# Semi-classical analysis of Schrödinger equation on H-type groups

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[Folland & Stein 1974], [Rotschild & Stein 1976], [Folland, 1977]: Analysis of sublaplacians can be performed using Lie groups theory, lifting procedure.

**Example**: Heisenberg group and the operator  $L$  on  $\mathbb{R}^3$ :

$$
L = X^2 + Y^2, \ \ X = \partial_x - \frac{y}{2}\partial_z, \ \ Y = \partial_y + \frac{x}{2}\partial_z, \ \ [X, Y] = \partial_z
$$

What would be a semiclassical/microlocal approach on graded Lie groups using theory of representations ? Which class of pseudodifferential operators ?

[Taylor 1984], [Beals & Greiner 1988], [Christ, Geller, Glowasky & Polin 1992] [Geller 1990] [Bahouri, FK & Gallagher 2012], [Fischer & Ruzhansky 2016], [FK & Fischer, 2018 & 2019]

# The question: Large time evolution of energy density of families of solutions of a Schrödinger equation

Let  $V_j$  be 2d vector fields of  $\mathbb{R}^{2d+p}$  and

$$
-\Delta_G = \sum_{1 \leq j \leq 2d} V_j^2.
$$

Assume  $\mathfrak{g} = \mathfrak{v} \oplus \mathfrak{z}$  is a Lie algebra of Heisenberg type with centre  $\mathfrak{z}$ ,  $\mathfrak{v} = \mathrm{Vect}(V_i, 1 \leq j \leq 2d\}, \ \ \mathfrak{z} = \mathrm{Vect}([V_i, V_j], 1 \leq i, j \leq 2d),$  $2d = \text{dim}\mathfrak{v}, \ \ \mathfrak{p} = \text{dim}\mathfrak{z}.$ 

Let  $(\psi_0^\varepsilon)_{\varepsilon>0}$  be bounded in  $L^2(\mathbb{R}^d)$  such that  $\exists s,\, \mathcal{C}_s>0,\, \forall \varepsilon>0,$  $\varepsilon^{\mathcal{S}} \| (-\Delta_{\mathcal{G}} )^{\frac{s}{2}} \psi_0^\varepsilon \|_{L^2(\mathcal{G})} + \varepsilon^{-s} \| (-\Delta_{\mathcal{G}} )^{-\frac{s}{2}} \psi_0^\varepsilon \|_{L^2(\mathcal{G})} \leq \mathcal{C}_s.$ 

Describe for  $\phi \in C_c^{\infty}(G)$ ,  $T \in \mathbb{R}$ ,  $\tau > 0$ 

$$
\lim_{\varepsilon\to 0}\frac{1}{T}\int_0^T\int_G \phi(x)|\mathrm{e}^{i\frac{t}{2\varepsilon^{\tau}}\varepsilon^2\Delta_G}\psi_0^{\varepsilon}(x)|^2dx\,dt.
$$

#### The same question in the Euclidean case

Let  $(\psi_0^\varepsilon)$  be a bounded family in  $L^2(\mathbb{R}^d)$  satisfying

 $\exists s, \mathcal{C}_s>0, \;\; \forall \varepsilon >0, \;\; \varepsilon^{s} \| (-\Delta)^{\frac{s}{2}} \psi_0^\varepsilon \|_{L^2(\mathbb{R}^d)} + \varepsilon^{-s} \| (-\Delta)^{-\frac{s}{2}} \psi_0^\varepsilon \|_{L^2(\mathbb{R}^d)} \leq \mathcal{C}_c.$ 

Then any limit point of the measure

$$
\left| e^{-i\frac{t}{2\varepsilon^{\tau}}\varepsilon^2\Delta}\psi_0^{\varepsilon} \right|^2 dx dt
$$

is of the form  $\varrho_t(\mathsf{x})dt$  where  $\varrho_t \in \mathcal{M}^+(\mathbb{R}^d)$  and

- **11** If  $\tau \in (0, 1)$ , then  $\rho_t = \rho_0$ .
- **2** If  $\tau = 1$  then  $\varrho_t(x) = \int_{\mathbb{R}^d} \mu_0(x t\xi, d\xi)$ .
- **3** If  $\tau > 1$  then  $\rho_t = 0$ .

