

Inverse problems for real principal type operators

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Outline

1. Motivation
2. Results for real principal type operators
3. Methods of proof

1. Calderón problem (elliptic PDE)

Laplace-Beltrami equation

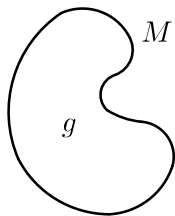
$$\begin{cases} \Delta_g u = 0 & \text{in } M, \\ u = f & \text{on } \partial M \end{cases}$$

where (M, g) is a compact Riemannian manifold with boundary ($g \iff$ electrical conductivity).

Boundary measurements given by the **Dirichlet-to-Neumann (DN) map**

$$\Lambda_g : C^\infty(\partial M) \rightarrow C^\infty(\partial M), f \mapsto \partial_\nu u|_{\partial M}.$$

Inverse problem: given Λ_g , recover g .



2. Gel'fand problem (hyperbolic PDE)

Wave equation ($g \leftrightarrow$ sound speed)

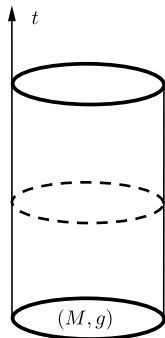
$$\begin{cases} (\partial_t^2 - \Delta_g)u = 0 & \text{in } M \times (0, T), \\ u = f & \text{on } \partial M \times (0, T) \\ u|_{\{t < 0\}} = 0 \end{cases}$$

where (M, g) is compact with boundary.

Boundary measurements given by the
hyperbolic DN map

$$\Lambda_g^{\text{Hyp}} : f \mapsto \partial_\nu u|_{\partial M \times (0, T)}.$$

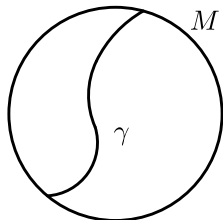
Inverse problem: given Λ_g^{Hyp} , recover g .



Ray transform/scattering relation (transport PDE)

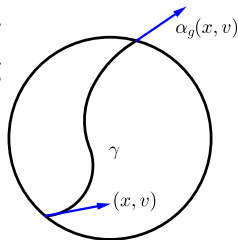
3. Try to recover a function f in (M, g) from its *geodesic X-ray transform* I_f , where

$$I_f(\gamma) = \int_{\gamma} f dt, \quad \gamma \text{ maximal geodesic.}$$



4. Related **scattering rigidity** problem: recover (M, g) from its **scattering relation** α_g , relating initial and final data of maximal geodesics:

$$\alpha_g : (x, v) \mapsto \alpha_g(x, v)$$



Can formulate both questions in terms of (transport) PDEs. These are **highly nonlinear questions** related to **linear PDEs**!

Connections

Unexpected connections in special geometries:

- ▶ Calderón problem reduces to geodesic X-ray transform
[Dos Santos-Kenig-S-Uhlmann 2009]
- ▶ Calderón problem reduces to Gel'fand problem
[Dos Santos-Kurylev-Lassas-S 2016]
- ▶ scattering rigidity problem reduces to Calderón problem
[Pestov-Uhlmann 2005]

What are the general structural conditions and mechanisms behind this?

Goal

Propose to study inverse problems for **general** differential operators. Hope to understand:

- ▶ structural conditions for treating classes of operators
- ▶ fundamental mechanisms for solving inverse problems
- ▶ the extent to which it is possible to push existing methods.

Approach in the spirit of [Hörmander, ALPDO vol. III-IV].
Earlier results for constant coefficients [Isakov 1991, . . . ,
Hitrik-Krupchyk-Ola-Päivärinta 2010].

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Real principal type operators

Let M be compact with boundary. A differential operator P on M , of order $m \geq 1$, is **real principal type** [Hörmander vol. IV] if

- ▶ it has **real principal symbol** $\sigma_{\text{pr}}(P) = p_m$, and
- ▶ the null bicharacteristic flow is **nontrapping**.

Null bicharacteristic curves are integral curves of H_{p_m} in $p_m^{-1}(0)$. In coordinates $\gamma(t) = (x(t), \xi(t))$ with

$$\begin{cases} \dot{x}(t) &= \nabla_{\xi} p_m(x(t), \xi(t)), \\ \dot{\xi}(t) &= -\nabla_x p_m(x(t), \xi(t)). \end{cases}$$

Nontrapping means that any such $\gamma(t)$ reaches ∂M in finite time in both directions.

