Semi-classical analysis of Schrödinger equation on *H*-type groups

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Recent developments in microlocal analysis Berkeley, October 18th. 2019 • [Folland & Stein 1974], [Rotschild & Stein 1976], [Folland, 1977]: Analysis of sublaplacians can be performed using Lie groups theory, lifting procedure.

Example: Heisenberg group and the operator *L* on \mathbb{R}^3 :

$$L = X^2 + Y^2, \ X = \partial_x - \frac{y}{2}\partial_z, \ Y = \partial_y + \frac{x}{2}\partial_z, \ [X, Y] = \partial_z$$

 What would be a semiclassical/microlocal approach on graded Lie groups using theory of representations ?
 Which class of pseudodifferential operators ?

[Taylor 1984], [Beals & Greiner 1988], [Christ, Geller, Glowasky & Polin 1992] [Geller 1990] [Bahouri, FK & Gallagher 2012], [Fischer & Ruzhansky 2016], [FK & Fischer, 2018 & 2019]

The question: Large time evolution of energy density of families of solutions of a Schrödinger equation

• Let V_j be 2*d* vector fields of \mathbb{R}^{2d+p} and

$$-\Delta_G = \sum_{1 \le j \le 2d} V_j^2.$$

• Assume $\mathfrak{g} = \mathfrak{v} \oplus \mathfrak{z}$ is a Lie algebra of Heisenberg type with centre \mathfrak{z} , $\mathfrak{v} = \operatorname{Vect}(V_i, \ 1 \le j \le 2d\}, \ \mathfrak{z} = \operatorname{Vect}([V_i, V_j], \ 1 \le i, j \le 2d),$ $2d = \dim \mathfrak{v}, \ p = \dim \mathfrak{z}.$

• Let $(\psi_0^{\varepsilon})_{\varepsilon>0}$ be bounded in $L^2(\mathbb{R}^d)$ such that $\exists s, C_s > 0, \forall \varepsilon > 0, \\ \varepsilon^s \|(-\Delta_G)^{\frac{s}{2}} \psi_0^{\varepsilon}\|_{L^2(G)} + \varepsilon^{-s} \|(-\Delta_G)^{-\frac{s}{2}} \psi_0^{\varepsilon}\|_{L^2(G)} \le C_s.$

Describe for $\phi \in \mathcal{C}^{\infty}_{c}(G)$, $T \in \mathbb{R}$, $\tau > 0$

$$\lim_{\varepsilon\to 0}\frac{1}{T}\int_0^T\int_G\phi(x)|\mathrm{e}^{i\frac{t}{2\varepsilon^\tau}\varepsilon^2\Delta_G}\psi_0^\varepsilon(x)|^2dx\,dt.$$

The same question in the Euclidean case

Let $(\psi_0^{arepsilon})$ be a bounded family in $L^2(\mathbb{R}^d)$ satisfying

 $\exists s, C_s > 0, \ \forall \varepsilon > 0, \ \varepsilon^s \| (-\Delta)^{\frac{s}{2}} \psi_0^{\varepsilon} \|_{L^2(\mathbb{R}^d)} + \varepsilon^{-s} \| (-\Delta)^{-\frac{s}{2}} \psi_0^{\varepsilon} \|_{L^2(\mathbb{R}^d)} \leq C_c.$

Then any limit point of the measure

$$\left| \mathrm{e}^{-i \frac{t}{2\varepsilon^{\tau}} \varepsilon^2 \Delta} \psi_0^{\varepsilon} \right|^2 dx dt$$

is of the form $\varrho_t(x)dt$ where $\varrho_t \in \mathcal{M}^+(\mathbb{R}^d)$ and

- 1 If $\tau \in (0,1)$, then $\varrho_t = \varrho_0$.
- 2 If $\tau = 1$ then $\varrho_t(x) = \int_{\mathbb{R}^d} \mu_0(x t\xi, d\xi)$.
- $If \tau > 1 then \varrho_t = 0.$

The result

Theorem (FK & Fischer 2019)

Any weak limit of $|e^{i\frac{t}{2\varepsilon^{\tau}}\varepsilon^{2}\Delta_{G}}\psi_{0}^{\varepsilon}(x)|^{2}dx dt$ writes $\varrho_{t}dt = \left(\varrho_{t}^{\mathfrak{v}^{*}} + \varrho_{t}^{\mathfrak{z}^{*}}\right)dt$ with

• If
$$\tau \in (0,1)$$
, for all $t \in \mathbb{R}$, $\varrho_t = \varrho_0$

• If
$$\tau = 1$$
, then $\varrho_t^{\mathfrak{z}^*} = \varrho_0^{\mathfrak{z}^*}$ and $\varrho_t^{\mathfrak{v}^*}(x) = \int_{\mathfrak{v}^*} \varsigma_0\left(\operatorname{Exp}(t\,\omega\cdot V)x,d\omega\right)$.

