

Gibbsian line ensembles and β -corners processes

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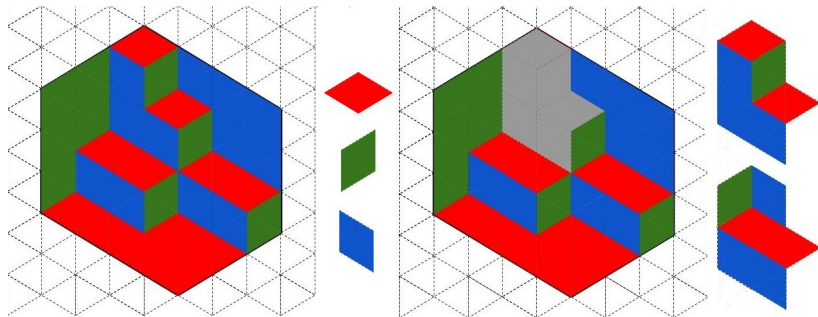
Workshop on Integrable Structures in Random Matrix Theory and Beyond
– MSRI, Berkeley

Outline

- 1 Gibbsian line ensembles
- 2 Convergence of Gibbsian line ensembles
- 3 Maximal free energy in the log-gamma polymer
- 4 β -corners processes

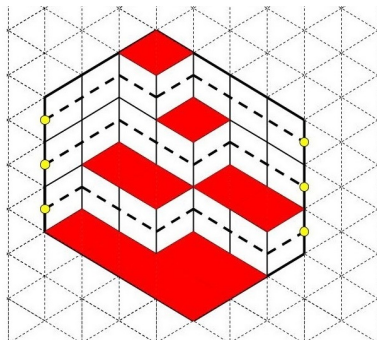
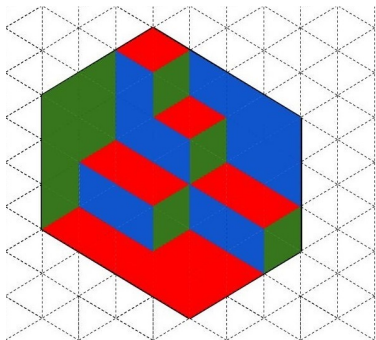
Gibbs measures

Gibbs property = internal consistency condition of a random model.
Uniform lozenge tilings of the hexagon



Gibbsian line ensembles

Line ensemble (LE) = finite or countably infinite collection of random continuous curves, defined on the same probability space.

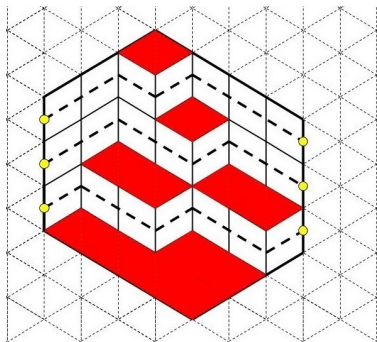
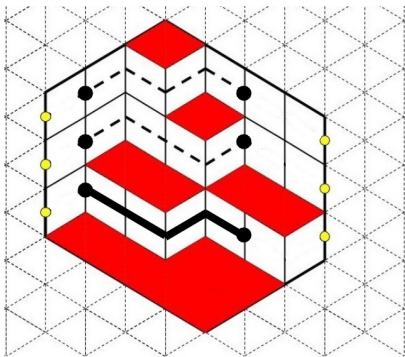


Tiling Gibbs property \implies LE Gibbs property.

LE Gibbs property = locally avoiding Bernoulli random walks (or bridges).

Gibbsian line ensembles

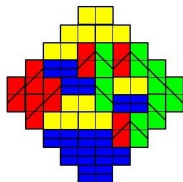
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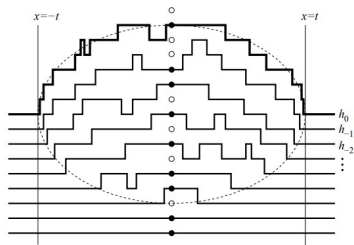
LE Gibbs property = locally avoiding Bernoulli random walks (or bridges).