### The result

#### Theorem (FK & Fischer 2019)

Any weak limit of  $|{\rm e}^{i\frac{t}{2\varepsilon^\tau}\varepsilon^2\Delta_\mathsf{G}}\psi^\varepsilon_0(x)|^2d\mathsf{x}\,dt$  writes  $\varrho_tdt=\left(\varrho_t^{\mathfrak{v}^*}+\varrho_t^{\mathfrak{z}^*}\right)$  $\left(\begin{smallmatrix} \mathfrak z^* \ t \end{smallmatrix}\right)$  dt with

• If  $\tau \in (0,1)$ , for all  $t \in \mathbb{R}$ ,  $\rho_t = \rho_0$ If  $\tau = 1$ , then  $\varrho_t^{\mathfrak z^*} = \varrho_0^{\mathfrak z^*}$  $\frac{3}{0}^*$  and  $\varrho_t^{\mathfrak{v}^*}$  $\int_{t}^{\mathfrak{v}^*}(x) = 1$  $\int_{\mathfrak{v}^*} \varsigma_0\left( \text{Exp}(t\,\omega\cdot V)\times\,,\mathit{d}\omega\right).$ If  $\tau \in (1,2)$ , then  $\varrho_t^{v^*} = 0$  and  $\partial_t \varrho_t^{3^*} = 0$  holds in  $\mathcal{D}'(\mathbb{R} \times G)$ . If  $\tau=2$ , then  $\varrho_t^{\mathfrak{v}^*}=0$  and  $\varrho_t^{\mathfrak{z}^*}=\sum$ n∈N Z  $\int\limits_{\mathfrak z^*\setminus\{0\}}\gamma_{n,t}(\mathsf x,d\lambda)$  where  $\left(\partial_t - \frac{2n + d}{2\lambda}\right)$  $\frac{n+d}{2|\lambda|} \mathcal{Z}^{(\lambda)}\bigg)\, \gamma_{n,t} = 0,$ 

where  $\mathcal{Z}(\lambda)$  is the left invariant vector field corresponding to  $\lambda \in \mathfrak{z}^*$ . • If  $\tau > 2$ , then  $\rho_t = 0$  for all  $t \in \mathbb{R}$ .

#### Remarks

- Quantum limits split into two parts, with different threshold indexes.
- **•** Dispersion takes longer than in the Euclidean case Bahouri, Gérard & Xu 2000], [Del Hierro 2005], [Bahouri, FK, Gallagher 2016]
- Splitting and invariance properties already noticed ( $p = 1$  contact manifolds, Grauert tubes) in [Zelditch 1997], [Colin de Verdière, Hillairet & Trélat 2018], [Burq & Sun 2019]
- **1** Graded, stratified Lie groups, H-type groups
- <sup>2</sup> Analysis on graded Lie groups and Fourier transform
- <sup>3</sup> Pseudodifferential operators, Egorov theorem
- 4 Semi-classical measures and Schrödinger equation.

Graded Lie groups

# Graded Lie groups - Definition

#### **Definition**

A simply connected Lie group G is graded if the Lie algebra  $\frak g$  of its left-invariant vector fields is graded :

$$
\mathfrak{g}=\mathfrak{g}_1\oplus\cdots\oplus\mathfrak{g}_n,\ \ [\mathfrak{g}_\ell,\mathfrak{g}_{\ell'}]\subset\mathfrak{g}_{\ell+\ell'}.
$$

The group G is stratified if

$$
\forall \ell \in \{1, \cdots n-1\}, \ \ [\mathfrak{g}_{\ell}, \mathfrak{g}_{1}]=\mathfrak{g}_{\ell+1}.
$$

• G identifies to g via the exponential map:

$$
\begin{array}{rccc}\n\exp & : & \mathfrak{g} & \rightarrow & \mathsf{G} \\
X & \mapsto & \exp(X)e_{\mathsf{G}}\n\end{array}
$$

 $\bullet$  The law group on G is a polynomial map (Campbell-Baker-Hausdorff).

## Graded Lie groups - Heisenberg group

- **The Heisenberg group:** 
	- The Lie algebra  $h = v \oplus \lambda$  is generated by

$$
X = \partial_x - \frac{y}{2}\partial_z
$$
,  $Y = \partial_y + \frac{x}{2}\partial_z$ ,  $Z = \partial_z = [X, Y]$ .

 $\implies$  The group  $\mathbb H$  is stratified and has two steps.

 $\bullet$  The points of G are the elements

$$
w = \mathrm{Exp}(xX + yY + sS), \ \ (x, y, s) \in \mathbb{R}^3.
$$

 $\mathbb{H}=\mathbb{R}^3$  with the product law

$$
(x,y,s)\cdot(x',y',s')=(x+x',y+y',s+s'+\frac{1}{2}(x\cdot y'-x'\cdot y)).
$$

• The center of  $\mathbb H$  is the set of points of the form  $(0, 0, s)$ .