• If $\tau \in (1,2)$, then $\varrho_t^{\mathfrak{v}^*} = 0$ and $\partial_t \varrho_t^{\mathfrak{z}^*} = 0$ holds in $\mathcal{D}'(\mathbb{R} \times G)$.

• If $\tau = 2$, then $\varrho_t^{\mathfrak{v}^*} = 0$ and $\varrho_t^{\mathfrak{z}^*} = \sum_{n \in \mathbb{N}} \int_{\mathfrak{z}^* \setminus \{0\}} \gamma_{n,t}(x, d\lambda)$ where

$$\left(\partial_t - \frac{2n+d}{2|\lambda|}\mathcal{Z}^{(\lambda)}\right)\gamma_{n,t} = 0,$$

where $\mathcal{Z}(\lambda)$ is the left invariant vector field corresponding to $\lambda \in \mathfrak{z}^*$. • If $\tau > 2$, then $\varrho_t = 0$ for all $t \in \mathbb{R}$.

Remarks

- Quantum limits split into two parts, with different threshold indexes.
- Dispersion takes longer than in the Euclidean case [Bahouri, Gérard & Xu 2000], [Del Hierro 2005], [Bahouri, FK, Gallagher 2016]
- Splitting and invariance properties already noticed (p = 1 contact manifolds, Grauert tubes) in [Zelditch 1997], [Colin de Verdière, Hillairet & Trélat 2018], [Burq & Sun 2019]

- Graded, stratified Lie groups, H-type groups
- Analysis on graded Lie groups and Fourier transform
- Seudodifferential operators, Egorov theorem
- Semi-classical measures and Schrödinger equation.

Graded Lie groups

Graded Lie groups - Definition

Definition

A simply connected Lie group G is graded if the Lie algebra \mathfrak{g} of its left-invariant vector fields is graded :

$$\mathfrak{g} = \mathfrak{g}_1 \oplus \cdots \oplus \mathfrak{g}_n, \ [\mathfrak{g}_\ell, \mathfrak{g}_{\ell'}] \subset \mathfrak{g}_{\ell+\ell'}.$$

The group G is stratified if

$$\forall \ell \in \{1, \cdots n-1\}, \ [\mathfrak{g}_{\ell}, \mathfrak{g}_1] = \mathfrak{g}_{\ell+1}.$$

• G identifies to g via the exponential map:

$$egin{array}{cccc} \exp & : & \mathfrak{g} &
ightarrow & G \ & X & \mapsto & \exp(X)e_G \end{array}$$

• The law group on G is a polynomial map (Campbell-Baker-Hausdorff).

Graded Lie groups - Heisenberg group

- The Heisenberg group:
 - The Lie algebra $\mathfrak{h}=\mathfrak{v}\oplus\mathfrak{z}\,$ is generated by

$$X = \partial_x - \frac{y}{2}\partial_z, \quad Y = \partial_y + \frac{x}{2}\partial_z, \quad Z = \partial_z = [X, Y].$$

 \implies The group \mathbb{H} is stratified and has two steps.

• The points of G are the elements

$$w = \operatorname{Exp}(xX + yY + sS), \ (x, y, s) \in \mathbb{R}^3.$$

• $\mathbb{H} = \mathbb{R}^3$ with the product law

$$(x, y, s) \cdot (x', y', s') = (x + x', y + y', s + s' + \frac{1}{2}(x \cdot y' - x' \cdot y)).$$

• The center of \mathbb{H} is the set of points of the form (0, 0, s).