Gibbsian line ensembles

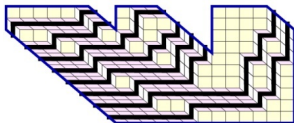
Gibbsian line ensembles = finite or countably infinite collection of random walk trajectories with *local* interactions



Domino tilings of the Aztec diamond [Johansson '02]



Multi-layer PNG model [Prähofer-Spohn '02]



Lozenge tilings of polygons [Petrov '14]

Gibbsian LEs appear in **random tilings**, **last passage percolation** and **directed random polymers**.

Asymptotics of Gibbsian line ensembles

Q: What happens when to a Gibbsian line ensemble $\{\mathcal{L}_i^N\}_{i=1}^N$ as $N \rightarrow \infty$?

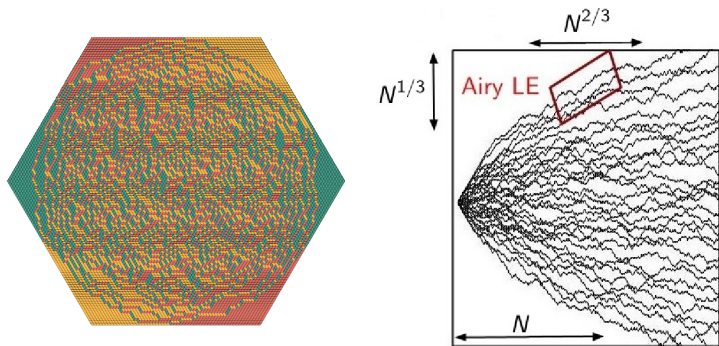
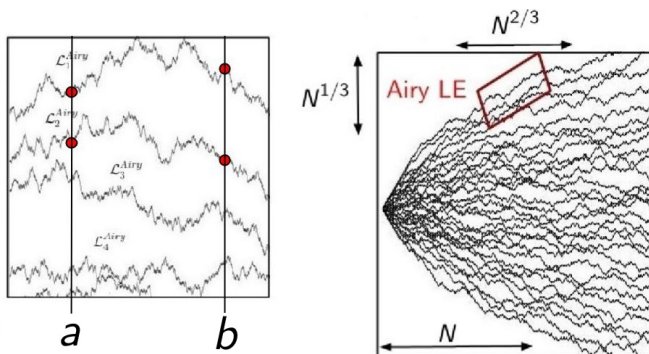


Figure: Simulation due to L. Petrov

We enter the [Kardar-Parisi-Zhang \(KPZ\) universality class](#)

Limiting object: [\(Parabolic\) Airy line ensemble](#) $\{\mathcal{L}_i^{Airy}\}_{i=1}^{\infty}$

The parabolic Airy line ensemble



(Parabolic) Airy line ensemble $\{\mathcal{L}_i^{Airy}\}_{i=1}^\infty$

\mathcal{L}_1^{Airy} = (parabolic) Airy process $\mathcal{L}_1^{Airy}(0) =$ GUE Tracy-Widom dist.

$\{\mathcal{L}_i^{Airy}\}_{i=1}^\infty$ has the **Brownian Gibbs property** (locally avoiding Brownian bridges)

Key questions for Gibbsian line ensembles

1 Tightness

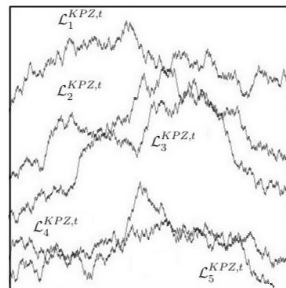
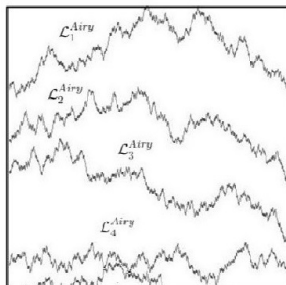
2 Characterization

