

10/29. Thomas Duyckaerts

Joint with Tristan Roy

$$(NLW) \begin{cases} \partial_t^2 u - \Delta u = |u|^{p-1} u \\ \vec{u}|_{t=0} = (u_0, u_1) \end{cases} \quad x \in \mathbb{R}^3$$

$$p > 5 \quad f = (f, \partial_t f)$$

Locally well-posed in  $\mathcal{H}^{s_c} = \dot{H}^{s_c}(\mathbb{R}^3) \times \dot{H}^{s_c-1}(\mathbb{R}^3)$

$$s_c = \frac{3}{2} - \frac{2}{p-1} \quad \text{conserved energy}$$

$$E(\vec{u}) = \frac{1}{2} \int_{\mathbb{R}^3} |\nabla_{t,x} u|^2 - \frac{1}{p+1} \int_{\mathbb{R}^3} |u|^{p+1}$$

$$|\nabla_{t,x} u|^2 = |\nabla_x u|^2 + |\nabla_t u|^2$$

$$p > 5 \Leftrightarrow s_c > 1 \quad p = 5 \Leftrightarrow s_c = 1$$

Finite speed of propagation.

$u(t, x)$  is determined by  $\vec{u}(0, y)$   $|y-x| < |t|$



- Profile decomposition (Bahouri-Gérard)
- Rigidity classification of non dispersive objects.  
- Merle 1995
- exterior energy estimate (T.D., Kenig, Merle)
- no soliton

II Typical examples of solutions

A Type I blow-up solutions

Def.  $u$  is a type I blow-up solution if  $T_+ < \infty$  and

$$\lim_{t \rightarrow T_+} \| \vec{u}(t) \|_{\mathcal{H}^{s_c}} = +\infty$$

ODE

$$y'' = |y|^{p-1} y$$

$$y(t) = \frac{c_p}{(T-t)^{\frac{2}{p-1}}}$$

Merle and Zaag and  
Dörmayer and Schörkhuber

B Linear equation.

$$(LW) \begin{cases} \partial_t^2 u - \Delta u = 0 \\ \vec{u}|_{t=0} = (u_0, u_1) \in \dot{H}^{sc} \end{cases}$$

Strichartz (Cinibre and Velo '95).

$$(u_0, u_1) \in \dot{H}^{sc} \Rightarrow v \in L^{2(p-1)}(\mathbb{R} \times \mathbb{R}^3)$$

Prop. If  $u$  is a solution of (NLW)

(i)  $u \in L^{2(p-1)}([0, +\infty)_t \times \mathbb{R}^3_x)$

(ii)  $\exists v$  solution of (LW) such that

$$\lim_{t \rightarrow +\infty} \| \vec{u}(t) - \vec{v}(t) \|_{\dot{H}^{sc}} = 0$$

= "scattering".

small data scattering.

C. Elliptic equation

$$(EL) \Delta Q = |Q|^{p-1} Q \quad Q \in \dot{H}^{sc}(\mathbb{R}^3)$$

$$p=5 \quad W = \frac{1}{(1 + \frac{|x|^2}{3})^{3/2}} \sim \frac{1}{|x|} \text{ "only" radial solution except 0.}$$

$p > 5$

no solution in  $\dot{H}^{sc}(\mathbb{R}^3)$  Joseph / Lundgren 70's.

Prop. There exist a solution of (EL) for  $n \neq 0$

such that

•  $Q(n) \sim \frac{1}{n} \quad n \rightarrow \infty$

•  $Q \notin \dot{H}^{sc}(\mathbb{R}^3)$

III. Main Result

$p=5$  (Energy critical case)

"soliton resolution" Any non-type I blow-up solution

decomposes as  $\underset{\substack{\uparrow \\ \text{solution of (LW)}}}{v(t)} + \text{finite sum of rescaled } W$

## Conjecture for $p \geq 5$

Any solution of (NLW) is scattering or Type I blow-up.

### Weaker version

(\*) if  $\limsup_{t \rightarrow +\infty} \|\vec{u}(t)\|_{\dot{H}^{s_c}} < \infty$ , then  $u$  scatters.

(\*) Kenig - Merle (2011) defocusing radial  
Killip - Visan nonradial

TD K M (2013) focusing radial  
higher dimension Bulut (2012)

Dodson Lawrie

• Navier - Stokes: Several works by Sverak, Kenig, Koch  
and coauthors

### Thm

(TD, Tristan Roy) If  $u$  is a solution of (NLW)

(i)  $\lim_{t \rightarrow T_+} \|\vec{u}(t)\|_{\dot{H}^{s_c}} = +\infty$  or

(ii)  $T_+ = +\infty$  and  $u$  scatters.

Krieger - Schlag constructed a global solution, non-scattering  
 $p = 7$   $s_c = \frac{7}{6}$   $\in \dot{H}^s$ ,  $s > \frac{7}{6}$ ;  $\notin \dot{H}^{7/6}$

• Colliot  $D \geq 11$   $p \gg 1$

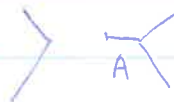
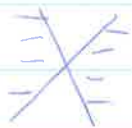
$T < +\infty$   $\dot{H}^s$   $s < s_c$  bounded,  $\dot{H}^{s_c}$  blows up.

### IV. Idea of proof.

A. Exterior energy estimates.

prop 1 Any solution  $v$  of (LW) satisfies  $\forall t \geq 0$  or  $t \leq 0$

$$\int_{|x| > |t|} |\nabla_{t,x} v(t,x)|^2 dx \geq \frac{1}{2} \int |v_0|^2 + v_0'^2$$



Prop 2 If  $v$  is radial and  $A > 0$ , then

$$\int_{A+1}^{\infty} |\partial_r(rv)|^2 + |\partial_t(rv)|^2 dr \geq \frac{1}{2} \int_A^{+\infty} |\partial_{r,t}(rv)|^2|_{t=0}$$

Prop 2  $\Rightarrow$  Prop 1.  $A \rightarrow 0$

$$\int_B^{+\infty} (\partial_r(rv))^2 = \int_B^{+\infty} r^2 (\partial_r v)^2 - Bv^2(B)$$

Prop 1.  $(\partial_t^2 - \partial_r^2)(rv) = 0$

$$rv = f(t+r) - f(t-r)$$

$$(\partial_r(rv))^2 + (\partial_t(rv))^2 = r|f(t+r)|^2 + r|f(t-r)|^2$$

= "modified energy"  $m = \frac{p-1}{2}$

$$E_m(\vec{v}(t)) = \int_0^{\infty} |\partial_r(rv)|^m + |\partial_t(rv)|^m dr$$

Prop

- $E_m(\vec{v}(t)) \approx E_m(\vec{v}(0))$
- $(v_0, v_1) \in \dot{H}^{sc} \Rightarrow E_m$  well-defined.
- exterior energy estimates.
- small data theory.

Assume  $T_+ < \infty$

$\|\vec{u}(t_n)\|_{\dot{H}^{sc}}$  bounded  $t_n \rightarrow T_+$



$\vec{u}(t_n)$  decomposes as sum of profiles.