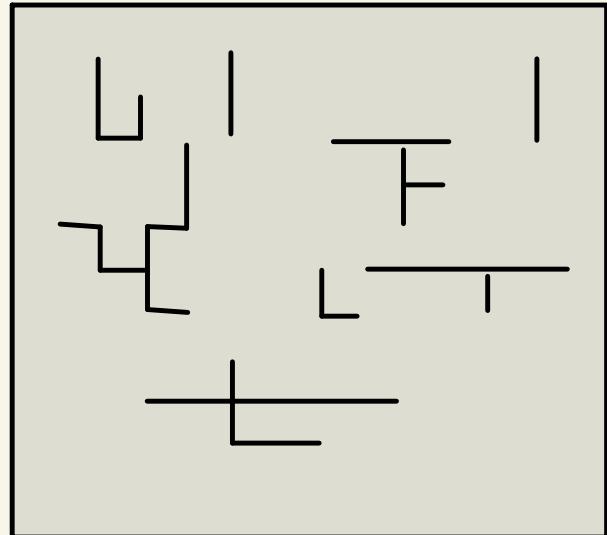


Random forests and non-linear sigma models

(and what these have to
do with random matrices)



Roland Bauerschmidt (Cambridge)

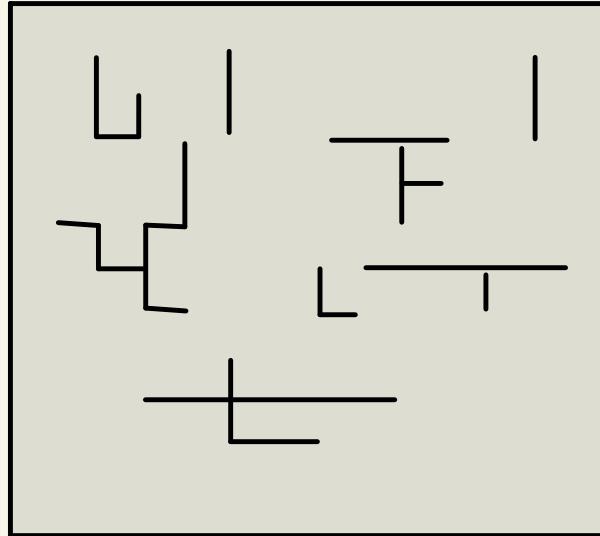
with N. Crawford, T. Helmuth, A. Swan

The arboreal gas

$G = (\Lambda, E)$ finite graph

$F = (\Lambda, E(F)) \subset G$ is a forest if it has no cycles

$$P_\beta^G(F) = \frac{1}{Z_\beta} \beta^{|E(F)|} \mathbf{1}(F \text{ is a forest})$$



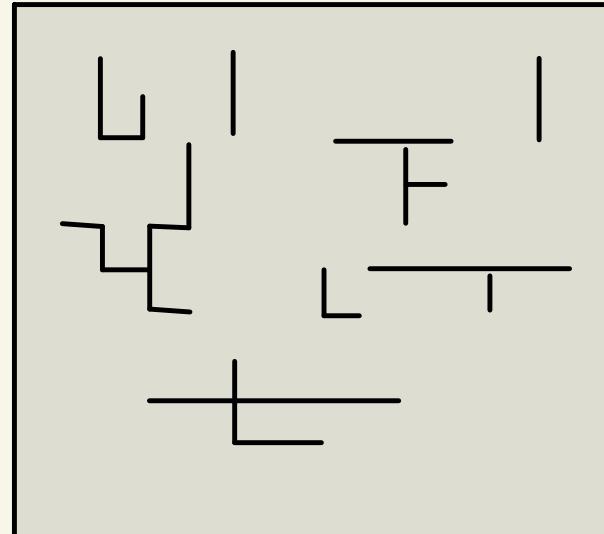
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unrooted