3 Convergence

- Zero temperature: $\{\mathcal{L}_i^{\text{Airy}}\}_{i=1}^{\infty}$
- Positive temperature: $\{\mathcal{L}_i^{\text{KPZ},t}\}_{i=1}^{\infty}$

4 Properties

5 Applications



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- 1 Gibbsian line ensembles
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Framework for proving convergence to the Airy LE

To prove $\mathcal{L}^N \xrightarrow{\text{u.c.}} \mathcal{L}^{\text{Airy}}$ one needs

- 1 Show that $\mathcal{L}^N \xrightarrow{\text{f.d.}} \mathcal{L}^{\text{Airy}}$.
- 2 Show that \mathcal{L}^N is tight

Zero temperature

- PNG [Prähofer-Spohn '02]
- Domino tilings Aztec diamond [Johansson '02]
- Schur processes [Okounkov-Reshetikhin '03]
- TASEP [Johansson '03]
- Expon. LPP [Borodin-Péché '08]
- Brownian watermelons [Corwin-Hammond '14]

Positive temperature Until 2020 most work was on $\mathcal{L}_1^N \xrightarrow{\text{1.p.}} \mathcal{L}_1^{\text{Airy}}$:

- ASEP [Tracy-Widom '08]
- KPZ eqn. [Sasamoto-Spohn '10]
- Macdonald processes [Borodin-Corwin '14]
- Log-gamma polymer [Borodin-Corwin-Remenik '13]
- S6V [Borodin-Corwin-Gorin '16]

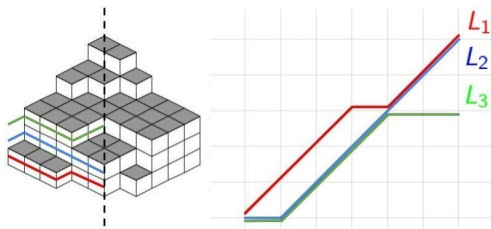
Until 2020: $\mathcal{L}_1^N \xrightarrow{\text{f.d.}} \mathcal{L}_1^{\text{Airy}}$

- Physics: **replica approach** (non-rigorous)
- Math: [Nguyen-Zygouras '16] (incomplete)

Finite dimensional convergence

Theorem (D. '20)

Let $\mathcal{L}^N = \{L_i^N\}_{i=1}^N$ be the *Hall-Littlewood Gibbsian LE* (M, N, a, t) . The two-point distribution of L_1^N converges to the two-point distribution of $\mathcal{L}_1^{\text{Airy}}$.



- 1 Result is limited to two points and parameter restrictions
- 2 When reflected L_1^N has the law of the height function of the *stochastic six-vertex model* [Borodin-Bufetov-Wheeler '16]
- 3 The *first* multi-point convergence result for a *non-determinantal (positive temperature) model*. Softer techniques were later developed by [Quastel-Sarkar '20] (*ASEP* and *KPZ*) and [Virág '20] (*polymer models*).

Finite dimensional convergence

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Use the method of *Macdonald difference operators* [Borodin-Corwin '14]

$$\mathbb{E} \left[\frac{1}{(u_1 t^{-L_1^N(n_1)}; t)_\infty} \frac{1}{(u_2 t^{-L_1^N(n_2)}; t)_\infty} \right] = \sum_{N_1=0}^{\infty} \sum_{N_2=0}^{\infty} I_M(N_1, N_2)$$

$$I_M(N_1, N_2) = \frac{1}{N_1! N_2!} \int_{\gamma_1^{N_1}} \int_{\gamma_2^{N_1}} \int_{\gamma_3^{N_2}} \int_{\gamma_4^{N_2}} D(\vec{w}, \vec{z}) G(\vec{w}, \vec{z}, n_1, u_1)$$

$$D(\vec{w}, \vec{z}) G(\vec{w}, \vec{z}, n_2, u_2) \cdot CT(\vec{w}, \vec{z}; \vec{w}, \vec{z}) \prod_{i=1}^{N_2} \frac{d\hat{w}_i}{2\pi t} \prod_{i=1}^{N_2} \frac{d\hat{z}_i}{2\pi t} \prod_{i=1}^{N_1} \frac{dw_i}{2\pi t} \prod_{i=1}^{N_1} \frac{dz_i}{2\pi t}$$