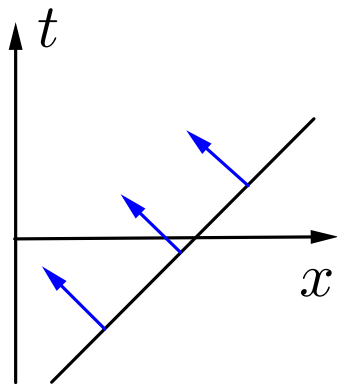
Real principal type operators

Examples:

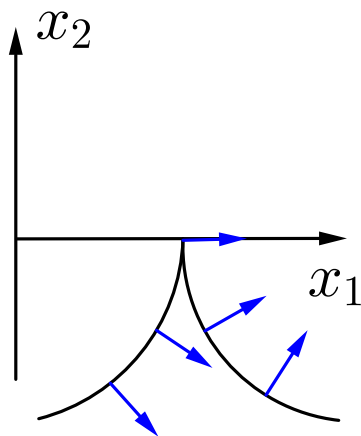
- ▶ real vector fields with no trapped integral curves
- ▶ wave operator in $M \times (0, T)$, Lorentzian wave operators, strictly hyperbolic operators with nontrapping condition
- ▶ Tricomi type operators, e.g. $x_2 D_{x_1}^2 + D_{x_2}^2$
- ▶ Schrödinger operator $i\partial_t + \Delta$, plate equation $\partial_t^2 + \Delta^2$ with suitable (anisotropic) weighting for ∂_t

Real principal type operators can be microlocally conjugated to normal form D_{x_1} . Singularities of solutions propagate along null bicharacteristics, solvability theory for $Pu = f$.

Null bicharacteristics



Wave operator $\partial_t^2 - \Delta$



Tricomi operator $x_2 D_{x_1}^2 + D_{x_2}^2$

Boundary measurements

It is not clear how to define an analogue of DN map for a general operator. However, we consider the **Cauchy data set**

$$C_P = \{(u|_{\partial M}, \dots, \nabla^{m-1} u|_{\partial M}); Pu = 0 \text{ in } M, u \in H^m(M)\}.$$

This is equivalent to knowing the DN map e.g. in the Calderón and Gel'fand problems.

Inverse problem: given C_P , determine information about P .

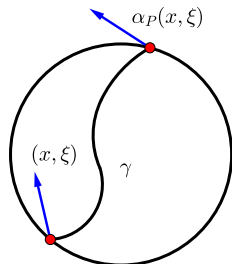
From now on, all operators will be **real principal type** in M .

Determining (sub)principal information

Theorem 1

If $C_{P_1} = C_{P_2}$ and if $P_1 = P_2$ to infinite order on ∂M , then

$$\alpha_{P_1} = \alpha_{P_2}$$



where α_P is the **bicharacteristic scattering relation**, mapping an initial point of a **maximal null bicharacteristic** to its final point.

Moreover, if P_1 and P_2 have the same principal symbol, then

$$\exp \left[i \int \sigma_{\text{sub}}(P_1)(\gamma(t)) dt \right] = \exp \left[i \int \sigma_{\text{sub}}(P_2)(\gamma(t)) dt \right]$$

for any **maximal null bicharacteristic** γ in T^*M .

Lower order coefficients

The conclusion $\exp[i \int \dots] = \exp[i \int \dots]$ is equivalent with

$$\int \sigma_{\text{sub}}(P_1)(\gamma(t)) dt = \int \sigma_{\text{sub}}(P_2)(\gamma(t)) dt \pmod{2\pi\mathbb{Z}}.$$

This is related to the **Aharonov-Bohm effect** in determining subprincipal terms on domains with nontrivial topology. For lower order coefficients, this effect does not appear:

Theorem 2 (Bicharacteristic ray transforms)

If $C_{P+Q_1} = C_{P+Q_2}$ where Q_j are operators of order $\leq m - 2$, then

$$\int \sigma_{\text{pr}}(Q_1)(\gamma(t)) dt = \int \sigma_{\text{pr}}(Q_2)(\gamma(t)) dt$$

for any maximal null bicharacteristic γ in T^*M .

Real principal type operators

The results are quite general: they extend known results for wave equations [Stefanov-Yang 2018], and are valid for

- ▶ operators of any order, with **real principal symbol** and **nontrapping** condition (no wellposedness assumptions)
- ▶ **any maximal bicharacteristic**, even with cusps (Tricomi) and tangential reflections (but not at endpoints for α_P)

However, the results are **conditional**: in order to recover coefficients of P , one still needs to analyze the scattering relation α_P or bicharacteristic ray transforms.