The arboreal gas

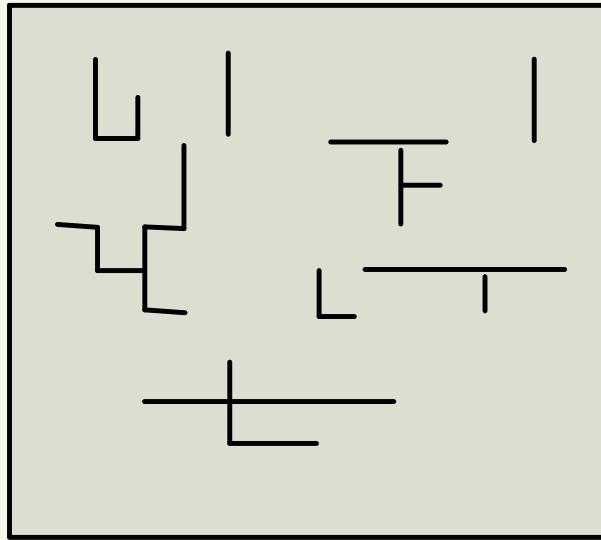
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- Fortuin-Kasteleyn: $q \rightarrow 0$ limit of random cluster model
- Lubensky-Isaacson: model for gelation of polymers
- Caracciolo-Jacobsen-Saleur-Sokal-Sportiello: SUSY
- Luczak-Pittel, Martin-Yeo: analysis on complete graph
- Pemantle, Kahn, Grimmett-Winkler: Conj. negative dependence
- Brändén-Huh, Anari et al.: a main example of Lorentzian polyn.

When does the arboreal gas percolate?



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| Prop. On any graph, the arboreal gas is stochastically dominated by bond percolation with $p = \beta/(1+\beta)$.

→ Arboreal gas does not percolate on \mathbb{Z}^d if β small

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→ Arboreal gas does not percolate on \mathbb{Z}^d if β small

| Thm (BCHS). In $d=2$, no percolation for any $\beta > 0$.

| Thm (BCH). In $d \geq 3$, percolation for $\beta > \beta_0$.

When does the arboreal gas percolate?

Thm. (BCHS). Let $d=2$. For any $\beta > 0$, $\exists c(\beta) > 0$ s.t.
 $P_\beta^\Lambda(0 \leftrightarrow x) \leq O(|x|^{-c(\beta)})$ for all $x \in \mathbb{Z}^2$.

Thm. (BCH). Let $d \geq 3$. There are $\beta_0 \in (0, \infty)$ and $\zeta(\beta) = (-O(1/\beta))$ s.t. for $\beta \geq \beta_0$,

$$P_\beta^\Lambda(0 \leftrightarrow x) = \zeta(\beta) + \frac{c(\beta)}{\beta |x|^{d-2}} + O\left(\frac{1}{\beta |x|^{d-2+\kappa}} + \frac{1}{\beta L^{3\kappa}}\right)$$

if Λ is a torus $\Lambda = \mathbb{Z}^d / L \mathbb{Z}^d$, $L \geq L_0(d)$.

Matrix tree theorem

delete a row and column

$$\#\text{spanning trees on } G = \det(-\Delta^0)$$

$$P_{UST}(e_1 \in T, \dots, e_k \in T) = \det(K(e_i, e_j))_{i,j}$$

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$\xi_x, \eta_x, x \in \Lambda$ anticommuting variables

$$e^{(\xi - \Delta^0 \eta)}$$

expand $e^{\sum \xi_x (-\Delta^0)_{xy} \eta_y}$

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$\xi_x, \gamma_x, x \in \Lambda$ anticommuting variables

$$\#ST = \int \partial_\gamma \partial_\xi e^{(\xi - \Delta^0 \gamma)}$$

Grassmann int.

project onto top degree coefficient

expand $e^{\sum \xi_x (-\Delta^0)_{xy} \gamma_y}$

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$\xi_x, \gamma_x, x \in \Lambda$ anticommuting variables

$$\#ST = \int \partial_\eta \partial_\xi e^{(\xi - \Delta^0 \gamma)} \text{ Grassmann int.}$$

$$P^{UST}(e_1 \in T, \dots, e_k \in T) = \left\langle \prod_{i=1}^k (\nabla_{e_i} \xi)(\nabla_{e_i} \gamma) \right\rangle_{FF}$$

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Nonlinear σ -model (NLSM)

Spins u_i , $i \in \Lambda$

$$\langle F \rangle_\beta = \frac{1}{Z_\beta} \int F e^{-\frac{\beta}{2} \sum_{ij \in E} (u_i - u_j) \cdot (u_i - u_j)} \prod_i |T_i| du_i$$

