. Deng**, **Burq-Lebeau**, **Bourgain-Bulut, Nahmod- S, Poiret-Robert-Thomann, Bényi-Oh-**
- ► KdV Equations: **Bourgain, T. Oh** and **Richards**.
- ▶ NLW Equations: **Burq-Tzvetkov, de Suzzoni, Bourgain-Bulut** and **Luehrmann-Mendelson**. Also see non-squeezing for 3D cubic NLKG **Mendelson**.
- ► Benjamin-Ono Equations: **Y. Deng** and **Y. Deng-Tzvetkov-Visciglia**.
- ► Navier-Stokes Equations: Nahmod-Pavlovic-S. (infinite energy weak solutions on T 3 ). Also work by **C.Deng-Cui** an[d](#page-18-0) **Z[ha](#page-20-0)[n](#page-18-0)[g-](#page-20-0)[F](#page-21-0)[a](#page-13-0)[n](#page-14-0)[g](#page-21-0)**

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- ▶ Schrödinger Equations: **Bourgain, Tzvetkov, Thomann, Thomann-Tzvetkov**, **Nahmod-Oh-Rey-Bellet-S.**, **Nahmod-Rey-Bellet-Sheffield-S.**, **Burq-Thomann-Tzevtkov**, **Y. Deng**, **Burq-Lebeau**, **Bourgain-Bulut**, **Nahmod- S**, **Poiret-Robert-Thomann**, **Benyi- Oh- ´ Pocovnicu**.
- ► KdV Equations: **Bourgain**, **T. Oh** and **Richards**.
- ▶ NLW Equations: **Burg-Tzvetkov, de Suzzoni, Bourgain-Bulut** and **Luehrmann-Mendelson**. Also see non-squeezing for 3D cubic NLKG **Mendelson**.
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#### <span id="page-21-0"></span>To sum up:

- When deterministic statements about existence, uniqueness and stability of solutions to certain evolution equations are **not** feasible/available:
	- $\rightarrow$  turn to a more probabilistic point of view
	- $\rightarrow$  within reach at this time: investigate these problems from a nondeterministic viewpoint; e.g. for random data.

#### Situations when such a point of view is desirable include:

- **o** supercritical regime
- when certain type of illposedness is present,
- when there still remains a gap between local and global wellposedness (subcritical regime relative to the scaling threshold),

- <span id="page-22-0"></span>*Result 1* : A.s. global well-posedness for 2D, cubic defocusing NLS in *H s* (T 2 ), 0 < *s* < 1. **(Nahmod-S.)**
- *Result 2* : Existence of large data global solutions to the 3D quintic NLS for supercritical data in *H* 1− (T 3 ). **(Nahmod-S.)**
- *Result 3* : A.s. global well-posedness for 1D, quintic (small mass) focusing NLS in *H s* (T), 1/2 < *s*. **(Nahmod-S.)**

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## On *Result 1*

### Theorem (Nahmod-S.)

*The 2D cubic defocusing NLS is a.s globally well-posed in*  $H^{s}(\mathbb{T}^{2}), 0 < s < 2/3.$ 

#### **What Was Known:**



- Deterministic methods: l.w.p for  $s > 0$  (**Bourgain**) and g.w.p.  $s > 2/3$ **(Bourgain; De Silva-Pavlovic-S.-Tzirakis)**.
- Methods exploiting data randomization and invariant Gibbs measure  $\mu$ : a.s. global well-posedness in *H* − , **(Bourgain)**.

**Remark:** The theorem is not trivial since any  $\Sigma \subset H^s$ ,  $s > 0$ , is such that for the Gibbs measure  $\mu$  one has  $\mu(\Sigma) = 0$ .  $QQ$ イロト イ母 トイラト イラト

## <span id="page-24-0"></span>Idea of the Proof

Let  $\{g_n(\omega)\}\$  be a sequence of complex i.i.d. zero mean Gaussian random variables on a probability space  $(\Omega, \mathcal{F}, \mathbb{P})$ . Then consider the data

$$
\phi_{\alpha}^{\omega}(x) = \sum \frac{1}{|n|^{\alpha}} \frac{g_n(\omega)}{|n|} e^{i\langle x, n \rangle}
$$

for  $\alpha > 0$ . Note that  $\phi^{\omega}_{\alpha} \in H^s$ ,  $0 < s < \alpha$ .