Graded Lie groups - H-type groups

● H-type groups (multidimensional versions of Ⅲ [Kaplan 80])

The Lie algebra of G is stratified with two steps g = v ⊕ 3.
 If B(λ) is the skew symmetric form defined on v × v by

 $\forall \lambda \in \mathfrak{z}^*, \ \forall U, V \in \mathfrak{v}, \ B(\lambda)(U, V) = \lambda([U, V]),$

then $B(\lambda)^2 = -|\lambda|^2 \text{Id.}$

• Notation (1): Let (V_1, \ldots, V_{2d}) be an ONB of v and (Z_1, \cdots, Z_p) of \mathfrak{z} , then $x \in G$ or $X \in \mathfrak{g}$ writes

 $x = \operatorname{Exp}(X), \qquad X = v_1 V_1 + \ldots + v_{2d} V_{2d} + z_1 Z_1 + \ldots + z_p Z_p.$

Graded Lie groups - H-type groups

• Notation (2): v decomposes in a λ -depending way as $v = p_{\lambda} \oplus q_{\lambda}$ with

$$\mathfrak{p} := \mathfrak{p}_{\lambda} := \mathsf{Span}\left(P_1, \dots, P_d\right), \quad \mathfrak{q} := \mathfrak{q}_{\lambda} := \mathsf{Span}\left(Q_1, \dots, Q_d\right).$$

where $(P_1, \ldots, P_d, Q_1, \ldots, Q_d)$ be an ONB of v such that

$$B(\lambda)(U,V) = |\lambda| U^t J V, \ \ J = egin{pmatrix} 0 & \mathrm{Id} \ -\mathrm{Id} & 0 \end{pmatrix}.$$

Then $x \in G$ or $X \in \mathfrak{g}$ write

 $x = \text{Exp}(X), \ X = p_1 P_1 + \ldots + p_d P_d + q_1 Q_1 + \ldots + q_d Q_d + z_1 Z_1 + \ldots + z_p Z_p.$

Besides

$$[P_j, Q_j] = \mathcal{Z}^{\lambda}, \quad \forall j \in \{1, \cdots, d\}$$

where \mathcal{Z}^{λ} corresponds to the vector λ in the identification $\mathfrak{z} \sim \mathfrak{z}^*$.

Analysis on *H*-type groups

Analysis on *H*-type groups

• The Haar measure, *dx* is deduced on *G* from the Lebesgue meas. on *g*. Associated Lebesgue spaces on *G*:

$$\|f\|_{L^p(G)} := \left(\int_G |f(x)|^p dx\right)^{\frac{1}{p}}, \ p \in [1,\infty[.$$

• The dilatation. For r > 0 and $X = V + Z \in \mathfrak{g}$, one sets $\delta_r X = rV + r^2 Z$.



 $r \cdot \operatorname{Exp}(V + Z) := \operatorname{Exp}(r V + r^2 Z), \ r > 0.$

 \implies Homogeneous functions and operators

• Homogeneous dimension. The Haar measure is *Q*-homogeneous:

 $Q := \dim \mathfrak{v} + 2 \dim \mathfrak{z} = 2d + 2p.$

Let \hat{G} be the set of irreducible representations of G:

 $\widehat{\mathcal{G}} = \{ \text{class of } \pi^{\lambda} \ : \ \lambda \in \mathfrak{z}^* \setminus \{0\} \} \sqcup \{ \text{class of } \pi^{0, \omega} \ : \ \omega \in \mathfrak{v}^* \}.$

Infinite dimensional representations are parametrized by *s*^{*} \ {0}: for λ ∈ *s*^{*} \ {0}, then *H*_λ = *L*²(*p*_λ) with for x = Exp(P + Q + Z) ∈ G,

$$\pi_x^{\lambda} \Phi(\xi) = \exp\left[i\lambda(Z) + \frac{i}{2}|\lambda| P \cdot Q + i\sqrt{|\lambda|} \xi \cdot Q\right] \Phi\left(\xi + \sqrt{|\lambda|}P\right).$$

• Finite dimensional representations are parametrized by v^* : for $\omega \in v^*$,

$$\pi^{0,\omega}_x = \mathrm{e}^{i\omega(V)}, \quad x = \mathrm{Exp}(V+Z) \in \mathcal{G}.$$

The dual set of H-type groups

• Plancherel measure: $d\mu(\pi^{\lambda}) = |\lambda|^d d\lambda$.

