

MSRI LECTURES ON PSEUDODIFFERENTIAL OPERATORS

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ABSTRACT. Rough notes for lectures at the MSRI introductory workshop in Fall 2019.

A large part of these notes is a shortened version of lecture notes by Richard Melrose, available at <http://math.mit.edu/~rbm/iml190c2.ps>.

1. PROLOGUE: WHY STUDY PSEUDODIFFERENTIAL OPERATORS?

In these two lectures we will mostly study pseudodifferential operators on \mathbb{R}^n . They do also work beautifully well on manifolds.

A very basic example of a pseudodifferential operator is the L^2 inverse to the shifted Laplacian

$$P = \Delta + 1, \quad \Delta := - \sum_{j=1}^n \partial_{x_j}^2.$$

This inverse is a Fourier multiplier (here \mathcal{S} denotes Schwartz functions):

$$P^{-1} = Q : \mathcal{S}(\mathbb{R}^n) \rightarrow \mathcal{S}(\mathbb{R}^n), \quad \widehat{Qu}(\xi) = \frac{1}{1 + |\xi|^2} \hat{u}(\xi).$$

Now, imagine that we instead have a *variable coefficient* operator, e.g.

$$- \sum_{j,k} p_{jk}(x) \partial_{x_j} \partial_{x_k},$$

where p_{jk} is a positive definite matrix depending on x . What would the inverse be? At the end of this lectures we construct an approximate inverse which is a *pseudo-differential operator*. This was the original motivation for studying pseudodifferential operators in the theory of PDE.

Pseudodifferential operators are a general class of operators which include differential operators, Fourier multipliers like Q above, approximate inverses to elliptic differential operators, and a lot more.

2. SPECIAL QUANTIZATION FORMULA

To get the formula for a general pseudodifferential operator, we first look at a differential operator of order m

$$A = \sum_{|\alpha| \leq m} a_\alpha(x) D_x^\alpha, \quad a_\alpha \in C^\infty(\mathbb{R}^n),$$

where we henceforth adopt the notation

$$D_x^\alpha = D_{x_1}^{\alpha_1} \cdots D_{x_n}^{\alpha_n}, \quad D_{x_j} = \frac{1}{i} \partial_{x_j}.$$

Take $u \in \mathcal{S}(\mathbb{R}^n)$ and write by the Fourier inversion formula

$$u(x) = (2\pi)^{-n} \int_{\mathbb{R}^n} e^{ix \cdot \xi} \hat{u}(\xi) d\xi.$$

Now let's differentiate under the integral sign to obtain

$$D_x^\alpha u(x) = (2\pi)^{-n} \int_{\mathbb{R}^n} e^{ix \cdot \xi} \xi^\alpha \hat{u}(\xi) d\xi.$$

From here we get

$$Au(x) = (2\pi)^{-n} \int_{\mathbb{R}^n} e^{ix \cdot \xi} a(x, \xi) \hat{u}(\xi) d\xi \quad (1)$$

where $a(x, \xi)$ is the *symbol* of the operator:

$$a(x, \xi) = \sum_{|\alpha| \leq m} a_\alpha(x) \xi^\alpha. \quad (2)$$

To obtain a pseudodifferential operator, we simply take a more general function $a(x, \xi)$ in (1), one which is not necessarily a polynomial in the ξ variables. The corresponding operator A is called the *quantization* of a , and we write

$$A = \text{Op}_0(a).$$

Here Op_0 stands for quantization for y -independent symbols (as opposed to general symbols defined later). Note that the Fourier multiplier Q defined above can now be written as

$$Q = \text{Op}_0 \left(\frac{1}{1 + |\xi|^2} \right).$$

To get good properties of $\text{Op}_0(a)$ we need to make certain assumption on the behavior of a at infinity. This brings us to symbol classes.

3. CLASSICAL KOHN–NIRENBERG SYMBOLS

We use the notation

$$\langle \xi \rangle := \sqrt{1 + |\xi|^2}.$$

This is asymptotic to $|\xi|$ as $\xi \rightarrow \infty$ and is also smooth at $\xi = 0$.