**Step 1:** Prove that a.s with these data the IVP is globally well-posed in *H*<sup>-ε</sup>.

Pick  $N \gg 1$  and write

$$
\phi_{\alpha}^{\omega}(x) = \sum_{|n| < N} \frac{g_n(\omega)}{|n|^{1+\alpha}} e^{in \cdot x} + \sum_{n \in \mathbb{Z}^2} a_n \frac{g_n(\omega)}{|n|} e^{in \cdot x} =: w_N(x) + \psi_1^{\omega}(x)
$$

where  $\|w_N\|_{H^{\epsilon}} < A$  and

$$
a_n = \begin{cases} 0 & \text{if} \quad |n| < N \\ \frac{1}{|n|^{\alpha}} & \text{if} \quad |n| \ge N, \end{cases} \qquad |a_n| \lesssim \frac{1}{N^{\alpha}} \quad \text{ for all } \quad n \in \mathbb{Z}^2.
$$

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Use Bourgain's result in  $H^{-\epsilon}_-$  to claim that there exists a set  $\tilde{\Omega}\subset \Omega$  such that  $\mathbb{P}(\tilde{\Omega}) = 1$  and for any  $\omega \in \tilde{\Omega}$  one has that

$$
\phi_{\beta_N}^{\omega}(x) = \beta_N \sum_{n \in \mathbb{Z}^2} \frac{g_n(\omega)}{|n|} e^{in \cdot x} + w_N(x) =: \psi_2^{\omega}(x) + w_N(x)
$$

where  $\beta_{\bm{N}}=\frac{1}{N^{\alpha}}$  and  $\phi_{\beta_{\bm{N}}}^{\omega}(x)$  evolves globally to a solution  $v_{\bm{N}}$  and  $v_N(t) \in S(t)(\phi_{\beta_N}^{\omega}) + H^{\epsilon}.$ 

Use the global solution  $v_N$  and a *perturbation argument* to prove the existence and uniqueness of the solution *u* for the original IVP

 $u(t) \in S(t)(\phi_\alpha^\omega) + H^\epsilon$ .

The **Key** point is that if we define  $\zeta = v_N - u$ , then  $\zeta$  solves a Schrödinger equation with nonlinearity of type

 $F([v_N + S(t)w_N] + S(t)\psi_1^{\omega} - F([v_N + S(t)w_N] + S(t)\psi_2^{\omega} - \zeta),$ 

 $\psi_1^{\omega}, \, \psi_2^{\omega}$  $\psi_1^{\omega}, \, \psi_2^{\omega}$  $\psi_1^{\omega}, \, \psi_2^{\omega}$  are *small* a[n](#page-24-0)d  $w_N$  is *uniformly boun[de](#page-24-0)d* in N[.](#page-26-0)

<span id="page-26-0"></span>**Step 2**: Recovery of regularity to claim:

 $u(t) \in S(t)(\phi_\alpha^\omega) + H^{\alpha+\epsilon}.$ 

From Step 1 we know that  $u(x, t) = S(t)(\phi^{\omega}_{\alpha}) + w(x, t)$  with  $||w||_{H^{\epsilon}} \leq A$ . We want to upgrade the regularity of *w* by *D* <sup>α</sup>. By the Duhamel principle we have to estimate

$$
D^{\alpha}w = \int_0^t S(t-t')D^{\alpha}[|S(t')(\phi^{\omega}_{\alpha}) + w|^2(S(t')(\phi^{\omega}_{\alpha}) + w)] dt' \sim \int_0^t S(t-t')[(S(t')(\phi^{\omega}) + D^{\alpha}w)(S(t')(\phi^{\omega}_{\alpha}) + w)^2] dt'
$$

For which the analysis of Bourgain in the random part  $S(t')(\phi^\omega)$  and the fact that *D* <sup>α</sup>*w* only appears **linearly** can be used to conclude.