 \implies The set of finite dimensional representations is of 0 Plancherel measure.

• Dilations on \hat{G} : $r \cdot \pi_x = \pi_{r,x}$

• for $r\in \mathbb{R}^*_+$, $\lambda\in \mathfrak{z}^*\setminus\{0\}$, $x\in {\mathcal G}$,

$$\pi_{r\cdot x}^{\lambda} = T_r \pi_x^{r^2 \lambda} T_r^*, \ T_r f(\xi) = r^{1/2} f(r\xi),$$

 $\implies r \cdot \lambda = r^2 \lambda$

• for $r \in \mathbb{R}^*_+$, $\omega \in \mathfrak{v}^*$, $x \in G$, $\pi^{0,\omega}_{r\cdot x} = \mathrm{e}^{ir\omega \cdot x}$

 \implies $r \cdot (0, \omega) = (0, r\omega)$

Fourier transform

Definition

Let $f \in L^1(G)$, the Fourier transform of f is the operator of $\mathcal{L}(\mathcal{H}_{\lambda})$,

$$\widehat{f}(\lambda) = \int_{\mathcal{G}} f(x) \pi^{\lambda}(x)^* dx.$$

- Notation: $\hat{f}(\lambda) = \hat{f}(\pi^{\lambda}), \ \hat{f}(0,\omega) = \hat{f}(\pi^{0,\omega}).$
- Extension to $L^2(G)$: if $f \in L^2(G)$, $\widehat{f}(\lambda) \in HS(\mathcal{H}_{\lambda})$ and Plancherel formula $\int_G |f(x)|^2 dx = c_0 \int_{\widehat{G}} \|\widehat{f}(\lambda)\|_{HS(\mathcal{H}_{\lambda})}^2 |\lambda|^d d\lambda.$
- Inversion formula : with ad-hoc assumptions

$$f(x) = c_0 \int_{\widehat{G}} \operatorname{tr}(\pi_x^{\lambda} \widehat{f}(\lambda)) |\lambda|^d d\lambda.$$

The sublalacian

The Sublaplacian:
$$-\Delta_G = \sum_{1 \le j \le 2d} V_j^2$$
,
 $\widehat{\Delta_G f}(\lambda) = H(\lambda)\widehat{f}(\lambda), \ \lambda \in \widehat{G}, \ f \in \mathcal{S}(G)$ with

$$egin{aligned} \mathcal{H}(\lambda) &:= |\lambda| \sum_{1 \leq j \leq d} (-\partial_{\xi_j}^2 + \xi_j^2) \ \ \ \ if \ \ \lambda \in \mathfrak{z}^* \setminus \{0\}, \ \mathcal{H}((0,\omega)) &= |\omega|^2, \ \ \ \ \ if \ \ \omega \in \mathfrak{v}^*. \end{aligned}$$

- Spectrum of $H(\lambda)$: $|\lambda|(2|\alpha| + d)$, $\alpha \in \mathbb{N}^d$
- Eigenprojectors of $H(\lambda)$:

$$\Pi_n = \sum_{|\alpha|=n} |h_{\alpha}\rangle \langle h_{\alpha}|,$$

with h_{α} Hermite functions $h_{\alpha}(\xi) = \prod_{1 \leq j \leq d} h_{\alpha_j}(\xi_j)$.

Pseudodifferential operators and Egorov Theorem

Pseudodifferential operators on *H*-type groups

• Using Fourier inversion formula for defining an operator:

$$f(x) = c_0 \int_{\hat{G}} \operatorname{tr} \left(\pi_x^{\lambda} \widehat{f}(\lambda) \right) |\lambda|^d d\lambda.$$
$$\operatorname{Op}(\sigma) f(x) = c_0 \int_{\hat{G}} \operatorname{tr} \left(\pi_x^{\lambda} \sigma(x, \lambda) \widehat{f}(\lambda) \right) |\lambda|^d d\lambda, \ \sigma(x, \lambda) \in \mathcal{L}(\mathcal{H}_{\lambda}).$$

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Symbols : σ(x, λ) ∈ S^m(G) iff σ is C[∞]_c in x and "homogeneous of degree m in λ" + some "differentiability" condition in π^λ (difference operators). Note that σ(x, (0, ω)) is a scalar and not an operator.