# Graded Lie groups - H-type groups

 $\bullet$  H-type groups (multidimensional versions of  $\mathbb H$  [Kaplan 80])

**1** The Lie algebra of G is stratified with two steps  $\mathfrak{g} = \mathfrak{v} \oplus \mathfrak{z}$ . **2** If  $B(\lambda)$  is the skew symmetric form defined on  $v \times v$  by

 $\forall \lambda \in \mathfrak{z}^*, \forall U, V \in \mathfrak{v}, \quad B(\lambda)(U, V) = \lambda([U, V]),$ 

then  $B(\lambda)^2 = -|\lambda|^2 \mathrm{Id}$ .

• Notation (1): Let  $(V_1, \ldots, V_{2d})$  be an ONB of v and  $(Z_1, \cdots, Z_p)$  of  $\chi$ , then  $x \in G$  or  $X \in \mathfrak{g}$  writes

 $x = \text{Exp}(X)$ ,  $X = v_1V_1 + ... + v_{2d}V_{2d} + z_1Z_1 + ... + z_nZ_n$ .

## Graded Lie groups - H-type groups

• Notation (2): v decomposes in a  $\lambda$ -depending way as  $\mathfrak{v} = \mathfrak{p}_{\lambda} \oplus \mathfrak{q}_{\lambda}$  with

$$
\mathfrak{p} := \mathfrak{p}_{\lambda} := \mathsf{Span} (P_1, \ldots, P_d), \quad \mathfrak{q} := \mathfrak{q}_{\lambda} := \mathsf{Span} (Q_1, \ldots, Q_d).
$$

where  $(P_1, \ldots, P_d, Q_1, \ldots, Q_d)$  be an ONB of v such that

$$
B(\lambda)(U, V) = |\lambda|U^tJV, \ \ J = \begin{pmatrix} 0 & \mathrm{Id} \\ -\mathrm{Id} & 0 \end{pmatrix}.
$$

Then  $x \in G$  or  $X \in \mathfrak{g}$  write

 $x = \text{Exp}(X)$ ,  $X = p_1P_1 + ... + p_dP_d + q_1Q_1 + ... + q_dQ_d + z_1Z_1 + ... + z_nZ_n$ .

**Besides** 

$$
[P_j,Q_j]=\mathcal{Z}^{\lambda},\ \forall j\in\{1,\cdots,d\}
$$

where  $\mathcal{Z}^{\lambda}$  corresponds to the vector  $\lambda$  in the identification  $\mathfrak{z}\sim\mathfrak{z}^*.$ 

# Analysis on H-type groups

# Analysis on H-type groups

 $\bullet$  The Haar measure, dx is deduced on G from the Lebesgue meas. on g. Associated Lebesgue spaces on G:

$$
||f||_{L^p(G)} := \left(\int_G |f(x)|^p\ dx\right)^{\frac{1}{p}},\quad p\in [1,\infty[.
$$

The dilatation. For  $r > 0$  and  $X = V + Z \in \mathfrak{g}$ , one sets  $\delta_r X = rV + r^2 Z$ .



 $r \cdot \text{Exp}(V + Z) := \text{Exp}(r V + r^2 Z), r > 0.$ 

=⇒ Homogeneous functions and operators

• Homogeneous dimension. The Haar measure is Q-homogeneous:

 $Q := \dim v + 2 \dim x = 2d + 2p$ .

Let  $\hat{G}$  be the set of irreducible representations of  $G$ :

 $\widehat{G} = {\text{class of }\pi^{\lambda} : \lambda \in \mathfrak{z}^* \setminus \{0\}} \sqcup {\text{class of }\pi^{0,\omega} : \omega \in \mathfrak{v}^*}.$ 

Infinite dimensional representations are parametrized by  $\mathfrak z^*\setminus\{0\}$ : for  $\lambda\in \mathfrak z^*\setminus\{0\}$ , then  $\mathcal H_\lambda = L^2(\mathfrak p_\lambda)$  with for  $x=\mathrm{Exp}(P+Q+Z)\in \mathsf G$ ,

$$
\pi_{x}^{\lambda}\Phi(\xi)=\exp\left[i\lambda(Z)+\frac{i}{2}|\lambda|P\cdot Q+i\sqrt{|\lambda|}\xi\cdot Q\right]\Phi\left(\xi+\sqrt{|\lambda|}P\right).
$$

Finite dimensional representations are parametrized by  $v^*$ : for  $\omega \in v^*$ ,

$$
\pi_x^{0,\omega} = e^{i\omega(V)}, \quad x = \text{Exp}(V + Z) \in G.
$$

# The dual set of H-type groups

Plancherel measure:  $d\mu(\pi^{\lambda}) = |\lambda|^d d\lambda$ .