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## On *Result 2*

We consider the energy-critical quintic nonlinear Schrödinger equation

$$
\begin{cases}\ni u_t + \Delta u = \lambda u |u|^4 & x \in \mathbb{T}^3 \quad \text{(Rational)}\\ \nu(0, x) = \phi(x) & \in H^\gamma(\mathbb{T}^3),\end{cases}
$$

below  $H^1(\mathbb{T}^3)$  (ie. for some  $\gamma < 1)$  and where  $\lambda = \pm 1$ 

- Herr, Tzvetkov and Tataru (10') proved small data global well posedness in  $H^1(\mathbb{T}^3)$ .
- Ionescu and Pausader (12') proved *large data* global well posedness in  $H^1(\mathbb{T}^3)$  in the defocusing case
	- Rely on large data GWP in  $\mathbb{R}^3$  for the energy-critical quintic NLS (by Colliander-Keel-S-Takaoka-Tao (03')).

Our interest is first to establish a local almost sure well posedness for random data *below H*<sup>1</sup>(T<sup>3</sup>) that is in the **supercritical** regime relative to scaling, and then address g.w.p.

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The problem we are considering here is the analogue of the supercritical well-posedness result proved by **Bourgain** for the periodic mass critical cubic NLS in 2D; that is - a.s for data in  $H^{-\epsilon}(\mathbb{T}^2),\, \epsilon>0$  mentioned above.

In our problem we consider data  $\phi \in H^{1-\varepsilon}(\mathbb{T}^{3})$  for any  $\varepsilon > 0$  of the form

$$
\phi(x)=\sum_{n\in\mathbb{Z}^3}\frac{1}{\langle n\rangle^{\frac{5}{2}}}e^{in\cdot x}\quad \xrightarrow{\text{randomization}}\quad \phi^\omega(x)=\sum_{n\in\mathbb{Z}^3}\frac{g_n(\omega)}{\langle n\rangle^{\frac{5}{2}}}e^{in\cdot x}
$$

where  $(g_n(\omega))_{n\in\mathbb{Z}^3}$  is a sequence of complex i.i.d centered Gaussian random variables, as above.

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## Heart of the matter

- Assume *u* solves our IVP, then we define  $w := u S(t) \phi^{\omega}$ , where  $S(t) \phi^{\omega}$  is the linear evolution of the initial profile  $\phi^\omega.$
- We study the IVP for *w* which solves a difference equation with nonlinearity

 $\tilde{N}(w) := |w + S(t) \phi^{\omega}|^4 (w + S(t) \phi^{\omega}).$ 

We expect to prove that *w* belongs to *H s* for some *s* > 1.

- **•** The heart of the matter is to prove multilinear deterministic/random estimates coming from  $N(w)$  to then be able to set up a contraction method to obtain well-posedness.
- $\bullet$  When the NLS equation is considered, **multilinear estimates for**  $\tilde{N}(w)$  can be carried out only after having removed certain *"double frequencies"* involved in the nonlinear part of the equation. In the **cubic** case a Wick ordering of the Hamiltonian was used (see **Bourgain (96'), Colliander-Oh (12')**).

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There are four major complications in the work that we present here compared to the work of Bourgain:

- a quintic nonlinearity increases quite substantially the different cases that needs to be analyzed,
- the counting lemmata in a 3*D* integer lattice are much less favorable than in a 2*D* lattice,
- **the Wick ordering is not sufficient to remove certain bad** *resonant* frequencies.
- We work on [HTT]'s atomic function spaces *X s* , *Y <sup>s</sup>* whose norms are not invariant if one replaces the Fourier transform with its absolute value.

### Theorem (Nahmod-S.)

Let  $\phi^\omega$  as above. Then there exists  $s>1$  and and there exists  $0<\delta_0\ll1$  and  $r = r(s) > 0$  *s.t. for any*  $\delta < \delta_0$ , *there exists*  $\Omega_{\delta}$  *with* 

 $\mathbb{P}(\Omega_{\delta}^{\boldsymbol{c}})<\boldsymbol{e}^{-\frac{1}{\delta^{\boldsymbol{r}}}},$ 

*and for each*  $\omega \in \Omega_{\delta}$  *there exists a unique solution u of the quintic NLS in the space*

 $S(t)\phi^{\omega} + X^{s}([0,\delta))_{d}$ 

*with initial condition*  $φ<sup>ω</sup>$ .

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Extending these solutions globally in time is hard since if there were an invariant Gibbs measure it would be supported in *H*<sup>−1/2−∈</sup>! Other possible routes could be:

- Low-High method of Bourgain and randomization. (See previous work of **Colliander-Oh** and **Luehrmann-Mendelson**). This though is not implementable at the moment because one would need to use the global result in *H* <sup>1</sup> by Ionescu-Pausader where the bounds for the "Strichartz norm" of the solution is super-exponential with respect to the energy.
- **The recent** *conditional* **argument of <b>Bényi- Oh- Pocovnicu** for the NLS.