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- Semi-classical pseudodifferential operators: Let $\sigma \in S^{-\infty}(G)$, $\varepsilon \ll 1$

$$\operatorname{Op}_{\varepsilon}(\sigma)f(x) = c_0 \int_{\widehat{G}} \operatorname{tr}\left(\pi_x^{\lambda}\sigma(x,\varepsilon\cdot\lambda)\widehat{f}(\lambda)\right) |\lambda|^d d\lambda.$$

[Bahouri, FK & Gallagher, 12], [Fischer & Ruzhansky, 16], [FK & Fischer, 18].

The kernel of a semi-classical pseudodifferential operator.

For $f \in \mathcal{S}(G)$,

$$\begin{aligned} \operatorname{Op}_{\varepsilon}(\sigma)f(x) &= c_0 \int_{\widehat{G}} \operatorname{Tr}\left(\pi_x^{\lambda} \sigma(x, \varepsilon \cdot \lambda) \mathcal{F}f(\lambda)\right) |\lambda|^d \, d\lambda \\ &= c_0 \, \varepsilon^{-Q} \int_{G \times \widehat{G}} \operatorname{Tr}\left(\pi_{\delta_{\varepsilon^{-1}}(y^{-1}x)}^{\lambda} \sigma(x, \lambda)\right) f(y) |\lambda|^d \, d\lambda \, dy \\ &= \int_{G} \kappa_x^{\varepsilon}(y^{-1}x) f(y) \, dy. \end{aligned}$$

The convolution kernel of $Op_{\varepsilon}(\sigma)$ is obtained by scaling from $\kappa_{x}(\cdot)$,

$$\kappa_{x}^{\varepsilon}(z) = \varepsilon^{-Q} \kappa_{x} \left(\delta_{\varepsilon^{-1}} z \right)$$

with

$$\kappa_{\mathrm{x}}(z) = c_0 \int_{\hat{G}} \mathrm{Tr}\left(\pi_z^{\lambda} \sigma(x,\lambda)\right) |\lambda|^d d\lambda \text{ i.e. } \mathcal{F}(\kappa_{\mathrm{x}})(\lambda) = \sigma(x,\lambda).$$

Action on $L^2(G)$

The set \mathcal{A}_0 is the set of operator-valued symbols $\sigma(x,\lambda) \in \mathcal{L}(\mathcal{H}_\lambda)$ such that

$$\sigma(x,\lambda) = \mathcal{F}\kappa_x(\lambda) = \int_G \kappa_x(z)(\pi_z^\lambda)^* dz,$$

where $x \mapsto \kappa_x(y)$ is in $\mathcal{C}^{\infty}_c(\mathcal{G}, \mathcal{S}(\mathcal{G}))$.

Proposition

The set A_0 is an algebra and $\exists C > 0$, $\forall \sigma \in A_0$, $\forall \varepsilon > 0$,

$$\|\operatorname{Op}_{\varepsilon}(\sigma)\|_{\mathcal{L}(L^{2}(G))} \leq C \int_{G} \sup_{x \in G} |\kappa_{x}(z)| dz.$$

$$\begin{array}{ll} \text{Proof:} & |\operatorname{Op}_{\varepsilon}(\sigma)f(x)| = \left| \int_{G} f(y)\kappa_{x}^{\varepsilon}(y^{-1}x)dy \right| \\ & \leq \int_{G} |f(y)| \sup_{x_{1} \in G} |\kappa_{x_{1}}^{\varepsilon}(y^{-1}x)| \ dy = |f| * \sup_{x_{1} \in G} |\kappa_{x_{1}}^{\varepsilon}(\cdot)|(x) \end{array}$$

Then Young conv. inequ. and $\|\sup_{x \in G} |\kappa_{x}^{\varepsilon}(\cdot)|\|_{L^{1}(G)} = \|\sup_{x \in G} |\kappa_{x}(\cdot)|\|_{L^{1}(G)}. \end{array}$

Egorov Theorem [FK, Fischer 19]