 $\implies$  The set of finite dimensional representations is of 0 Plancherel measure.

- Dilations on  $\hat{G} \cdot r \cdot \pi_{\rm v} = \pi_{\rm v} \cdot r$ 
	- for  $r \in \mathbb{R}_+^*$ ,  $\lambda \in \mathfrak{z}^* \setminus \{0\}$ ,  $x \in \mathcal{G}$ ,

$$
\pi_{r\cdot x}^{\lambda} = T_r \pi_x^{r^2 \lambda} T_r^*, \quad T_r f(\xi) = r^{1/2} f(r\xi),
$$

 $\implies$   $r \cdot \lambda = r^2 \lambda$ 

for  $r \in \mathbb{R}_+^*$ ,  $\omega \in \mathfrak{v}^*$ ,  $x \in \mathcal{G}$ ,  $\pi_{r \cdot x}^{0,\omega} = e^{ir\omega \cdot x}$ 

 $\implies$   $r \cdot (0, \omega) = (0, r\omega)$ 

### Fourier transform

#### **Definition**

Let  $f\in L^1(G)$ , the Fourier transform of  $f$  is the operator of  $\mathcal{L}(\mathcal{H}_\lambda)$ ,

$$
\widehat{f}(\lambda)=\int_G f(x)\pi^{\lambda}(x)^*dx.
$$

• Notation: 
$$
\widehat{f}(\lambda) = \widehat{f}(\pi^{\lambda}), \ \widehat{f}(0,\omega) = \widehat{f}(\pi^{0,\omega}).
$$

Extension to  $L^2(G)$  : if  $f \in L^2(G)$ ,  $\hat{f}(\lambda) \in HS(H_\lambda)$  and Plancherel formula

$$
\int_G |f(x)|^2 dx = c_0 \int_{\hat{G}} ||\hat{f}(\lambda)||_{HS(\mathcal{H}_{\lambda})}^2 |\lambda|^d d\lambda.
$$

• Inversion formula : with ad-hoc assumptions

$$
f(x) = c_0 \int_{\hat{G}} \text{tr}(\pi_x^{\lambda} \widehat{f}(\lambda)) |\lambda|^d d\lambda.
$$

# The sublalacian

The Sublaplacian: 
$$
-\Delta_G = \sum_{1 \le j \le 2d} V_j^2
$$
,  
\n
$$
\widehat{\Delta_G f}(\lambda) = H(\lambda) \widehat{f}(\lambda), \ \lambda \in \widehat{G}, \ f \in \mathcal{S}(G)
$$
\nwith

$$
H(\lambda) := |\lambda| \sum_{1 \leq j \leq d} (-\partial_{\xi_j}^2 + \xi_j^2) \text{ if } \lambda \in \mathfrak{z}^* \setminus \{0\},
$$
  

$$
H((0,\omega)) = |\omega|^2, \text{ if } \omega \in \mathfrak{v}^*.
$$

- Spectrum of  $H(\lambda)$ :  $|\lambda|(2|\alpha|+d)$ ,  $\alpha \in \mathbb{N}^d$
- **•** Eigenprojectors of  $H(\lambda)$ :

$$
\Pi_n = \sum_{|\alpha|=n} |h_{\alpha}\rangle \langle h_{\alpha}|,
$$

with  $h_\alpha$  Hermite functions  $h_\alpha(\xi)=\Pi_{1\leq j\leq d}h_{\alpha_j}(\xi_j).$ 

# Pseudodifferential operators and Egorov Theorem

### Pseudodifferential operators on H-type groups

Using Fourier inversion formula for defining an operator:

$$
f(x) = c_0 \int_{\hat{G}} \operatorname{tr} \left( \pi_x^{\lambda} \hat{f}(\lambda) \right) |\lambda|^d d\lambda.
$$
  
Op( $\sigma$ ) $f(x) = c_0 \int_{\hat{G}} \operatorname{tr} \left( \pi_x^{\lambda} \sigma(x, \lambda) \hat{f}(\lambda) \right) |\lambda|^d d\lambda, \ \sigma(x, \lambda) \in \mathcal{L}(\mathcal{H}_{\lambda}).$ 

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Symbols :  $\sigma(x, \lambda) \in S^m(G)$  iff  $\sigma$  is  $\mathcal{C}^\infty_c$  in  $x$  and "homogeneous of degree  $m$ in  $\lambda^{\prime\prime}$  + some "differentiability" condition in  $\pi^\lambda$  (difference operators). Note that  $\sigma(x,(0,\omega))$  is a scalar and not an operator.