See also the recent results of **Pocovnicu** and **Oh-Pocovnicu** similar to the one above but for NLW. Here they are able to remove the *conditional assumption* by using a "probabilistic" energy bound on the difference equation.

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### On Global Solutions

If we are interested on claiming that there are "some" large data evolving to solutions for large times, then this can be done. Again assume

$$
\phi^{\omega} = \sum_{n \in \mathbb{Z}^3} \frac{g_n(\omega)}{\langle n \rangle^{\frac{5}{2}}} e^{in \cdot x}.
$$

We have the following theorem:

#### Theorem (Nahmod-S.)

*Let s* > 1 *and* φ <sup>ω</sup> *as above. Fix a large interval of time* [0, *T*]*. Then there*  $\textsf{exists}~0<\delta\sim\mathcal{T}^{-\frac{1}{4}}$  and there exists  $\Omega_\delta$  with

 $\mathbb{P}(\Omega_\delta^\mathcal{c}) < \boldsymbol{e}^{-\delta}$ 

*and for each*  $\omega \in \Omega_{\delta}$  *there exists a unique solution u of the quintic NLS in the space*

$$
S(t)\phi^{\omega}+X^s([0,T))_d,
$$

with initial condition  $\phi^{\omega}$ .

#### Remark

- *This is a large data result.*
- *As T* → ∞ *the size of the set of initial data giving rise to solutions on the whole interval* [0, *T*] *shrinks to zero.*

**Idea of the proof:** It is a combination of an iterated continuity argument and the fact that the random term can be made small via Large Deviation Estimates.

### Remark

**Krieger-Schlag** *considered the septic NLW in* R <sup>3</sup>+<sup>1</sup> *and proved the existence of a class of global smooth solutions with infinite critical norm*  $\dot{H}^{7/6} \times \dot{H}^{1/6}$ *. This is a constructive purely deterministic approach.*

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### Theorem (Nahmod-S.)

*The focusing 1D quintic NLS, with small mass, is a.s globally well-posed in*  $H^{s}(\mathbb{T}^{1}), 1/2 < s.$ 

#### **What Was Known:**

- Deterministic methods, focusing and defocusing: l.w.p for *s* > 0, **(Bourgain)**.
- For defocusing g.w.p for *s* > 4/9, **(Bourgain), (De Silva-Pavlovic-S.-Tzirakis)**.
- Methods exploiting data randomization: a.s. global well-posedness in *H* 1/2− for defocusing and focusing when mass small, **(Bourgain)**.

**Remark:** Also in this case the theorem is not trivial since any  $\Sigma \subset H^s, s > 1/2$ , is such that for the Gibbs measure  $\mu$  one has  $\mu(\Sigma) = 0$ .

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## <span id="page-36-0"></span>Idea of the Proof

Let  ${g_n(\omega)}$  be a sequence of complex i.i.d. zero mean Gaussian random variables on a probability space  $(\Omega, \mathcal{F}, \mathbb{P})$ . Then consider the data

$$
\phi_{\alpha}^{\omega}(x) = \sum \frac{1}{|n|^{\alpha}} \frac{g_n(\omega)}{|n|} e^{i\langle x, n \rangle}
$$

for  $\alpha > 0$ . Note that  $\phi_{\alpha}^{\omega} \in H^{\mathbf{s}},\, \mathbf{0} < \mathbf{s} < 1/2 + \alpha$ .

**Step 1:** Prove that a.s with these data the IVP is globally well-posed in *H*<sup>1/2− $\epsilon$ . In particular show that the solution *u* can be written as</sup>

 $u(x, t) = S(t)\phi_{\alpha}^{\omega}(x) + w(x, t), \quad ||w(t)||_{H^{1/2+\epsilon}} < A, \quad \forall t.$ 

### Remark

*This step doesn't follow from Bourgain's proof since there he uses the deterministic l.w.p result available for s* > 0*. In order to obtain this step one has to repeat the argument for the quintic IVP, (gauge transformation etc). Here the analysis is simpler since the counting lemmata are trivial.*

**• Step 2:** Recovery of regularity.

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