Let $\psi \in L^2(G)$ and $\theta \in \mathcal{C}^{\infty}_{c}(\mathbb{R})$, $\sigma \in \mathcal{A}$ with $\sigma = 0$ close to $\lambda = 0$ and $Q^{\varepsilon}_{\sigma}(t) = e^{-i\frac{t}{2\varepsilon^{\tau}}\varepsilon^2\Delta_G} \operatorname{op}_{\varepsilon}(\sigma) e^{i\frac{t}{2\varepsilon^{\tau}}\varepsilon^2\Delta_G}.$

• If $[\sigma, H(\lambda)] \neq 0$, then $\int_{\mathbb{R}} \theta(t) (Q_{\sigma}^{\varepsilon}(t)\psi, \psi)_{L^{2}(G)} dt = O(\varepsilon^{\min(\tau, 1)} \|\psi\|^{2}).$

• If $[\sigma, H(\lambda)] = 0$, $\sigma = \prod_n \sigma \prod_n$ • if $\tau \in (0, 2)$, $\int_{\mathbb{R}} \theta'(t) (Q_{\sigma}^{\varepsilon}(t)\psi, \psi) dt = O(\varepsilon^{\min(1, \tau - 2)} \|\psi\|^2)$, • if $\tau = 2$, for all $s \in \mathbb{R}$ (transport) $\int_{\mathbb{R}} \theta(t) (Q_{\sigma}^{\varepsilon}(t)\psi, \psi) dt = \int_{\mathbb{R}} \theta(t+s) \left(Q_{\Phi_n^{-s}(\sigma)}^{\varepsilon}(t)\psi, \psi\right) dt + O(\varepsilon \|\psi\|^2)$,

3 if $\tau > 2$, for all $s \in \mathbb{R}$ (invariance) $\int_{\mathbb{R}} \theta(t) \left(Q_{\sigma}^{\varepsilon}(t)\psi, \psi \right) dt = \int_{\mathbb{R}} \theta(t) \left(Q_{\Phi_{n}^{-s}(\sigma)}^{\varepsilon}\psi, \psi \right) dt + O(\varepsilon^{\min(1,2-\tau)} \|\psi\|^{2}).$

Above,
$$\Phi_n^s(\sigma) = \sigma\left(\exp(\frac{2n+d}{2|\lambda|} \mathcal{Z}^{(\lambda)} s) x, \lambda \right).$$

Proof of the Egorov Theorem

One writes

$$egin{aligned} &arepsilon^ au \, rac{d}{dt} \left(\mathrm{op}_arepsilon(\sigma) \psi^arepsilon(t), \psi^arepsilon(t)) = \left(\mathrm{op}_arepsilon([\sigma, \mathcal{H}(\lambda)]) \psi^arepsilon(t), \psi^arepsilon(t), \psi^arepsilon(t))
ight. \ &+ arepsilon \left(\mathrm{op}_arepsilon(\mathcal{V}.\pi^\lambda(\mathcal{V})\sigma) \psi^arepsilon(t), \psi^arepsilon(t)) - arepsilon^2 \left(\mathrm{op}_arepsilon(\Delta_G \sigma)(t) \psi^arepsilon(t), \psi^arepsilon(t), \psi^arepsilon(t))
ight. \end{aligned}$$

2 facts:

1 There exists $\sigma_1 \in \mathcal{A}$ such that

 $V.\pi^{\lambda}(V)\sigma = [\sigma_1, H(\lambda)].$

2 For this σ_1 , one has

$$\Pi_n\left(V.\pi^{\lambda}(V)\sigma_1 - \frac{1}{2}\Delta_G\sigma\right)\Pi_n = \frac{2n+d}{2|\lambda|}\Pi_n\mathcal{Z}^{(\lambda)}\sigma\Pi_n.$$

Proof of the Egorov Theorem

One writes for $\sigma = \prod_n \sigma \prod_n$

$$\varepsilon^{\tau} \int \theta(t) \frac{d}{dt} \left(\operatorname{op}_{\varepsilon}(\sigma) \psi^{\varepsilon}(t), \psi^{\varepsilon}(t) \right) dt = \\ \varepsilon^{2} \int \theta(t) \left(\operatorname{op}_{\varepsilon} \left(\frac{2n+d}{2|\lambda|} \Pi_{n} \mathcal{Z}^{(\lambda)} \sigma \Pi_{n} \right) \psi^{\varepsilon}(t), \psi^{\varepsilon}(t) \right) dt.$$

2 facts:

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2 For this σ_1 , one has

$$\Pi_n\left(V.\pi^{\lambda}(V)\sigma_1 - \frac{1}{2}\Delta_G\sigma\right)\Pi_n = \frac{2n+d}{2|\lambda|}\Pi_n\mathcal{Z}^{(\lambda)}\sigma\Pi_n.$$

Semi-classical measures and Schrödinger equation

Existence and structure of semi-classical measures

Let $\mathcal{M}_1^+(G \times \widehat{G})$ be the set of pairs (Γ, γ) consisting of a positive Radon measure γ on $G \times \widehat{G}$ and a family of positive trace 1 operators $\Gamma(x, \lambda)$ on \mathcal{H}_{λ} .

Proposition

Let $(\psi^{\varepsilon}(t))$ be a bounded family in $L^{\infty}(\mathbb{R}, L^{2}(\mathbb{R}^{d}))$. There exist $\varepsilon_{k} \xrightarrow[k \to +\infty]{} 0$ and $t \mapsto \Gamma_{t} d\gamma_{t} \in L^{\infty}(\mathbb{R}, \mathcal{M}_{1}^{+}(G \times \widehat{G}))$ such that for all $\theta \in L^{1}(\mathbb{R})$ and $\sigma \in \mathcal{A}_{0}$,

$$\int \theta(t) \left(\operatorname{Op}_{\varepsilon}(\sigma) \psi^{\varepsilon}(t), \psi^{\varepsilon}(t) \right) dt \underset{\varepsilon_k \to 0}{\longrightarrow} \int_{\mathbb{R} \times G \times \widehat{G}} \theta(t) \operatorname{Tr} \left(\sigma(x, \lambda) \Gamma_t(x, \lambda) \right) d\gamma_t(x, \lambda) dt.$$

Remark:

- **1.** $\Gamma_t d\gamma_t$ has two parts:
 - a non commutative part described by an operator valued measure on $G \times \mathfrak{z}^* \setminus \{0\}$ $(\mathcal{H}_{\lambda} = L^2(\mathfrak{p}_{\lambda})).$
 - an Euclidean part described by a (scalar) positive Radon measure on v^{*} (H_λ = C).

Existence and structure of semi-classical measures

Proposition

Let $(\psi^{\varepsilon}(t))$ be a bounded family in $L^{\infty}(\mathbb{R}, L^{2}(\mathbb{R}^{d}))$. There exist $\varepsilon_{k} \xrightarrow[k \to +\infty]{} 0$ and $t \mapsto \Gamma_{t} d\gamma_{t} \in L^{\infty}(\mathbb{R}, \mathcal{M}_{1}^{+}(G \times \widehat{G}))$ such that for all $\theta \in L^{1}(\mathbb{R})$ and $\sigma \in \mathcal{A}$, $\int \theta(t) (\operatorname{Op}_{\varepsilon}(\sigma)\psi^{\varepsilon}(t), \psi^{\varepsilon}(t)) dt \xrightarrow[\varepsilon_{k} \to 0]{} \int_{\mathbb{R} \times G \times \widehat{G}} \theta(t) \operatorname{Tr}(\sigma(x, \lambda)\Gamma_{t}(x, \lambda)) d\gamma_{t}(x, \lambda) dt.$

Remark:

2. Link with the weak limit of the energy density: If $\|(-\varepsilon^2 \Delta_G)^{\frac{s}{2}} \psi_0^{\varepsilon}\|_{L^2(G)} \leq C$, then $\forall \phi \in \mathcal{C}^{\infty}_c(G), \ \theta \in \mathcal{C}(\mathbb{R}),$ $\lim_{k \to +\infty} \int \theta(t) \int_G \phi(x) |\psi^{\varepsilon_k}(t,x)|^2 dx dt = \int_{\mathbb{R} \times G \times \widehat{G}} \theta(t) \phi(x) \operatorname{Tr}(\Gamma(x,\lambda)) d\gamma(x,\lambda)$

A few words about the proof...

- Difficulty : No Garding's inequality...
 ⇒ Use a "C*-algebra" approach.
- The key: Consider the C*-algebra \mathcal{A} obtained by completion of \mathcal{A}_0 with

$$\sigma \mapsto \sup_{(x,\pi)\in G\times \hat{G}} \|\sigma(x,\pi)\|_{\mathcal{L}(\mathcal{H}_{\pi})}.$$

The set of the states of \mathcal{A} coincide with $\mathcal{M}_1^+(G \times \hat{G})$.