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- Semi-classical pseudodifferential operators: Let  $\sigma\in S^{-\infty}(G)$ ,  $\varepsilon\ll 1$

$$
\mathrm{Op}_{\varepsilon}(\sigma)f(x)=c_0\int_{\hat{G}}\mathrm{tr}\left(\pi_{x}^{\lambda}\sigma(x,\varepsilon\cdot\lambda)\widehat{f}(\lambda)\right)|\lambda|^{d}d\lambda.
$$

[Bahouri, FK & Gallagher, 12], [Fischer & Ruzhansky, 16], [FK & Fischer, 18].

# The kernel of a semi-classical pseudodifferential operator.

For  $f \in \mathcal{S}(G)$ ,

$$
Op_{\varepsilon}(\sigma) f(x) = c_0 \int_{\widehat{G}} \operatorname{Tr} \left( \pi_x^{\lambda} \sigma(x, \varepsilon \cdot \lambda) \mathcal{F} f(\lambda) \right) |\lambda|^{d} d\lambda
$$
  
\n
$$
= c_0 \varepsilon^{-Q} \int_{G \times \widehat{G}} \operatorname{Tr} \left( \pi_{\delta_{\varepsilon^{-1}}(y^{-1}x)}^{\lambda} \sigma(x, \lambda) \right) f(y) |\lambda|^{d} d\lambda dy
$$
  
\n
$$
= \int_{G} \kappa_x^{\varepsilon} (y^{-1}x) f(y) dy.
$$

The convolution kernel of  ${\rm Op}_\varepsilon(\sigma)$  is obtained by scaling from  $\kappa_\mathsf{x}(\cdot)$ ,

$$
\kappa_{\mathsf{x}}^{\varepsilon}(z) = \varepsilon^{-Q} \kappa_{\mathsf{x}} \left( \delta_{\varepsilon^{-1}} z \right)
$$

with

$$
\kappa_x(z) = c_0 \int_{\hat{G}} \text{Tr} \left( \pi_z^{\lambda} \sigma(x, \lambda) \right) |\lambda|^d d\lambda \text{ i.e. } \mathcal{F}(\kappa_x)(\lambda) = \sigma(x, \lambda).
$$

# Action on  $L^2(G)$

The set  $\mathcal{A}_0$  is the set of operator-valued symbols  $\sigma(x, \lambda) \in \mathcal{L}(\mathcal{H}_\lambda)$  such that

$$
\sigma(x,\lambda)=\mathcal{F}\kappa_x(\lambda)=\int_G \kappa_x(z)(\pi_z^{\lambda})^*dz,
$$

where  $x \mapsto \kappa_x(y)$  is in  $C_c^{\infty}(G, S(G))$ .

#### Proposition

The set  $A_0$  is an algebra and  $\exists C > 0$ ,  $\forall \sigma \in A_0$ ,  $\forall \varepsilon > 0$ ,

$$
\|\mathrm{Op}_{\varepsilon}(\sigma)\|_{\mathcal{L}(L^2(G))}\leq C\,\int_G\sup_{x\in G}|\kappa_x(z)|dz.
$$

**Proof:** 
$$
|Op_{\varepsilon}(\sigma)f(x)| = \left| \int_{G} f(y)\kappa_{x}^{\varepsilon}(y^{-1}x)dy \right|
$$
  
\n
$$
\leq \int_{G} |f(y)| \sup_{x_1 \in G} |\kappa_{x_1}^{\varepsilon}(y^{-1}x)| dy = |f| * \sup_{x_1 \in G} |\kappa_{x_1}^{\varepsilon}(.)|(x)
$$
\nThen Young conv. inequ. and  $||sup_{x \in G} |\kappa_{x}^{\varepsilon}(.)||_{L^{1}(G)} = ||sup_{x \in G} |\kappa_{x}(\cdot)||_{L^{1}(G)}$ .