$$D(\vec{w}, \vec{z}) = \det \left[\frac{1}{z_i - w_j} \right]_{i,j=1}^{N_1}, \quad CT = \prod_{i=1}^{N_1} \prod_{j=1}^{N_2} \frac{(\hat{z}_j z_i^{-1}; t)_\infty}{(\hat{w}_j z_i^{-1}; t)_\infty} \frac{(\hat{w}_j w_i^{-1}; t)_\infty}{(\hat{z}_j w_i^{-1}; t)_\infty}$$

$$G(\vec{w}, \vec{z}, n_1, u_1) = \prod_{i=1}^{N_1} \left(\frac{1+az_i}{1+aw_i} \right)^M \left(\frac{1+aw_i^{-1}}{1+az_i^{-1}} \right)^{n_1} \frac{S(\log w_i - \log z_i; u_1, t)}{[-\log t]_{w_i}}$$

$$(a; t)_\infty = \prod_{m=0}^{\infty} (1 - at^m) - t\text{-Pochhammer symbol}$$

Framework for proving convergence to the Airy LE

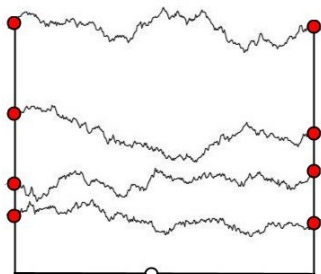
To prove $\mathcal{L}^N \xrightarrow{\text{u.c.}} \mathcal{L}^{\text{Airy}}$ one needs

- 1 Show that $\mathcal{L}^N \xrightarrow{\text{f.d.}} \mathcal{L}^{\text{Airy}}$.
- 2 Show that \mathcal{L}^N is tight

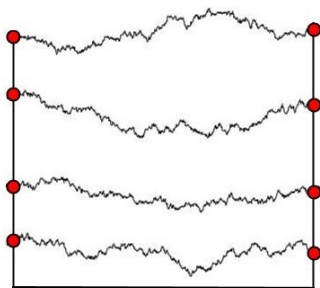
In [D.-Matetski '20] we proposed the following alternate framework.

- 1 Show that $\mathcal{L}_1^N \xrightarrow{\text{f.d.}} \mathcal{L}_1^{\text{Airy}}$ **parabolic Airy process**.
 - 2 Show that \mathcal{L}^N is tight and that all subsequential limits satisfy the **Brownian Gibbs property** (locally avoiding Brownian bridges).
 - 3 Show that $\mathcal{L}^{\text{Airy}}$ is the unique line ensemble which satisfies the Brownian Gibbs property and has the parabolic Airy process as its top curve.
- The framework reduces the quantitative information we need from \mathcal{L}^N to \mathcal{L}_1^N (useful for models that are **non-determinantal**)
 - \mathcal{L}_1^N is frequently special: KPZ line ensemble $\{\mathcal{L}_i^{\text{KPZ},t}\}_{i=1}^\infty$ ($\mathcal{L}_1^{\text{KPZ},t}$ is the **solution to the narrow wedge KPZ equation at time t**)

The Princess and the Pea



The Pea



Q: Can the Princess feel the pea? A: Yes, and more!

Classification of Brownian Gibbsian line ensembles

Theorem (D.-Matetski '20)

Suppose that \mathcal{L}^1 and \mathcal{L}^2 are Brownian Gibbsian line ensembles with laws \mathbb{P}_1 and \mathbb{P}_2 , respectively. Suppose further that for every $k \in \mathbb{N}$, $t_1 < t_2 < \dots < t_k$ and $x_1, \dots, x_k \in \mathbb{R}$ we have

$$\mathbb{P}_1 (\mathcal{L}_1^1(t_1) \leq x_1, \dots, \mathcal{L}_1^1(t_k) \leq x_k) = \mathbb{P}_2 (\mathcal{L}_1^2(t_1) \leq x_1, \dots, \mathcal{L}_1^2(t_k) \leq x_k).$$

Then $\mathbb{P}_1 = \mathbb{P}_2$.