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Haar measure

Ising model:

$$u_i \in \{\pm 1\}$$

XY model:

$$u_i \in S^1$$

Heisenberg model:

$$u_i \in S^2$$

Hyperbolic σ -models: $u_i \in \mathbb{H}^n$

spins take values in
a symmetric space

Nonlinear σ -model (NLSM)

Spins u_i , $i \in \Lambda$

$$\langle F \rangle_{\beta, h} = \frac{1}{Z_{\beta, h}} \int F e^{-\frac{\beta}{2} \sum_{ij \in E} (u_i - u_j) \cdot (u_i - u_j)} e^{-h \sum_{i \in \Lambda} z_i} \prod_i du_i$$

$z_j = u_j \cdot e$, some fixed direction e

Ising model:

$$u_i \in \{\pm 1\}$$

XY model:

$$u_i \in S^1$$

Heisenberg model: $u_i \in S^2$

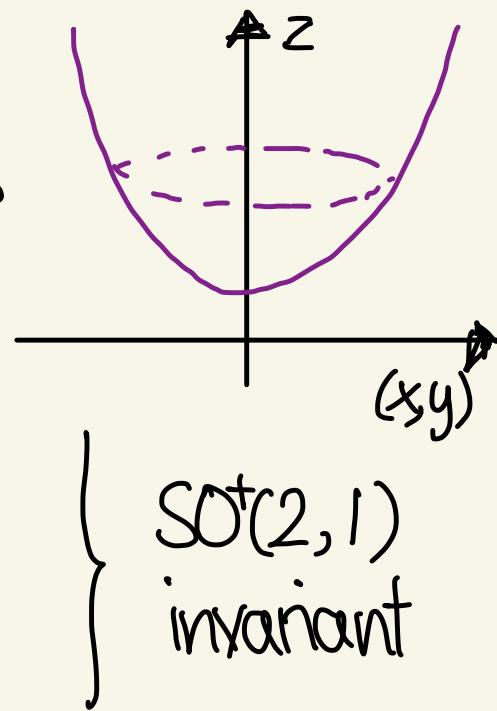
Hyperbolic σ -models: $u_i \in H^n$

Hyperbolic plane \mathbb{H}^2

Vectors $u_i = (x_i, y_i, z_i)$ in \mathbb{R}^3
 \uparrow
 $z_i > 0$

with $u_i \cdot u_i = -1$

where $u_i \cdot u_j = x_i x_j + y_i y_j - z_i z_j$



$\left. \begin{array}{l} \text{SO}^+(2,1) \\ \text{invariant} \end{array} \right\}$

Concretely $z_i = \sqrt{1 + x_i^2 + y_i^2}$

$$\int_{\mathbb{H}^2} F = \int_{\text{Lebesgue integral}} dx_i dy_i \frac{1}{z_i} F$$

$\text{SO}(2,1)$ inv. volume form

Fermionic hyperbolic plane $\mathbb{H}^{0|2}$

Supervectors $u_i = (\underbrace{\xi_i}_{\text{odd}}, \underbrace{\gamma_i}_{\text{even}}, \underbrace{z_i}_{\text{even}})$ in $\mathbb{R}^{1|2}$

$$\text{with } u_i \cdot u_i = -1$$

$$\text{where } u_i \cdot u_j = \gamma_i \xi_j + \gamma_j \xi_i - z_i z_j$$

$\left. \begin{array}{l} \text{OSp}(1|2) \\ \text{invariant} \end{array} \right\}$

$$\text{Concretely } z_i = \sqrt{1 - 2 \xi_i \gamma_i} = 1 - \xi_i \gamma_i$$

$$\int_{\mathbb{H}^{0|2}} F = \int \partial \gamma_i \partial \xi_i \frac{1}{z_i} F$$

↑
Grassmann integral
↑
 $\text{OSp}(1|2)$ inv. volume form

$\mathbb{H}^{0|2}$ model \rightarrow arboreal gas

$$\langle F \rangle_{\beta, h} \propto \int_{(\mathbb{H}^{0|2})^\Lambda} e^{\frac{\beta}{2}(u, \Delta u) - h \sum_i z_i} F$$

$(u, -\Delta u) = \sum_{ij} (u_i - u_j) \cdot (u_i - u_j)$
 \uparrow
OSp(1|2) inner prod.