# Egorov Theorem [FK, Fischer 19]

Let  $\psi\in L^2(\mathsf{G})$  and  $\theta\in\mathcal{C}^\infty_c(\mathbb{R}),\ \sigma\in\mathcal{A}$  with  $\sigma=0$  close to  $\lambda=0$  and  $Q_{\sigma}^{\varepsilon}(t) = e^{-i\frac{t}{2\varepsilon^{\tau}}\varepsilon^{2}\Delta_{\mathcal{G}}} \exp_{\varepsilon}(\sigma) e^{i\frac{t}{2\varepsilon^{\tau}}\varepsilon^{2}\Delta_{\mathcal{G}}}.$ 

If  $[\sigma, H(\lambda)] \neq 0$ , then  $\displaystyle \int_{\mathbb{R}} \theta(t) \left( Q^{\varepsilon}_{\sigma}(t) \psi, \psi \right)_{L^{2}(G)} dt = O(\varepsilon^{\min(\tau,1)} \Vert \psi \Vert^{2}).$ 

- **•** If  $[\sigma, H(\lambda)] = 0$ ,  $\sigma = \prod_{p} \sigma \prod_{p}$ 
	- **1** if  $\tau \in (0, 2)$ ,  $\int\limits_{\mathbb R} \theta'(t)\left(Q^\varepsilon_\sigma(t)\psi,\psi\right)dt = O(\varepsilon^{\min(1,\tau-2)}\|\psi\|^2),$ 2 if  $\tau = 2$ , for all  $s \in \mathbb{R}$  (transport)

Z  $\int\limits_{\mathbb R} \theta(t) \left( Q^\varepsilon_\sigma(t) \psi, \psi \right) dt = \int\limits_{\mathbb R}$  $\int\limits_{\mathbb R} \theta(t+s) \, \Big( Q^\varepsilon_\Phi$  $\underset{\Phi_{n}^{-s}(\sigma)}{\varepsilon}(t)\psi,\psi\Big)\,dt + O(\varepsilon\|\psi\|^{2}),$ 

**3** if  $\tau > 2$ , for all  $s \in \mathbb{R}$  (invariance) Z  $\int\limits_{\mathbb R} \theta(t) \left( Q^\varepsilon_\sigma(t) \psi, \psi \right) dt = \int\limits_{\mathbb R}$  $\int\limits_{\mathbb R} \theta(t) \, \Big( Q_\Phi^\varepsilon$  $\left(\Phi_{\sigma}^{-s}(\sigma)\psi,\psi\right)dt+O(\varepsilon^{\min(1,2-\tau)}\|\psi\|^{2}).$ 

Above, 
$$
\Phi_n^s(\sigma) = \sigma \left( \exp(\frac{2n+d}{2|\lambda|} \mathcal{Z}^{(\lambda)} s) x, \lambda \right).
$$

# Proof of the Egorov Theorem

One writes

$$
\varepsilon^{\tau} \frac{d}{dt} \left( \operatorname{op}_{\varepsilon}(\sigma) \psi^{\varepsilon}(t), \psi^{\varepsilon}(t) \right) = \left( \operatorname{op}_{\varepsilon}([\sigma, H(\lambda)]) \psi^{\varepsilon}(t), \psi^{\varepsilon}(t) \right) \quad + \varepsilon \left( \operatorname{op}_{\varepsilon}(V. \pi^{\lambda}(V)\sigma) \psi^{\varepsilon}(t), \psi^{\varepsilon}(t) \right) - \varepsilon^{2} \left( \operatorname{op}_{\varepsilon}(\Delta_{\mathcal{G}} \sigma)(t) \psi^{\varepsilon}(t), \psi^{\varepsilon}(t) \right).
$$

2 facts:

**1** There exists  $\sigma_1 \in \mathcal{A}$  such that

$$
V.\pi^{\lambda}(V)\sigma=[\sigma_1,H(\lambda)].
$$

**2** For this  $\sigma_1$ , one has

$$
\Pi_n\left(V.\pi^\lambda(V)\sigma_1-\frac{1}{2}\Delta_G\sigma\right)\Pi_n=\frac{2n+d}{2|\lambda|}\Pi_n\mathcal{Z}^{(\lambda)}\sigma\Pi_n.
$$

# Proof of the Egorov Theorem

One writes for  $\sigma = \prod_{p} \sigma \prod_{p}$ 

$$
\varepsilon^{\tau} \int \theta(t) \frac{d}{dt} \left( \operatorname{op}_{\varepsilon}(\sigma) \psi^{\varepsilon}(t), \psi^{\varepsilon}(t) \right) dt =
$$
  

$$
\varepsilon^{2} \int \theta(t) \left( \operatorname{op}_{\varepsilon} \left( \frac{2n + d}{2|\lambda|} \Pi_{n} \mathcal{Z}^{(\lambda)} \sigma \Pi_{n} \right) \psi^{\varepsilon}(t), \psi^{\varepsilon}(t) \right) dt.
$$

