

# MSRI LECTURE ON GENERAL RELATIVITY

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ABSTRACT. Rough notes for the lecture on general relativity at the MSRI introductory workshop in Fall 2019.

- Einstein Vacuum Equations: (EVE)  $\text{Ein}(g) + \Lambda g = 0$  where  $\Lambda \in \mathbb{R}$ ,  $g$  a Lorentzian metric on  $M$  and  $\text{Ein}(g) = \text{Ric}(g) - \frac{1}{2}\text{scal}_g g$ .
  - Upon taking traces, EVE is equivalent to  $\text{Ric}(g) - \Lambda g = 0$
- Examples  $\Lambda = 0$ :
  - (1) Minkowski space:  $M = \mathbb{R}_t \times \mathbb{R}_x^3$ ,  $g_0 = -dt^2 + dx^2$ 
    - The local model of any Lorentzian spacetime
  - (2) Schwarzschild black hole:  $m > 0$  (mass of the black hole),  $M = \mathbb{R}_t \times (2m, \infty)_r \times \mathbb{S}_\omega^2$  and  $g_m = -(1 - \frac{2m}{r})dt^2 + (1 - \frac{2m}{r})^{-1}dr^2 + r^2g_{\mathbb{S}^2}$ 
    - Unique vacuum spherically symmetric family
    - The apparent singularity in the second term of the metric is only a coordinate singularity and can be removed by a change of coordinates
      - \* Light cones begin to tilt toward the black hole ( $r = 0$ ) as they approach  $r = 2m$
      - \* Light or massive observer must reach  $r = 0$  in finite time once their distance from  $r = 0$  is less than or equal to  $2m$
    - Trapping occurs at  $r = 3m$  (photosphere); in phase space  $T^*M$ , we have normally hyperbolic trapping
  - (3) Kerr black hole:  $m > 0$ ,  $0 \leq a < m$ ,  $g_{m,a} = \dots$ 
    - In  $T^*M$ , there is normally hyperbolic trapping (Wunsch-Zworski, Dyatlov)
- Examples  $\Lambda > 0$ :
  - (1) De-Sitter space:  $g = \frac{-d\tau^2 + dy^2}{\tau^2}$  for  $\Lambda = 3$
  - (2) Schwarzschild de-Sitter:
    - The metric is  $g_m = -(1 - \frac{2m}{r} - \frac{\Lambda r^2}{3})dt^2 + (1 - \frac{2m}{r} - \frac{\Lambda r^2}{3})^{-1}dr^2 + r^2g_{\mathbb{S}^2}$
    - Event horizon at  $r \approx 2m$  and cosmological horizon and  $r \approx \sqrt{\frac{3}{\Lambda}}$
  - (3) Kerr de-Sitter
- The Initial Value Problem:
  - $\text{Ric}(g)_{ij} = g^{kl}\partial_k\partial_l g_{ij} + g\partial^2 g + \text{l.o.t.}$

- One sees that the first term is  $\square_g g_{ij}$ .

**Theorem 1.** (Choquet-Bruhat)  $\Lambda \in \mathbb{R}$ . Given  $(\Sigma^3, \gamma)$  and Riemannian manifold,  $K$  a symmetric two-tensor on  $\Sigma$  satisfying the constraint equations, there exists  $(M^4, g)$  such that  $g$  satisfies the EVE and  $\iota : \Sigma \hookrightarrow M$  such that  $\gamma = \iota^*g$  and  $K = II_{\iota(\Sigma)}^g$  which is unique up to a diffeomorphism.

- Proof sketch:

- \* Take  $\Lambda = 0$ .
- \* Guess  $M = (\epsilon, \epsilon)_t \times \Sigma$
- \* Solve  $\text{Ric}(g) = 0$  along with  $W(g) = (1\text{-form in } g \text{ and } \partial g) = 0$
- \* Typically  $W(g) = \text{tr}_g(\nabla^g - \nabla^{g_0}) \sim \square_{g, g_0} \mathbb{1}$
- \* So,  $W(g) = 0$  means  $\mathbb{1} : (M, g) \rightarrow (M, g_0)$  is a wave map
- \* Consider the reduced Einstein vacuum equations (REVE):

$$\text{Ric}(g) - \delta_g^* W(g) = 0$$

where  $\delta_g^*$  is the symmetric gradient

- \* The second term kills off the non-wave term in the EVE
- \* REVE is a quasilinear wave equation
- \* Given  $(\gamma, K) \in (S^2(T^*\Sigma))^2$ , cook up  $(g^0, g^1) \in (S^2(T^*M))^2$  such that  $\gamma = \iota^*g$  and  $K = II_{\iota(\Sigma)}^g$  and  $W(g^0 + tg^1) = 0$  at  $t = 0$
- \* Solve REVE using the constraint equations and  $\mathcal{L}_{\partial_t} W(g) = 0$  at  $t = 0$
- \* Use second Bianchi identity to deduce that  $\text{Ric}(g) = 0$

- Stability Problems:

- Minkowski space:

- \*  $\Sigma = \mathbb{R}_x^3$
- \*  $|\gamma_{ij} - (g_m)_{ij}| \leq \epsilon r^{-1-\alpha}$ ,  $r \geq 1$ ,  $|m| < \epsilon$
- \*  $|k_{ij}| \leq \epsilon \langle r \rangle^{-2-\alpha}$
- \* Analogous assumptions for derivatives

**Theorem 2.** (Christodoulou-Klainerman, Friedrich, Lindblad-Rodnianski, H-Vasy) There exists a geodesically complete  $g$  with  $\text{Ric}(g) = 0$  on  $\mathbb{R}^4$  with

$$|g_{ij} - (g_0)_{ij}| \leq (1 + |t| + r)^{-1}$$

**Theorem 3.** (H-Vasy) If  $(\gamma, K)$  are polyhomogeneous (phg) on  $\overline{\mathbb{R}^3}$  then  $g$  is phg on the compactification of  $\mathbb{R}^4$ .

- Kerr stability: (open)

- \* Decay rate in Minkowski is polynomial
- \* Taturu's work implies decay is  $t^{-3}$

- $\Lambda > 0$  gives better decay

**Theorem 4.** (*H-Vasy*) (*KdS stability*)  $(\gamma, K)$  close to data of  $g_{m_i, a_i}^{KdS}$  and  $|a_i| \ll m_i$ . There there exists  $g$  on  $\Omega$  such that  $\text{Ric}(g) - \Lambda g = 0$  and  $(m_f, a_f)$  such that  $g = g_{m_f, a_f}^{KdS} + \tilde{g}$  where  $\tilde{g}$  decays and  $g_{m_f, a_f}^{KdS}$  is constant.

\* Analysis in  ${}^bT^*M$  and is as nonlinear waves lecture

\* Sketch of proof:

· Solve REVE

$$\begin{cases} P(g) = (\text{Ric} - \Lambda)(g) - \delta_g^* W(g) = 0 \\ i.c.(g^0, g^1) \end{cases} \quad (1)$$

which is a quasilinear wave equation

- Use iteration scheme
- Show for  $g^{(k)} = g_{dS} + \tau^\alpha H_b^s$  that  $(D_g P)^{-1} \sim (\square_g + \text{l.o.t.})^{-1} : \tau^\alpha H_b^{s-1} \rightarrow \tau^\alpha H_b^s$
- Run the iteration scheme
- Obtain decaying solution
- Suffices to compute the resonances of  $D_{g_{dS}} P$
- Bad resonances are artifacts of gauge invariance
- Change gauge in consistent way to work around bad resonance