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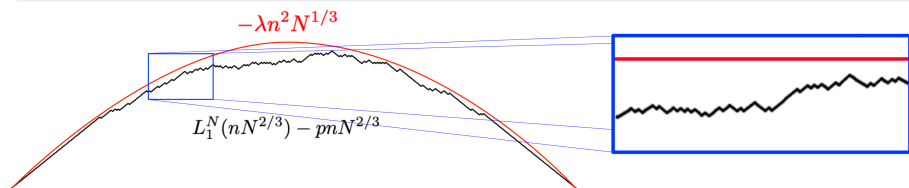
Then $\mathbb{P}_1 = \mathbb{P}_2$.

- 1 Theorem says that a Brownian Gibbsian line ensemble is completely characterized by its top curve. The (parabolic) Airy line ensemble $\{\mathcal{L}_i^{\text{Airy}}\}_{i=1}^\infty$ is characterized by (1) the Brownian Gibbs property and (2) $\mathcal{L}_1^{\text{Airy}} =$ (parabolic) Airy process
- 2 Proof is non-constructive (no formulas for \mathcal{L}_2 etc.)
- 3 Positive temperature analogue recently proved in [D. '21]. Specifically, the KPZ_t line ensemble is characterized by its lowest-indexed curve being the narrow wedge KPZ equation and its Gibbs property (called *H-Brownian Gibbs property*)

Tightness of Gibbsian line ensembles

Theorem (Meta-)

Let $\mathfrak{L}^N = \{L_i^N\}_{i=1}^N$ be a sequence of Gibbsian line ensembles. Suppose that for some constants $p \in \mathbb{R}$ and each $n \in \mathbb{Z}$ the r.v.'s $\frac{L_1^N(nN^{2/3}) - pnN^{2/3}}{N^{1/3}}$ are **tight** and **globally parabolic**. Then $\frac{L_i^N(xN^{2/3}) - pxN^{2/3}}{N^{1/3}}$ are tight in $C(\mathbb{N} \times \mathbb{R})$ and all subsequential limits satisfy the Brownian Gibbs property.



- 1-point tightness + Gibbs property \implies tightness of whole ensemble
- Avoiding Bernoulli random walkers ([proof of concept](#))
[D.-Fang-Fesser-Serio-Teitler-Wang-Zhu '20]
- (H, H^{RW}) -Gibbsian line ensembles ([log-gamma polymer](#)) [D.-Wu '21]
- Key ingredient: [KMT coupling for random walk bridges](#) [D.-Wu '19]

Framework for proving convergence to the Airy LE

- 1 Show that $\mathcal{L}_1^N \xrightarrow{\text{f.d.}} \mathcal{L}_1^{\text{Airy}}$ **parabolic Airy process**.
- 2 Show that \mathcal{L}^N is tight and that all subsequential limits satisfy the **Brownian Gibbs property** (locally avoiding Brownian bridges).
- 3 Show that $\mathcal{L}^{\text{Airy}}$ is the unique line ensemble which satisfies the Brownian Gibbs property and has the parabolic Airy process on top.

Framework was used to show $\{\mathcal{L}_i^{\text{KPZ},t}\}_{i=1}^\infty \implies \{\mathcal{L}_i^{\text{Airy}}\}_{i=1}^\infty$ as $t \rightarrow \infty$.
Step 1 in [Quastel-Sarkar '20], [Virág '20], Step 2 in [Wu '21], and Step 3 in [D.-Matetski '20].

Currently applying it to log-gamma polymer – Step 2 in [D.-Wu '21].