(Λ finite lattice)

Thm. (CJSSS, BCHS)

Nonlinear Matrix Tree Thm.

$$P_{\beta,0}(0 \leftrightarrow x) = -\langle u_0 \cdot u_x \rangle_{\beta,0}$$



$$P_{\beta,0}(e_1 \in F, \dots, e_k \in F) = \left(-\frac{\beta}{2}\right)^k \left\langle \prod_{i=1}^k (\nabla_{e_i} u) \cdot (\nabla_{e_i} u) \right\rangle_{\beta,0}$$

⋮

H^{012} model \rightarrow arboreal gas

$$P_{\beta,0}(0 \leftrightarrow x) = -\langle u_0 \cdot u_x \rangle_{\beta,0}$$

\Rightarrow percolation \leftrightarrow spontaneously broken symmetry

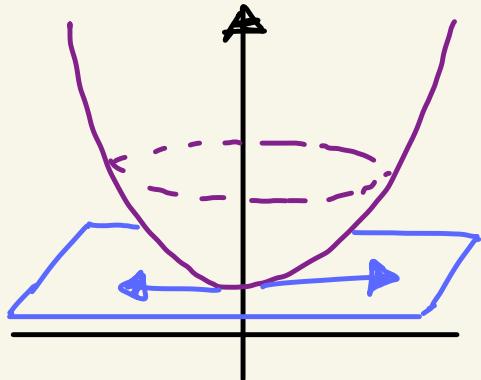
"spins point in same direction over large dist."

H^{O12} model \rightarrow arboreal gas

$$P_{\beta,0}(0 \leftrightarrow x) = -\langle u_0 \cdot u_x \rangle_{\beta,0}$$

\Rightarrow percolation \leftrightarrow spontaneously broken symmetry

"spins point in same direction over large dist."



Continuous symmetry
 \rightarrow expect diffusive corrections
(free field)

Ward identity: $\frac{\langle z_i \rangle_{\beta,h}}{h} = \sum_j \langle z_i \eta_j \rangle_{\beta,h}$

'Magic formula' \rightarrow No percolation in $d=2$

Mermin-Wagner theorem: "Continuous symmetries cannot be spontaneously broken in $2d$ ".

\rightarrow Does not directly apply to superspaces, but...

Ihm (BCHS). For the arboreal gas,

$$P_\beta(0 \leftrightarrow x) \propto \int e^{tx} e^{-\sum_{i \neq j} \beta \cosh l t_i - t_j} \left(\det(-\Delta_{\beta(t)}^0) \right)^{3/2} \prod_{i \neq 0} e^{-3t_i} dt_i.$$

'Magic formula'
(c.f. Sotov-Tarnes,
Diaconis)

\rightarrow Versions of Mermin-Wagner can be adapted to RHS.

Review of Anderson transition

$$\Lambda_L = \mathbb{Z}^d / L\mathbb{Z}^d, d \geq 2$$

$$H_L = -\Delta + \lambda V, V \text{ i.i.d. (or other finite range RM)}$$

DOS: $S_{E,h} = \mathbb{E} \operatorname{Im}(H-E-i\hbar)^{-1}(0,0)$

2-pt funct: $\tau_{E,h}(j) = \mathbb{E} |(H-E-i\hbar)^{-1}(0,j)|^2, j \in \Lambda$

Ward identity: $\frac{1}{\hbar} S_{E,h} = \sum_j \tau_{E,h}(j)$

Non-linear σ -model: $\frac{1}{\hbar} \langle Z_0 \rangle_{\beta,h} = \sum_j \langle X_0 X_j \rangle_{\beta,h}$

Review of Anderson transition

$$\Lambda_L = \mathbb{Z}^d / L\mathbb{Z}^d, d \geq 2$$

$$H_L = -\Delta + \lambda V, V \text{ i.i.d. Gaussian}$$

DOS: $S_{E,h} = \mathbb{E} \operatorname{Im}(H-E-i\hbar)^{-1}(0,0)$

2-pt funct: $\tau_{E,h}(x) = \mathbb{E} |(H-E-i\hbar)^{-1}(0,x)|^2$

Localisation: $\tau_{E,h}(x) \lesssim \frac{1}{h} e^{-|x|/\ell}$ (FS: $\lambda \gg 1$)

Delocalisation: $\tau_{E,h}(0) \lesssim 1$ as $L \rightarrow \infty$ then $h \downarrow 0$

Diffusion: $\tau_{E,h}(x) \sim D(E) |x|^{-(d-2)}$ (Conj.: $\lambda \ll 1, S_E > 0$)

Many questions

Universality: Is the order statistics of the component sizes on $\mathbb{Z}^d / L\mathbb{Z}^d$, $d \geq 3$, $\beta > \beta_0$, the same as on the complete graph?