• The arguments: The quantities

$$\ell_{\varepsilon}(\sigma) = (\operatorname{Op}_{\varepsilon}(\sigma)\psi^{\varepsilon}, \psi^{\varepsilon})$$

satisfy:

- Solution For any σ ∈ A₀, ℓ_ε(σ) is bounded and there exists (ε_k(σ))_{k∈N} such that ℓ_{ε_k(σ)}(θ, σ) has a limit ℓ(σ).
- **2** Using the separability of \mathcal{A}_0 and a diagonal extraction, one finds $(\varepsilon_k)_{k\in\mathbb{N}}$ such that for all $\sigma \in \mathcal{A}_0$, $(\ell_{\varepsilon_k}(\sigma))_{k\in\mathbb{N}}$ has a limit $\ell(\sigma)$.
- **3** The map $\sigma \mapsto \ell(\sigma)$ extends to \mathcal{A} and is a state of \mathcal{A} ($\ell(\sigma^*\sigma) \ge 0$).

Semi-classical measures and Schrödinger equation

Consider solutions of the Schrödinger equation

 $i\varepsilon^{ au}\partial_t\psi^{arepsilon}=-\varepsilon^2\Delta_G\psi^{arepsilon},\ \psi^{arepsilon}(0)=\psi^{arepsilon}_0.$

Let $\Gamma_t d\gamma_t$ be the semi-classical measure of a family $\psi^{\varepsilon}(t)$.

• The non-commutative part

Theorem (FK & Fischer 2019)

(i) For
$$d\gamma_t$$
-a. a. $(x,\lambda) \in G imes \widehat{G}$ and a.a. $t \in \mathbb{R}$,

$$\Gamma_t(x,\lambda) = \sum_{n \in \mathbb{N}} \Gamma_{n,t}(x,\lambda)$$
 with $\Gamma_{n,t}(x,\lambda) := \Pi_n \Gamma_t(x,\lambda) \Pi_n$

where Π_n are spectral projections of $H(\lambda)$ for the eigenvalues $|\lambda|(2n + d)$.

(ii) - if $\tau \in (0, 2)$, $\partial_t (\Gamma_{n,t}(x, \lambda) d\gamma_t) = 0$,

- if $\tau = 2$, $\Gamma_{n,t}(x,\lambda)d\gamma_t(x,\lambda)$ satisfies $(\partial_t - \frac{2n+d}{2|\lambda|}\mathcal{Z}^{(\lambda)})(\Gamma_{n,t}(x,\lambda)d\gamma_t) = 0$

- if $\tau > 2$, $\mathcal{Z}^{(\lambda)}(\Gamma_{n,t}(x,\lambda)d\gamma_t) = 0$ (invariance).

 $\mathcal{Z}^{(\lambda)} \in \mathfrak{z}$ is the vector corresponding to $\lambda \in \mathfrak{z}^*$.

Semi-classical measures and Schrödinger equation

- The non-commutative part (...)
- The Euclidean part

Theorem (FK & Fischer 2019)

(iii) Above $\lambda = 0$, set $d\varsigma_t(x, \omega) = \Gamma_t(x, (0, \omega)) d\gamma_t(x, (0, \omega))$

- if $\tau \in (0,1)$, $\partial_t \varsigma_t(x,(0,\omega)) = 0$,
- if $\tau = 1$, $(\partial_t \omega \cdot V)\varsigma_t(x, \omega) = 0$ (transport),
- if $\tau > 1$, $\omega \cdot V\varsigma_t = 0$ (invariance).

Proof: Same strategy than the proof of Egorov Theorem + limit $\varepsilon \rightarrow 0$ + identification of each part of the measure.

- We have extended the microlocal/semiclassical approach to the setting of graded Lie groups.
- In the non-semiclassical framework, one obtains (without pain...) compensated compactness theorems ([Baldi, Franchi, 13], [Baldi, Franchi, Tchou & Tesi 10], [FK & Fischer 18])
- Application to the analysis of eigenfunctions of (complicated) sub-Laplacians ?

Thank you for your attention !