Definition 3.1. *Let $m \in \mathbb{R}$. We say $a(z, \xi) \in C^\infty(\mathbb{R}^p \times \mathbb{R}^n)$ lies in the symbol class $S^m(\mathbb{R}^p; \mathbb{R}^n)$, if for all multiindices α, β*

$$|\partial_z^\alpha \partial_\xi^\beta a(z, \xi)| \leq C_{\alpha\beta} \langle \xi \rangle^{m-|\beta|}$$

Here we allow $p \neq n$ for future use, but in the original formula (1) we have $p = n$.

The derivative bounds above mean the following: $a = \mathcal{O}(\langle \xi \rangle^m)$, differentiation in z does not change the growth of a , but differentiation in ξ gives decay by a power of ξ . As an exercise, you can check that:

- a polynomial of the form (2) lies in $S^m(\mathbb{R}^n; \mathbb{R}^n)$ if we assume that all derivatives of a_α lie in L^∞ , and
- the symbol of the operator Q above, $a(x, \xi) = \langle \xi \rangle^{-2}$, lies in $S^{-2}(\mathbb{R}^n; \mathbb{R}^n)$.

4. GENERAL QUANTIZATION FORMULA
AND THE POWER OF INTEGRATION BY PARTS

We now return to the formula (1). Let $a \in S^m(\mathbb{R}^n; \mathbb{R}^n)$. Recalling the definition of the Fourier transform, we rewrite (1) as

$$Au(x) = (2\pi)^{-n} \int_{\mathbb{R}^{2n}} e^{i(x-y)\cdot\xi} a(x, \xi) u(y) dy d\xi, \quad u \in \mathcal{S}(\mathbb{R}^n).$$

(This only makes sense for $m < -n$, more on that later.)

We arrive to the general quantization formula by allowing a to also depend on y . (This will be useful in deriving properties of quantization.) Namely, for $a \in S^m(\mathbb{R}^{2n}; \mathbb{R}^n)$ we define

$$\text{Op}(a)u(x) = (2\pi)^{-n} \int_{\mathbb{R}^{2n}} e^{i(x-y)\cdot\xi} a(x, y, \xi) u(y) dy d\xi. \quad (3)$$

For $m < -n$, the integral in (3) converges and we get

$$a \in S^m(\mathbb{R}^{2n}; \mathbb{R}^n), \quad m < -n \quad \implies \quad \text{Op}(a) : \mathcal{S}(\mathbb{R}^n) \rightarrow L^\infty(\mathbb{R}^n). \quad (4)$$

We now make sense of the oscillatory integral (3) for general a (in particular, for a which is a polynomial of the form (2)), by integrating by parts in y . Let's just do it

one time. We write

$$\begin{aligned}
\text{Op}(a)u(x) &= (2\pi)^{-n} \int_{\mathbb{R}^{2n}} (e^{i(x-y)\cdot\xi} \langle \xi \rangle^2) (\langle \xi \rangle^{-2} a(x, y, \xi) u(y)) dy d\xi \\
&= (2\pi)^{-n} \int_{\mathbb{R}^{2n}} ((1 - \xi \cdot D_y) e^{i(x-y)\cdot\xi}) (\langle \xi \rangle^{-2} a(x, y, \xi) u(y)) dy d\xi \quad (5) \\
&= (2\pi)^{-n} \int_{\mathbb{R}^{2n}} e^{i(x-y)\cdot\xi} (1 + \xi \cdot D_y) (\langle \xi \rangle^{-2} a(x, y, \xi) u(y)) dy d\xi.
\end{aligned}$$