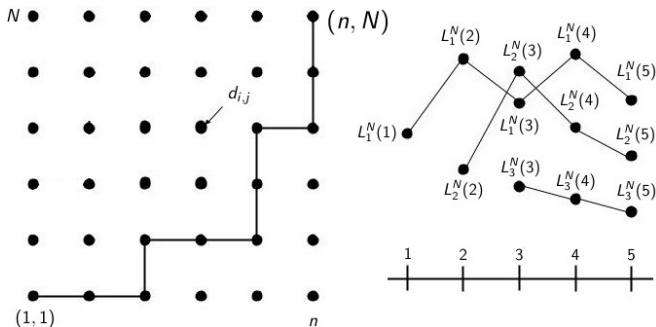
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The log-gamma polymer [Seppäläinen '09]

$d_{i,j}$ are i.i.d. inverse-gamma: $f_\theta(x) = \mathbf{1}\{x > 0\} \Gamma(\theta)^{-1} x^{-\theta-1} \exp(-x^{-1})$.

Partition functions: $Z^N(n) = \sum_{\pi: (1,1) \rightarrow (n,N)} w(\pi)$, $w(\pi) = \prod_{(i,j) \in \pi} d_{i,j}$



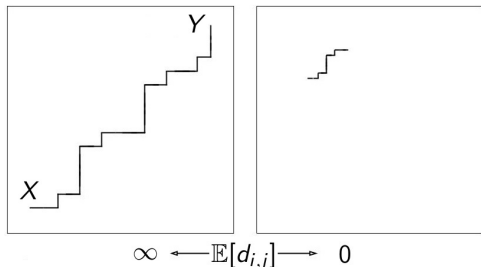
Free energies: $L_1^N(n) = \log Z^N(n)$

Using *geometric RSK* [Corwin-O'Connell-Seppäläinen-Zygouras '14] showed that L_1^N embeds as the lowest indexed curve in a line ensemble with a *nice Gibbs property* (uses inverse-gamma weights - very special)

Maximal free energy F_N

$$Z(X; Y) = \sum_{\pi: X \rightarrow Y} w(\pi), \quad w(\pi) = \prod_{(i,j) \in \pi} d_{i,j}$$

Maximal free energy: $F_N = \max_{(1,1) \leq X \leq Y \leq (N,N)} \log Z(X; Y).$



$$\mathbb{P}_d(\pi) = \frac{w(\pi)}{Z_d}, \quad Z_d = \sum_{\pi \in \square} w(\pi)$$

$$\mathbb{P}_d(\pi : X \rightarrow Y) = \frac{Z(X; Y)}{Z_d}$$

Why study F_N ?

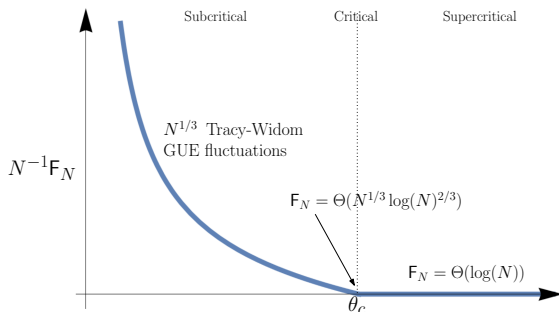
- 1 Interesting **phase transition** for F_N from $\theta < \theta_c$ to $\theta > \theta_c$ (next slide).
- 2 F_N is related to the smallest **singular value** of a **random operator on the honeycomb lattice** [Kotowski-Virág '19].
- 3 F_N is a proxy for studying a **free directed polymer path measure**.

Maximal free energy F_N

Theorem (Barraquand-Corwin-D. '21abc)

If $\theta_c = 2\Psi^{-1}(0) > 0$ (Ψ is the digamma function) then

- For $\theta < \theta_c$, $F_N + 2\Psi(\theta/2)N$ has order $N^{1/3}$ GUE Tracy-Widom fluctuations;
- For $\theta = \theta_c$, $F_N = \Theta(N^{1/3}(\log N)^{2/3})$;
- For $\theta > \theta_c$, $F_N = \Theta(\log N)$.