Boundary determination

Determine Taylor series of coefficients of P at null points $(x, \xi) \in T^*(\partial M)$, based on zeros of **characteristic polynomial**

$$t \mapsto p_m(x, \xi + t\nu).$$

Two methods:

1. **Elliptic region.** If there is a simple non-real zero, use exponentially decaying solutions (analogue of boundary determination for Laplace equation).
2. **Hyperbolic region.** If there are two distinct real zeros, use solutions concentrating near two null bicharacteristics (analogue of boundary determination for wave equation).

Boundary determination

Theorem 3 (Determining Taylor series of a potential)

If $V_1, V_2 \in C^\infty(M)$ and $C_{P+V_1} = C_{P+V_2}$, then

$$\nabla^k V_1(x_0) = \nabla^k V_2(x_0), \quad k \geq 0,$$

at any $x_0 \in \partial M$ so that for some $\xi \in T_x^*(\partial M)$, the map $t \mapsto p_m(x_0, \xi + t\nu)$ either has a simple non-real root, or two distinct real roots¹.

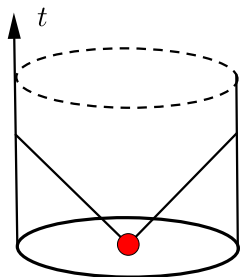
In particular, if M and V_j are **real-analytic** and there is one such x_0 , then $V_1 = V_2$ **everywhere in M** .

¹with corresponding bicharacteristics intersecting nicely at x_0

Boundary determination

Observations:

- ▶ boundary determination in general not possible for $m = 1$
- ▶ even for wave equation, can do boundary determination in the elliptic region as for elliptic operators (local argument)
- ▶ boundary determination in the hyperbolic region is global in character



$$P = \partial_t^2 - \Delta + V$$

Nonlinear equations

If $q \in C^\infty(M)$, consider the semilinear equation

$$Pu + q(x)u^k = 0 \text{ in } M.$$

Let C_q^{small} be the Cauchy data set for small solutions.

Theorem 4 (Semilinear equations)

Let $q_1, q_2 \in C^\infty(M)$ and $k \geq 3$. If $C_{q_1}^{\text{small}} = C_{q_2}^{\text{small}}$, then $q_1 = q_2$ in B where

$B = \{x \in M; \text{ there are two null bicharacteristics that intersect only once at } x \text{ transversally}\}.$

Nonlinearity helps (proof fails if $k = 1$)! Wave equations:

[Kurylev-Lassas-Uhlmann 2018, Lassas-Uhlmann-Wang 2018, Hintz-Uhlmann 2018]

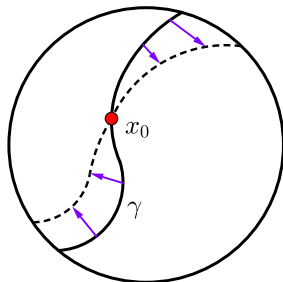
Nonlinear equations

Can recover the coefficient $q(x)$ in the set

$B = \{x_0 \in M; \text{there are two null bicharacteristics that intersect only once at } x_0 \text{ transversally}\}$.

If there is a nice¹ bicharacteristic $\gamma(t)$ through x_0 having a variation field only vanishing at $t = 0$, can recover $q(x_0)$.

Works e.g. if some $\gamma(t)$ through x_0 has "no conjugate points". May fail if there is a "maximally conjugate" point.



¹nontangential, no cusp at x_0 , $x(t)$ does not self-intersect

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Methods

1. Use Cauchy data of **special solutions** concentrating along a **null bicharacteristic** (propagation of singularities).
2. For boundary determination, also use **exponentially decaying solutions** concentrating at a boundary point.
3. Use **integral identities** and a **mix-and-match construction** to pass from Cauchy data set C_P to scattering relation / bicharacteristic ray transforms / pointwise information.

Quasimode construction

Theorem 5

Let P have **real principal symbol** in an open mfld X , and let $\gamma : [0, T] \rightarrow T^*X$ be an **injective** null bicharacteristic segment. There is $u = u_h \in C_c^\infty(X)$ with

$$WF_{\text{scl}}(u) = \gamma([0, T]), \quad WF_{\text{scl}}(Pu) = \gamma(0) \cup \gamma(T),$$