#### 2 facts:

**1** There exists  $\sigma_1 \in \mathcal{A}$  such that

$$
V.\pi^{\lambda}(V)\sigma=[\sigma_1,H(\lambda)].
$$

2 For this  $\sigma_1$ , one has

$$
\Pi_n\left(V.\pi^\lambda(V)\sigma_1-\frac{1}{2}\Delta_G\sigma\right)\Pi_n=\frac{2n+d}{2|\lambda|}\Pi_n\mathcal{Z}^{(\lambda)}\sigma\Pi_n.
$$

# Semi-classical measures and Schrödinger equation

## Existence and structure of semi-classical measures

Let  $\mathcal{M}_1^+(G\times \widehat{G})$  be the set of pairs  $(\Gamma,\gamma)$  consisting of a positive Radon measure  $\gamma$  on  $G \times \widehat{G}$  and a family of positive trace 1 operators  $\Gamma(x, \lambda)$  on  $\mathcal{H}_{\lambda}$ .

#### Proposition

Let  $(\psi^\varepsilon(t))$  be a bounded family in  $L^\infty(\mathbb R,L^2(\mathbb R^d)).$  There exist  $\varepsilon_k \underset{k \to +\infty}{\longrightarrow} 0$  and  $t \mapsto \Gamma_t d\gamma_t \in L^{\infty}(\mathbb{R},\mathcal{M}_1^+(G\times \widehat{G}))$  such that for all  $\theta \in L^1(\mathbb{R})$  and  $\sigma \in \mathcal{A}_0$ ,

$$
\int \theta(t) \left(\mathrm{Op}_{\varepsilon}(\sigma)\psi^{\varepsilon}(t),\psi^{\varepsilon}(t)\right)dt \underset{\varepsilon_k\to 0}{\longrightarrow} \int_{\mathbb{R}\times G\times \widehat{G}} \theta(t) \mathrm{Tr}\left(\sigma(x,\lambda)\Gamma_t(x,\lambda)\right) d\gamma_t(x,\lambda)dt.
$$

#### Remark:

- 1.  $\Gamma_t d\gamma_t$  has two parts:
	- a non commutative part described by an operator valued measure on  $G \times \mathfrak{z}^* \setminus \{0\}$   $(\mathcal{H}_\lambda = L^2(\mathfrak{p}_\lambda)).$
	- an Euclidean part described by a (scalar) positive Radon measure on  $v^*$  $(\mathcal{H}_\lambda = \mathbb{C}).$

### Existence and structure of semi-classical measures

#### Proposition

Let  $(\psi^\varepsilon(t))$  be a bounded family in  $L^\infty(\mathbb R,L^2(\mathbb R^d)).$  There exist  $\varepsilon_k \underset{k \to +\infty}{\longrightarrow} 0$  and  $t \mapsto \Gamma_t d\gamma_t \in L^{\infty}(\mathbb{R},\mathcal{M}_1^+(G\times \widehat{G}))$  such that for all  $\theta \in L^1(\mathbb{R})$  and  $\sigma \in \mathcal{A}$ ,

 $\int \theta(t) \,(\operatorname{Op}_{\varepsilon}(\sigma)\psi^{\varepsilon}(t),\psi^{\varepsilon}(t))\,dt \mathop{\longrightarrow}\limits_{{\varepsilon_{k}\rightarrow 0}}$  $\mathbb{R}\times G\times \widehat{G}$  $\theta(t)\text{Tr}\left(\sigma(x,\lambda)\Gamma_t(x,\lambda)\right)d\gamma_t(x,\lambda)dt.$ 

#### Remark:

2. Link with the weak limit of the energy density: If  $\|(-\varepsilon^2\Delta_G)^{\frac{s}{2}}\psi_0^{\varepsilon}\|_{L^2(G)}\leq C$ , then  $\forall\phi\in\mathcal{C}^\infty_c(G)$ ,  $\theta\in\mathcal{C}(\mathbb{R})$ , lim sup<br> $k\rightarrow+\infty$  $\int \theta(t)$ G  $\phi(x)|\psi^{\varepsilon_k}(t,x)|^2 dxdt =$  $\mathbb{R}\times G\times \widehat{G}$  $\theta(t)\phi(x) \text{Tr}\left(\mathsf{\Gamma}(x,\lambda)\right)d\gamma(x,\lambda)$ 

## A few words about the proof...

- **O Difficulty** : No Garding's inequality...  $\implies$  Use a "C\*-algebra" approach.
- The key: Consider the  $C^*$ -algebra  $\mathcal A$  obtained by completion of  $\mathcal A_0$  with

$$
\sigma \mapsto \sup_{(x,\pi)\in G\times \hat{G}} \|\sigma(x,\pi)\|_{\mathcal{L}(\mathcal{H}_{\pi})}.
$$