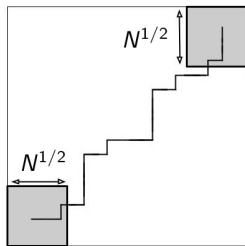
Theorem extends [Kotowski-Virág '19], which considers $\theta < \theta_c/2$.

Ideas behind the proof: subcritical case $\theta < \theta_c$

Maximal free energy: $F_N = \max_{(1,1) \leq X \leq Y \leq (N,N)} \log Z(X; Y)$.

Lower bound: $F_N \geq \log Z(1, 1; N, N)$ and $\log Z(1, 1; N, N) + 2\Psi(\theta/2)N$ has order $N^{1/3}$ GUE Tracy-Widom fluctuations [Borodin-Corwin-Remenik '13], [Krishnan-Quastel '18] and [Barraquand-Corwin-D. '21a].

Upper bound: $N^{-1/3}(F_N + 2\Psi(\theta/2)N)$.



Key ingredients:

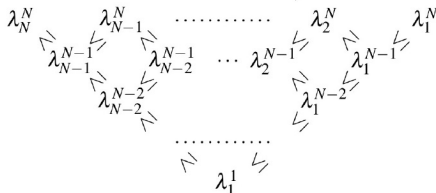
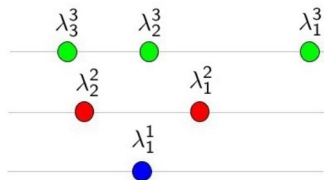
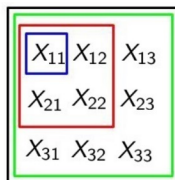
- **Moderate deviation estimates** for $\log Z(X; Y)$ from [Barraquand-Corwin-D. '21a]
- **Tightness** of $\log Z(X; \dots)$ and $\log Z(X; \cdot)$ [Barraquand-Corwin-D. '21b] and [D.-Wu '21]
- **Polymer structure**

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Gaussian Unitary Ensemble. $\{\xi_{ij}, \eta_{ij}\}_{i,j=1}^{\infty}$ are i.i.d. $N(0, 1)$ variables.

$$X_{ij} = \begin{cases} \xi_{ij} & \text{if } i = j \\ 2^{-1/2}(\xi_{ij} + \sqrt{-1}\eta_{ij}) & \text{if } i < j \\ 2^{-1/2}(\xi_{ij} - \sqrt{-1}\eta_{ij}) & \text{if } i > j \end{cases}$$



GUE-corners process: ($\beta = 2$ corners process). Gibbs property:
 $\mathbb{P}(\lambda^1, \dots, \lambda^{k-1} | \lambda^k) = \prod_{i=1}^{k-1} I(\lambda^i, \lambda^{i+1})$ with $I(\lambda^i, \lambda^{i+1}) = \mathbf{1}\{\lambda^i \preceq \lambda^{i+1}\}$

β -corners processes

Gaussian Orthogonal/Unitary/Symplectic Ensemble X_{ij} are real ($\beta = 1$), complex ($\beta = 2$) and quaternion ($\beta = 4$).

Continuous β -corners process:

$$\mathbb{P}(\lambda^1, \dots, \lambda^{k-1} | \lambda^k) = \prod_{i=1}^{k-1} I(\lambda^i, \lambda^{i+1})$$

$$I(\lambda^j, \lambda^{j+1}) = \mathbf{1}\{\lambda^j \preceq \lambda^{j+1}\} \prod_{1 \leq b < a \leq j} (\lambda_b^j - \lambda_a^j)^{2-\beta} \prod_{a=1}^j \prod_{b=1}^{j+1} |\lambda_a^j - \lambda_b^{j+1}|^{\frac{\beta}{2}-1}$$

Discrete β -corners process:

$$\mathbb{P}(\lambda^1, \dots, \lambda^{k-1} | \lambda^k) = \prod_{i=1}^{k-1} I(\lambda^i, \lambda^{i+1})$$

$I(\lambda^j, \lambda^{j+1}) = J_{\lambda^{j+1}/\lambda^j}(1)$ (skew) Jack symmetric function with $\theta = \beta/2$. When $\beta = 2$ then $J_{\lambda^{j+1}/\lambda^j}(1) = S_{\lambda^{j+1}/\lambda^j}(1) = \mathbf{1}\{\lambda^j \preceq \lambda^{j+1}\}$.