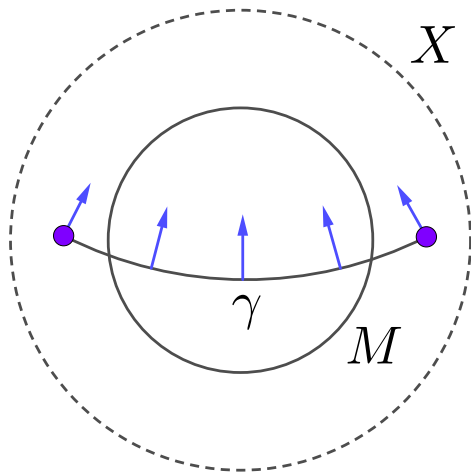
having semiclassical defect measure (with $c_\gamma(t) > 0$)

$$\lim_{h \rightarrow 0} (Op_h(a)u_h, u_h)_{L^2(X)} = \int_0^T a(\gamma(t))c_\gamma(t) dt$$

whenever $a \in S^0$ vanishes near endpoints of γ .

Semiclassical counterpart of [\[Duistermaat-Hörmander 1972\]](#).

Quasimode construction



$$WF_{\text{scl}}(Pu) = \gamma(0) \cup \gamma(T) \implies Pu = O(h^\infty) \text{ in } M$$

Methods for constructing quasimodes

1. Locally, enough to use **geometrical optics**:

$$u_h(x) = e^{i\varphi(x)/h} a(x)$$

where φ is real and $p_m(x, d\varphi(x)) = 0$ (**eikonal equation**).

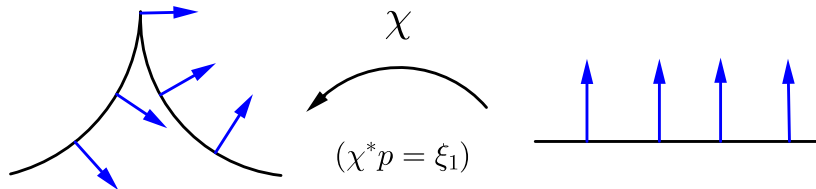
2. No cusps ($\dot{x}(t) \neq 0$): can use a **Gaussian beam** construction

$$u_h(x) = e^{i\Phi(x)/h} a(x)$$

where Φ is **complex** and solves eikonal equation to infinite order on the curve $x(t)$. Main point: $\nabla^2\Phi|_{x(t)}$ solves a matrix Riccati equation and $\text{Im}(\nabla^2\Phi) > 0$.

Quasimodes [Duistermaat-Hörmander 1972]

3. If $\gamma(t)$ is injective but may have cusps, can straighten $\gamma(t)$ in phase space by a canonical transformation χ .



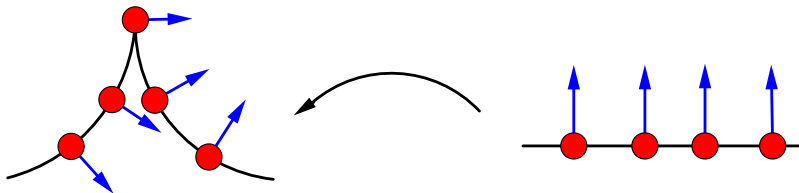
Multiply P by an elliptic Ψ DO so that P becomes of order 1. Construct Fourier integral operators A, B such that

$$BPA \approx D_{x_1} \text{ microlocally near } \gamma.$$

Quasimode U for $D_{x_1} \implies u = AU$ is a quasimode for P .

Quasimodes (direct construction)

4. Think of quasimodes as superpositions of wave packets $\approx e^{i\frac{\xi(t)\cdot(x-x(t))}{h}} e^{-\frac{|x-x(t)|^2}{2h}}$ at $x(t)$ oscillating in direction $\xi(t)$.



Look for u_h with $Pu_h = O(h^\infty)$ directly in the form

$$u_h(x) = \int_0^T e^{i\Phi(x,t)/h} a(x,t) dt.$$

Cf. Gaussian beam construction along $(x(t), t)$ in $X \times \mathbb{R}$.

Future directions

1. Inversion of scattering relation α_P ? If $P = X_g$,¹ studied in [Pestov-Uhlmann 2005, Stefanov-Uhlmann-Vasy 2017].
2. Inversion of bicharacteristic ray transform? Cf. geodesic ray transform [Uhlmann-Vasy 2016, Paternain-S-Uhlmann 2015] ($P = X_g$), and light ray transform [Lassas et al 2019] ($P = \square_g$).
3. Results for mild trapping? For $P = X_g$ and hyperbolic trapping, studied in [Guillarmou 2017].
4. Can one associate a symbol directly to C_P ?
5. The results are in the spirit of using singularities of the integral kernel of DN map. Can one extract information from the C^∞ part of the kernel?

¹geodesic vector field on unit sphere bundle