Expected similarly to universality of Wigner-Dyson.

Distributions known explicitly for complete graph:

$$(T^{(1)} - \alpha N) N^{-2/3} \rightarrow \text{limiting distribution}$$

k-th largest component

$$T^{(k)} N^{-2/3} \rightarrow \text{limiting distribution}$$

Many questions

Negative correlation conjecture. (Remanette, Grimmett-Winkler,
Kahn)

$$P_\beta(e_1 \in F, e_2 \in F) \leq P_\beta(e_1 \in F) P_\beta(e_2 \in F)$$

Known only in limit $\beta \rightarrow \infty$ (UST).

Bränden - Huh proved

$$P_\beta(e_1 \in F, e_2 \in F) \leq 2 P_\beta(e_1 \in F) P_\beta(e_2 \in F)$$

and the partition function with edge dependent weights is log concave (Lorentzian polynomial).

Thank you!

Magic formula for ERRW and VRSP

ERRW: Edge e has weight $\alpha + N_e(t)$

parameter

of crossings of edge e

α small: strong reinforcement
 α large: weak reinforcement

Magic formula for ERRW and VRSP

ERRW: Edge e has weight $\alpha + N_e(t)$

Diaconis-Freedman

& Coppersmith-Diaconis:

walk starting
at vertex O

$$P_{\text{ERRW}}(\alpha) = \int P_{\text{SRW}}(c) d\mu_\alpha(c)$$

random conductances
mixing measure
given by 'magic formula'

Magic formula for ERRW and VRSP

ERRW: Edge e has weight $\alpha + N_e(t)$

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$$P_0^{\text{ERRW}(\alpha)} = \int_{\Omega} P_0^{\text{SRW}(c)} d\mu_\alpha(c)$$

random conductances

walk starting at vertex 0

mixing measure given by 'magic formula'

VRSP: Vertex reinforced version

Sabot - Tarrès:

$$P_0^{\text{VRSP}(\beta)} = \int_{\Omega} P_0^{\text{SRW}(ct)} d\nu_\beta(t)$$

Thm (Sabot-Tarrès). Mixing measure for VRJP:

$$d\nu_\beta(t) \propto e^{-\sum_{i \neq j} \beta \cosh(t_i - t_j)} (\det^o(-\Delta_{\beta(t)}))^{1/2} \prod_{i \neq 0} e^{-t_i} dt_i$$

$(\Delta_{\beta(t)} f)_i = \sum_{j \neq i} \beta e^{t_i + t_j} (f_j - f_i)$

Thm (BCHS). For the arboreal gas,

$$P_\beta(0 \leftrightarrow x) \propto \int e^{tx} e^{-\sum_{i \neq j} \beta \cosh(t_i - t_j)} (\det^o(-\Delta_{\beta(t)}))^{3/2} \prod_{i \neq 0} e^{-3t_i} dt_i.$$

Three sources for the magic.

Hyperbolic σ -model

$$\propto e^{-\frac{\beta}{2} \sum_{ij \in E} (u_i - u_j) \cdot (u_i - u_j)} \prod_i du_i$$
$$= e^{-\beta \sum_{ij \in E} (\cosh(t_i - t_j) + \frac{1}{2} e^{t_i + t_j} (s_i - s_j)^2)} \prod_i T_i e^{-(n-1)t_i} dt_i ds_i$$

↑
Gaussian in s !

Hyperbolic σ -model

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Thus t -marginal is

$$e^{-\sum_{ij} \beta \cosh(t_i - t_j) (\det^o(-\Delta_{\beta(t)}))^{-1}} \prod_{i \neq 0} T_i e^{-(n-1)t_i} dt_i$$

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$n=0$

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$n=-2$

What is H^n if $n \leq 0$?

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$n=-2$

What is H^n if $n \leq 0$? H^{p+2q} has dim. $n=p-2q$.