The set of the states of  ${\mathcal A}$  coincide with  ${\mathcal M}_1^+(G\times \hat G).$ 

**o** The arguments: The quantities

$$
\ell_\varepsilon(\sigma) = (\operatorname{Op}_\varepsilon(\sigma) \psi^\varepsilon, \psi^\varepsilon)
$$

satisfy:

- **1** For any  $\sigma \in A_0$ ,  $\ell_{\varepsilon}(\sigma)$  is bounded and there exists  $(\varepsilon_k(\sigma))_{k\in\mathbb{N}}$  such that  $\ell_{\epsilon_k (\sigma)}(\theta, \sigma)$  has a limit  $\ell(\sigma)$ .
- 2 Using the separability of  $A_0$  and a diagonal extraction, one finds  $(\varepsilon_k)_{k\in\mathbb{N}}$  such that for all  $\sigma\in\mathcal{A}_0$ ,  $(\ell_{\varepsilon_k}(\sigma))_{k\in\mathbb{N}}$  has a limit  $\ell(\sigma).$
- **3** The map  $\sigma \mapsto \ell(\sigma)$  extends to  ${\cal A}$  and is a state of  ${\cal A}$   $(\ell(\sigma^*\sigma) \geq 0).$

## Semi-classical measures and Schrödinger equation

Consider solutions of the Schrödinger equation

 $i\varepsilon^{\tau}\partial_t\psi^{\varepsilon}=-\varepsilon^2\Delta_G\psi^{\varepsilon}, \ \ \psi^{\varepsilon}(0)=\psi^{\varepsilon}_0.$ 

Let  $\Gamma_t d\gamma_t$  be the semi-classical measure of a family  $\psi^\varepsilon(t).$ 

The non-commutative part

Theorem (FK & Fischer 2019)

(i) For 
$$
d\gamma_t
$$
-a. a.  $(x, \lambda) \in G \times \widehat{G}$  and a.a.  $t \in \mathbb{R}$ ,

$$
\Gamma_t(x,\lambda)=\sum_{n\in\mathbb{N}}\Gamma_{n,t}(x,\lambda) \text{ with }\Gamma_{n,t}(x,\lambda):=\Pi_n\Gamma_t(x,\lambda)\Pi_n,
$$

where  $\Pi_n$  are spectral projections of  $H(\lambda)$  for the eigenvalues  $|\lambda|(2n+d)$ .

- (ii) if  $\tau \in (0, 2)$ ,  $\partial_t (\Gamma_{n,t}(x, \lambda) d\gamma_t) = 0$ ,
	- if τ = 2, Γn,t(x, λ)dγt(x, λ) satisfies (∂<sup>t</sup> − 2n+d <sup>2</sup>|λ<sup>|</sup> Z (λ) ) (Γn,t(x, λ)dγt) = 0

- if  $\tau>$  2 ,  $\, {\cal Z}^{(\lambda)} \,(\Gamma_{n,t}(\mathsf{x},\mathsf{\lambda}) \mathsf{d} \gamma_t) = 0 \,$  (invariance).

 $\mathcal{Z}^{(\lambda)} \in \mathfrak{z}$  is the vector corresponding to  $\lambda \in \mathfrak{z}^*.$ 

## Semi-classical measures and Schrödinger equation

- The non-commutative part (...)
- The Euclidean part

#### Theorem (FK & Fischer 2019)

(iii) Above  $\lambda = 0$ , set  $d\varsigma_t(x,\omega) = \Gamma_t(x,(0,\omega))d\gamma_t(x,(0,\omega))$ 

- $-$  if  $\tau \in (0,1)$ ,  $\partial_t \varsigma_t(x,(0,\omega)) = 0$ ,
- if  $\tau = 1$ ,  $(\partial_t \omega \cdot V)_{S_t}(x, \omega) = 0$  (transport),
- if  $\tau > 1$ ,  $\omega \cdot V \zeta_t = 0$  (invariance).

**Proof**: Same strategy than the proof of Egorov Theorem + limit  $\varepsilon \to 0$  + identification of each part of the measure.

- We have extended the microlocal/semiclassical approach to the setting of graded Lie groups.
- **In the non-semiclassical framework, one obtains (without pain...)** compensated compactness theorems ([Baldi, Franchi, 13], [Baldi, Franchi, Tchou & Tesi 10], [FK & Fischer 18])
- Application to the analysis of eigenfunctions of (complicated) sub-Laplacians ?

Thank you for your attention !