The integration by parts in the third line above does make sense for $a \in S^m$ and $m < -n$. However, in the last line we compute

$$(1 + \xi \cdot D_y) (\langle \xi \rangle^{-2} a(x, y, \xi) u(y)) = \mathcal{O}(\langle \xi \rangle^{m-1} \langle y \rangle^{-n-1})$$

so the last integral actually converges when $m < 1-n$, which is better than the original definition of quantization! We can now integrate by parts repeatedly to arrive to

$$a \in S^m(\mathbb{R}^{2n}; \mathbb{R}^n), \text{ any } m \implies \text{Op}(a) : \mathcal{S}(\mathbb{R}^n) \rightarrow L^\infty(\mathbb{R}^n). \quad (6)$$

Some caution is needed here: what we really mean is that the linear operation $\text{Op}(a)$ is extended to the symbol class S^m by continuity from, say, S^{-n-1} . Such an extension is necessarily unique, and we can prove identities for $a \in S^m$ by just proving them for rapidly decaying a and arguing by approximation. (There are some subtleties here regarding “approximating by nice symbols”).

We can upgrade (6) further as follows:

$$a \in S^m(\mathbb{R}^{2n}; \mathbb{R}^n), \text{ any } m \implies \text{Op}(a) : \mathcal{S}(\mathbb{R}^n) \rightarrow \mathcal{S}(\mathbb{R}^n). \quad (7)$$

For that we need to apply the operators x_j, D_{x_j} to $\text{Op}(a)$ and use the identities

$$\begin{aligned}
x_j \text{Op}(a) &= \text{Op}(a)x_j - \text{Op}(D_{\xi_j} a), \\
D_{x_j} \text{Op}(a) &= \text{Op}(a)D_{x_j} + \text{Op}(D_{x_j} a + D_{y_j} a)
\end{aligned}$$

the first of which is proved by integrating by parts in ξ_j and the second one, by integrating by parts in y_j . We show the first one in a bit more detail since it will be used again later:

$$\begin{aligned}
(x_j \text{Op}(a) - \text{Op}(a)x_j)u(x) &= (2\pi)^{-n} \int_{\mathbb{R}^{2n}} e^{i(x-y)\cdot\xi} (x_j - y_j) a(x, y, \xi) u(y) dy d\xi \\
&= (2\pi)^{-n} \int_{\mathbb{R}^{2n}} (D_{\xi_j} e^{i(x-y)\cdot\xi}) a(x, y, \xi) u(y) dy d\xi \quad (8) \\
&= -(2\pi)^{-n} \int_{\mathbb{R}^{2n}} e^{i(x-y)\cdot\xi} D_{\xi_j} a(x, y, \xi) u(y) dy d\xi.
\end{aligned}$$

Since $D_{\xi_j} a, D_{x_j} a, D_{y_j} a$ still lie in S^m , we see that $x_j \text{Op}(a), D_{x_j} \text{Op}(a) : \mathcal{S}(\mathbb{R}^n) \rightarrow L^\infty(\mathbb{R}^n)$, and iteration gives $x^\alpha D_x^\beta \text{Op}(a) : \mathcal{S}(\mathbb{R}^n) \rightarrow L^\infty(\mathbb{R}^n)$, which implies (7).

The above discussion leads to the following statement:

Proposition 4.1. *Assume that $a \in S^m(\mathbb{R}^{2n}; \mathbb{R}^n)$. Then we may define*

$$\text{Op}(a) : \mathcal{S}(\mathbb{R}^n) \rightarrow \mathcal{S}(\mathbb{R}^n), \quad \text{Op}(a) : \mathcal{S}'(\mathbb{R}^n) \rightarrow \mathcal{S}'(\mathbb{R}^n).$$

Here the additional powers of x, y are handled similarly to the previous argument. To get the mapping property on tempered distributions, we use a definition by duality:

$$\langle \text{Op}(a)u, \varphi \rangle_{L^2} := \langle u, \text{Op}(a)^*\varphi \rangle_{L^2}, \quad u \in \mathcal{S}'(\mathbb{R}^n), \quad \varphi \in \mathcal{S}(\mathbb{R}^n)$$

where the adjoint $\text{Op}(a)^*$ has the form (3) with a different symbol (see below) and thus maps $\mathcal{S}(\mathbb{R}^n)$ to itself.