Discrete β -corners process are *integrable discretizations* of continuous β -corners processes. Special cases of **ascending Macdonald processes** [Borodin-Corwin '14]. Appear in distributions on **irreducible representations** [Bufetov-Gorin '18].

Discrete β -ensembles

Projecting β -corners processes to their top level we get

Continuous: $\mathbb{P}(\lambda_1, \dots, \lambda_N) \propto \prod_{1 \leq i < j \leq N} (\lambda_i - \lambda_j)^{2\beta} \prod_{i=1}^N e^{-V_N(\lambda_i)}$.

Discrete: $\mathbb{P}(\ell_1, \dots, \ell_N) \propto \prod_{1 \leq i < j \leq N} Q_\theta(\ell_i - \ell_j) \prod_{i=1}^N e^{-V_N(\ell_i)}$, $Q_\theta(x) := \frac{\Gamma(x+1)\Gamma(x+\theta)}{\Gamma(x)\Gamma(x+1-\theta)}$.

β -log gas and *Discrete β -ensembles* - [Borodin-Gorin-Guionnet '17]

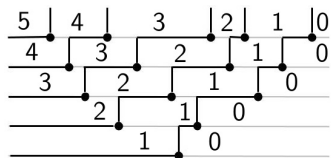
- **Law of large numbers:** $\mu_N = N^{-1} \sum_{i=1}^N \delta_{\ell_i/N}$ concentrate near μ_{eq} (**Wigner's Semicircle Law**)
- **Global Central Limit Theorem:** $\sum_{i=1}^N f(\ell_i/N) - \mathbb{E} \left[\sum_{i=1}^N f(\ell_i/N) \right]$ is asymptotically Gaussian [Borodin-Gorin-Guionnet '17]
- **Edge universality:** [Bourgade-Erdős-Yau '14],[Guionnet-Huang '19]
- **Edge large deviations:** [Johansson '98, '00], [Féral '08]

Theorem (Das-D. '21)

For general V_N 's the random variables ℓ_1/N satisfy a large deviation principle.

Upper tail rate is N and lower tail rate is N^2 . Rate functions are explicit but *different* from continuous β -log gases (due to **discreteness** of the model).

Global asymptotics for β -corners processes



β -corners process \implies height function $H(x, y)$

Conjecture: $H(x, y)$ converges to a suitable pullback of the *Gaussian Free field on \mathbb{H}* .

Known for *Wigner matrices* [Borodin '10].

Stieltjes transform: $G(z, s) = \sum_{i=1}^s \frac{1}{z - \ell_i^s}$.

[D.-Knizel '21]: *multi-level loop equations* or *Nekrasov equations* [Nekrasov '16] = functional equations relating joint cumulants of $G(z_1, s_1), \dots, G(z_k, s_k)$.

- 1 Generalize single level loop equations in [Borot-Guionnet '13] and Nekrasov equations in [Borodin-Gorin-Guionnet '17]
- 2 [D.-Knizel '19]: $G(z_1, N), G(z_2, N - 1), N^{1/2}[G(z_3, N) - G(z_3, N - 1)]$ have joint Gaussian limits.

Explanation: Top two levels converge to the same 1D slice of the 2D GFF, and their difference to a certain directional derivative of the GFF. The analogue for *Wigner matrices* is proved in [Erdős-Schröder '18].

Gibbs properties + integrable (algebraic) input + analytic tools

- KPZ universality for Gibbsian line ensembles
- Asymptotics for polymer models
- Asymptotics for β -corners processes
- (Did not discuss) Convergence of six-vertex models to the GUE corners process [D. '18], [D.-Rychnovsky '20]

Thank you!