We call the resulting class of operators $\text{Op}(a)$ *pseudodifferential operators*. In particular, we denote by

$$\Psi^m(\mathbb{R}^n)$$

operators of the form $\text{Op}(a)$ where $a \in S^m(\mathbb{R}^{2n}; \mathbb{R}^n)$.

Define the residual operator class

$$\Psi^{-\infty}(\mathbb{R}^n) := \bigcap_{m \in \mathbb{R}} \Psi^m(\mathbb{R}^n).$$

One can show (with a bit of annoying technical work) that every element of $\Psi^{-\infty}$ has the form $\text{Op}(a)$ where a lies in the residual symbol class

$$S^{-\infty}(\mathbb{R}^{2n}; \mathbb{R}^n) = \bigcap_m S^m(\mathbb{R}^{2n}; \mathbb{R}^n).$$

Note that $a \in S^{-\infty}$ simply means that each derivative of a decays like $\mathcal{O}(\langle \xi \rangle^{-\infty})$. Then the integral kernel of (3) converges with all the x, y derivatives, implying that every $A \in \Psi^{-\infty}$ is a *smoothing operator*:

$$Au(x) = \int K(x, y)u(y) dy \quad \text{where} \quad K(x, y) \in C^\infty(\mathbb{R}^{2n}).$$

In particular we have the mapping property $A : \mathcal{S}' \rightarrow C^\infty(\mathbb{R}^n)$.

5. REDUCTION TO y -INDEPENDENT SYMBOLS

We now show that the general quantization procedure Op from (3) actually gives the same class of operators as the special quantization procedure Op_0 from (1), and get a useful asymptotic expansion:

Theorem 1. *Assume that $a \in S^m(\mathbb{R}^{2n}; \mathbb{R}^n)$. Then there exists $\tilde{a} \in S^m(\mathbb{R}^n; \mathbb{R}^n)$ such that*

$$\text{Op}(a) = \text{Op}_0(\tilde{a}).$$

Moreover, we have the asymptotic expansion

$$\tilde{a}(x, \xi) \sim \sum_{k=0}^{\infty} \tilde{a}_k, \quad \tilde{a}_k := \frac{1}{k!} (-i\partial_y \cdot \partial_\xi)^k a(x, y, \xi)|_{y=x}, \quad \partial_y \cdot \partial_\xi := \sum_{j=1}^n \partial_{y_j} \partial_{\xi_j} \quad (9)$$

in the following sense: for each N ,

$$\tilde{a}(x, \xi) - \sum_{k=0}^{N-1} \dots \in S^{m-N}(\mathbb{R}^n; \mathbb{R}^n).$$

Note here that the expansion does make sense: the k -th term in the expansion is in S^{m-k} due to the fact that Kohn–Nirenberg symbols improve by a power of ξ when differentiated in ξ . (This is the first time we use this fact, actually.) Note also that the above is an asymptotic expansion, not a convergent series! Asymptotic expansions like the one above are very common in microlocal analysis.

We will not prove Theorem 1 (see Melrose’s notes for a proof). We instead prove a simpler statement (from which the full theorem follows, but after a good amount of annoying technical work):

Proposition 5.1. *For each N we may write*

$$\text{Op}(a) = \text{Op}_0 \left(\sum_{k=0}^{N-1} \tilde{a}_k \right) + \text{Op}(r_N), \quad r_N \in S^{m-N}(\mathbb{R}^{2n}; \mathbb{R}^n)$$

where \tilde{a}_k are the terms in the expansion (9).

Proof. We just show the cases $N = 1, N = 2$, with higher N obtained similarly. We first do $N = 1$. The symbol

$$a(x, y, \xi) - \tilde{a}_0(x, \xi) = a(x, y, \xi) - a(x, x, \xi)$$

vanishes on the partial diagonal $\{x = y\}$. We can then write

$$\begin{aligned} a(x, y, \xi) - \tilde{a}_0(x, \xi) &= \int_0^1 \partial_t (a(x, x + t(y - x), \xi)) dt = \sum_{j=1}^n (y_j - x_j) b_j(x, y, \xi), \\ b_j(x, y, \xi) &= \int_0^1 (\partial_{y_j} a)(x, x + t(y - x), \xi) dt. \end{aligned} \tag{10}$$

From the definition of b_j we see that $b_j \in S^m(\mathbb{R}^{2n}; \mathbb{R}^n)$. We now use the key identity proved by (8):

$$b \in S^m(\mathbb{R}^{2n}; \mathbb{R}^n) \implies \text{Op}((y_j - x_j)b) = \text{Op}(D_{\xi_j} b). \tag{11}$$

We get then

$$\text{Op}(a) - \text{Op}_0(\tilde{a}_0) = \text{Op}(r_1), \quad r_1(x, y, \xi) := \sum_{j=1}^n D_{\xi_j} b_j(x, y, \xi)$$

and r_1 does lie in $S^{m-1}(\mathbb{R}^{2n}; \mathbb{R}^n)$ owing to the differentiation in ξ .

To do $N = 2$, we iterate this process further, applying it now to the symbol r_1 . We see that the next term in the expansion should be the restriction of r_1 to $\{x = y\}$;

indeed, $r_1(x, y, \xi) - r_1(x, x, \xi)$ can be again written in the form (10). It is easy to compute that

$$r_1(x, x, \xi) = -i(\partial_y \cdot \partial_\xi)a(x, y, \xi)|_{y=x} = \tilde{a}_1(x, \xi).$$

So we get

$$\text{Op}(a) - \text{Op}_0(\tilde{a}_0 + \tilde{a}_1) = \text{Op}(r_2), \quad r_2 \in S^{m-2}(\mathbb{R}^{2n}; \mathbb{R}^n).$$

□

For $A = \text{Op}(a) \in \Psi^m(\mathbb{R}^n)$, where $a \in S^m(\mathbb{R}^{2n}; \mathbb{R}^n)$, we define the *principal symbol* $\sigma^m(A)$ as follows:

$$\sigma^m(A) = [a(x, x, \xi)] \in \frac{S^m(\mathbb{R}^n; \mathbb{R}^n)}{S^{m-1}(\mathbb{R}^n; \mathbb{R}^n)}. \quad (12)$$

The principal symbol will have nice algebraic properties as we will see soon. What we see immediately from Proposition 5.1 is the following statement: if $A \in \Psi^m(\mathbb{R}^n)$, then

$$\sigma^m(A) = 0 \iff A \in \Psi^{m-1}(\mathbb{R}^n).$$

So the principal symbol does determine A modulo a lower order term. One often suppresses the order of the operator in the notation, writing σ instead of σ^m .

6. ADJOINTS

We now discuss algebraic properties of the classes $\Psi^m(\mathbb{R}^n)$. One algebraic property that we can do easily is closure under adjoints:

Theorem 2. *Assume that $A \in \Psi^m(\mathbb{R}^n)$. Then $A^* \in \Psi^m(\mathbb{R}^n)$ and $\sigma(A^*) = \overline{\sigma(A)}$. Here the adjoint is understood in the following sense:*

$$\langle Au, v \rangle_{L^2} = \langle u, A^*v \rangle_{L^2} \quad \text{for all } u, v \in \mathcal{S}(\mathbb{R}^n).$$

Proof. Let $A = \text{Op}(a)$ where $a \in S^m(\mathbb{R}^{2n}; \mathbb{R}^n)$. Then we have the following representation of the adjoint:

$$\text{Op}(a)^*v(x) = (2\pi)^{-n} \int_{\mathbb{R}^{2n}} e^{i(x-y)\cdot\xi} \overline{a(y, x, \xi)} u(y) dy d\xi.$$

From here the result follows immediately since $\text{Op}(a)^* = \text{Op}(a^*)$ where $a^*(x, y, \xi) = \overline{a(y, x, \xi)}$. □